

Very low frequency oscillations of heat load and recycling flux in steady-state tokamak discharge in TRIAM-1M

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Plasma wall interaction (PWI) driven relaxation oscillations are investigated in the steady state discharge for 5 hours. The oscillation frequency was about 10^{-3} Hz and each perturbation lasted for ~ 300 s. The heat load, recycling flux and impurity influx were varied from a few % to several tens of % . The largest variation of 70% was seen on the Mo XIII, although the influx of Mo I was only 20 % . Although the input rf power is kept constant during the discharge, the coupling between the rf and plasma was increased by $\sim 10\%$. The current drive efficiency is decreased by 24 % in spite of current ramp. The toroidal and poloidal profiles of the recycling flux were also changed. During the last relaxation phase, the plasma was finally terminated. The current reduction (> 4 kA) was not recovered by intense local perturbation of the recycling superposed on the relaxation oscillation.

1. Steady State Tokamak Operation

With the goal of addressing the critical issues of steady-state tokamak operation (SSTO) [1] of future fusion devices, experiments aiming at ‘day long operation at high performance’ have been carried out in a limiter configuration of TRIAM-1M plasma ($R_0=0.84$ m, $a\sim 0.12$ m and $B=6-7$ T). The plasma was maintained by lower hybrid wave LHW whose power of ~ 10 kW and the density was $1 \times 10^{18} \text{ m}^{-3}$. The record value of the discharge duration was updated to 5 h and 16 min (Fig. 1). In order to investigate the main PWI and its effect on SSTO from a view point of heat load, particle recycling and impurity influx distributions, a rail limiter (movable limiter; ML) was inserted vertically into the plasma and thus the main contribution of the PWI could be localized on this. Other plasma facing components PFCs consist of three D-shaped limiters PLs, divertor plates DPs, vacuum vessel, the cover of bellows, and LH antenna. The

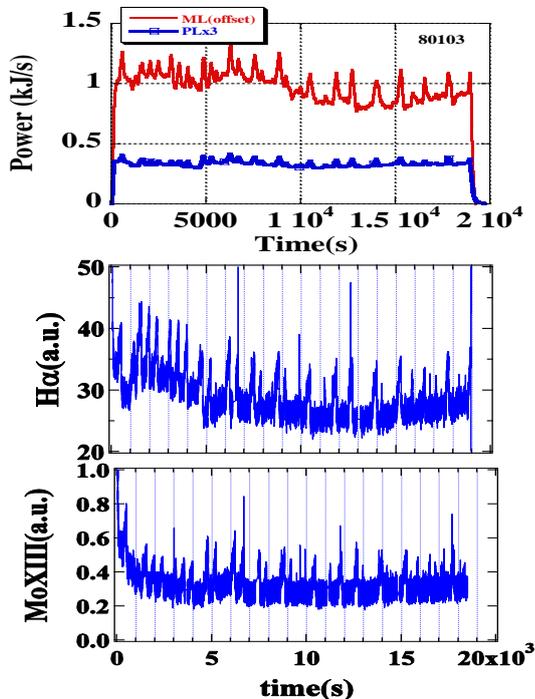


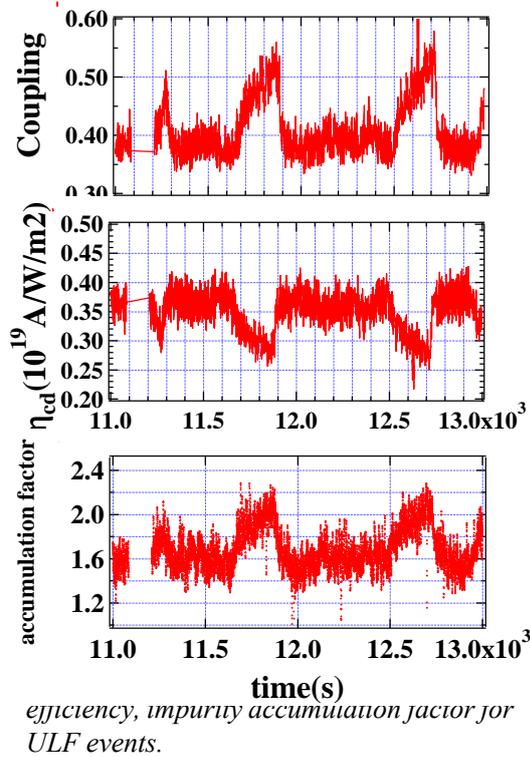
Fig. 1 Heat load, recycling flux and Mo XIII in five-hour discharge. Perturbations with $1-2 \times 10^{-3}$ Hz are seen on these signals.

ML and the one of PLs are located at the same port. The ML, PLs and DPs are made from molybdenum, and others are stainless-steel. The closest distances from the LCFS determined by the ML are ~ 5 mm for the inner mid-plane part of PL, ~ 24 mm for outer one, ~ 46 mm for DPs, 22 mm for inner wall, ~ 41 mm for outer one, and ~ 46 mm for LH antenna, respectively.

Heat and particle recycling fraction were determined by calorimetric measurement and H_α arrays around the torus[2,3]. A fraction of 34 % of the total heat load was deposited on the ML and the rest was distributed on PLs (10%), and wall (34%) and DPs (22%). The fraction of the recycling particle rate is $\sim 40\%$ from ML, 30% from PLs, and 30% from the wall and DPs. Thus the main heat load and particle flux are deposited mainly on the small area (a few cm^2) of the ML (3 cm in width and 16 cm in length along the major radius). The main metal impurity ions are Mo. Visible (Mo I 386 nm) and VUV (Mo XIII 34.1nm) spectrometers and soft X-ray array are used to investigate the impurity influx and accumulation in the core region. The influx of Mo I is determined $\sim 3 \times 10^{16}$ Mo atoms/ $\text{m}^2 \text{ s}$. The surface and dust temperatures are measured by IR camera at $3.5 \mu\text{m}$ and a near IR spectrometer from

0.9 μm to 1.65 μm . By assuming bulk and dust temperatures (T_b and T_d) and surface coverage factor of dust for the viewing area, the observed near IR spectrum shows that T_d can reach the melting temperature of Mo and dusts cover a few % of the viewing area. During the discharge, the frequency of Mo spikes measured by influx monitor and soft X-rays increases toward the end. This shows that PWI during the discharge is not constant in time [3].

2. Ultra Low Frequency events



As shown in Fig. 1, PWI driven relaxation oscillations (ULF events) are observed on heat load, recycling flux and impurity behaviour. These oscillations are semi-regular and should be distinguished from the random off-normal events, for example above mentioned impurity bursts. We are concerned with this oscillation because the 5 hour discharge has terminated during the last event (§3). ULF events are characterized by slow variations and quick decay. “Slow phase” lasts for 150 s ~ 200 s and “quick decay phase” ~ 100 s. ULF events are repeated at the frequency of $1-2 \times 10^{-3}$ Hz. The unignorable power loss through the antenna made of stainless steel is taken into account to evaluate the net coupled power and it is observed that the coupling between the incident LHW power and the plasma maintains ~ 40 % whole duration, but increases to 50 % during the ULF event. As a consequence I_p and hard X-ray intensity integrated above 30 keV increase by ~ 3%. Thus, the current drive efficiency ($\eta_{CD} = I_p R n_e / P$) with an assumption of the constant density shows a reduction from 0.37 to 0.28×10^{19} A/Wm². This unfavourable aspect is considered to be due to the impurity

accumulation phenomenon. Here the accumulation factor, shown in Fig. 2, is defined as a ratio of MoXIII to Mo I, and as a consequence of the relatively large increment of Mo XIII it denotes that the injected impurity neutrals are accumulated during ULF events and lead to reduction in η_{CD} by enhanced collisions. The core impurity transport change during the ULF event, which is considered to be triggered by enhanced PWI, is under study.

A calorimeter system with 32 thermistors for cooling water temperature of PFCs were used to evaluate the distribution of the heat load deposited on the PFCs. The heat load on the ML starts to increase linearly at the rate of $\sim 1.5 \pm 0.2$ W/s when the ULF event occurs, and then decays at the same or lower rate. This aspect is

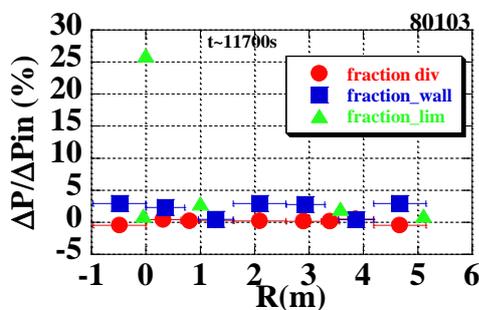


Fig. 3 The fraction distribution of the variation in heat load deposited on PFCs during the ULF event.

different from those of the impurity or recycling flux. When we assume the interaction area of ~ 3 cm² for ML from CCD view the heat flux increases from 2.8 MW/m² to 3.3 MW/m² at the event. The heat load on other PFCs changes in similar manner at various changing rate. From the heat load measurement the heat fluxes $\langle q \rangle$ averaged on their interactive surface area are evaluated on each PFCs. The total heat load on the wall is the summation of the six part of the wall. The horizontal bar indicates each covering area. There seems an n=2-like toroidal structure.

Although the covering areas of the two parts ($Z \sim 1.2\text{m}$ and $Z \sim 3.8\text{m}$) of the wall are smaller than others, the $\langle \Delta q(Z) \rangle$ distribution also shows the similar structure. This structure remained for ULF events. The total power change ΔP_{in} in event is evaluated of $\sim 1\text{ kW}$ from the net power evaluation, since the lost power in the antenna is decreased during ULF events. The summing up the total power change in all PFCs gives 0.44 kW , half the ΔP_{in} . The fraction of the heat load deposited on PFCs during the ULF events is shown in Fig.3. Although the missing power exists, the largest fraction was still localized on the ML.

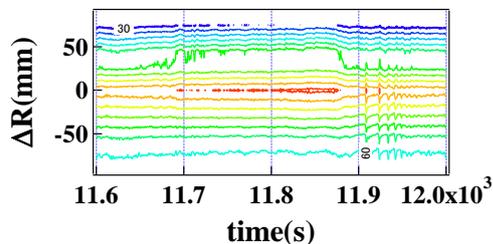


Fig. 4 The contour plot of $H_\alpha(\Delta R)$ during the ULF event.

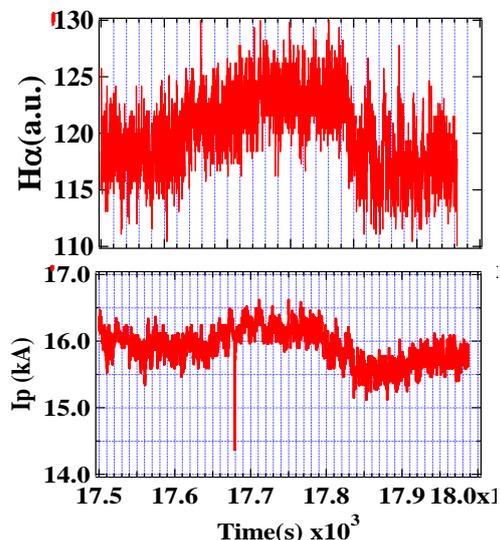


Fig.5 The one cycle of ULF event. Slow rise, quick decay, recovery phases are shown for H_α and I_p .

3. Termination of the steady state discharge

The 5 hour discharge was stopped at $t = t_{\text{end}} (=19873.8\text{ s})$ without a sign of termination. The last ULF event starts $\sim -150\text{ s}$ before t_{end} and seems to continue similar to the previous ones. The heat load on the PFCs increase and the H_α also increases. Both $H_\alpha(Z)$ and $H_\alpha(\Delta R)$ profiles behave similarly. These behaviors are within the variations of the previously repeated ones. An intense perturbation occurs $t_1 = -2.1\text{ s}$ before t_{end} , as shown in Fig. 6. The current dropped by $\sim 25\%$ within 0.2 s , but it could be recovered to the previous level of 16 kA within continuing 0.8 s . However, the second intense perturbation may trigger the current quench at $t_2 = -1.1\text{ s}$. The current decayed gradually to $\sim 10\text{ kA}$ for 1 s at the averaged quench rate of $dI/dt = -6\text{ kA/s}$ and suddenly disrupted within 0.1 s . There is a different point in the current decay, namely for the first 0.2 s at both perturbations, the former decay rate is -20 kA/s and the latter one is only 1 kA/s . At the first perturbation the intense changes in the recycling fluxes are observed at the same time, but they delay by 0.2 s or 0.3 s with respect to the time at the current decays. Although the main recycling flux is supplied from the ML, the changes in the recycling flux are observed simultaneously around the torus. Figure 7 shows the $\Delta H_\alpha(Z)$ and $\Delta H_\alpha(\Delta R)$ profiles, the changed quantities from the averaged H_α levels before the ULF event is triggered. For reference the profiles at the ULF event ($t = -3.8\text{ s} \sim -2.8\text{ s}$) are shown. Outward shift

An H_α measurement system consists of fiber optics, narrow H_α filters and photomultipliers is used to measure toroidal (six positions) and poloidal (seven vertical chords) distributions. The ULF events are analyzed for the toroidal $H_\alpha(Z)$ ($Z = \text{a circumference of the torus}$) profile at $\Delta R = 0$ and horizontal $H_\alpha(\Delta R)$ profile at $Z = 3.5\text{m}$. As shown in fig.4, during the ULF event H_α increases in the longer R side and decreases in the shorter side, that is, $H_\alpha(\Delta R)$ shifts outward, which is favorable for rf coupling. The variation of H_α is 8% at $\Delta R = 0$ and 35% at $\Delta R = 70\text{ mm}$. After a lapse of $\sim 200\text{ s}$ from the trigger time of ULF event $H_\alpha(\Delta R)$ shifts back to the previous one within $\sim 50\text{ sec}$, but after that several bursts are seen. Their amplitudes are larger than the slow variation and decay in time for $\sim 50\text{ s}$. The each burst increases with a time constant of $\sim 0.2\text{ s}$ and then decays with 0.7 s . The $H_\alpha(\Delta R)$ shifts inward at the burst, in the opposite direction compared to the slow variation. The $H_\alpha(Z)$ shows a highly localized profile near the ML, but the variation seems to be uniform around the torus, as shown in Fig. 7(a). In summary it is noted that the ULF event consists of three phases, slow increase and rapid decay following the relaxation phase. Namely, during the rapid decay phase the H_α and I_p are below the averaged ones, and therefore after that they recover gradually during the H_α bursts phase. This cycle is shown in Fig. 5. These cycles are repeated as ULF events.

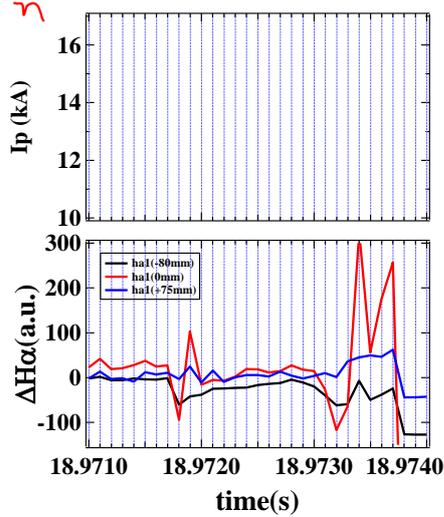


Fig.6 The two intense events ~ 2 s before the termination. The I_p drop, and relative variation of H_α at the ML are shown.

of $\Delta H_\alpha^9(\Delta R)$ and $\Delta H_\alpha^1(\Delta R)$ (not shown) and the almost uniform increment in H_α are similar to the above mentioned ones. At the first perturbation, $\Delta H_\alpha^9(\Delta R)$ shows a strong in-out asymmetry but $\Delta H_\alpha^1(\Delta R)$ does not show such one, indicating that the enhanced PWI in the outboard side occurs locally at the particular PFCs. The toroidal structure is quite complicated. At the ML the $\Delta H_\alpha(Z)$ is negative at t_1 and then becomes positive. $\Delta H_\alpha(Z)$ at the PL is always positive. The negative variation indicates the reduction in PWI. At the second perturbation above mentioned aspects in both profiles of $\Delta H_\alpha^9(\Delta R)$ and $\Delta H_\alpha(Z)$ are the same as the first one, but the amplitudes grows much resulting in final current quench.

Summary

The PWI driven relaxation oscillations and the termination of the discharge were investigated from a view point of SSTO. The 30~40 % of heat load and recycling particle flux are localized on the quite small area of the PFCs. The perturbations on heat load and particle recycling are repeated at the frequency of $1-2 \times 10^{-3}$ Hz and the each event lasts for ~ 300 s. During the slow rise phase in the last event, intense two perturbations occurred at the particular location and the plasma could not be survived.

Acknowledgement

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References

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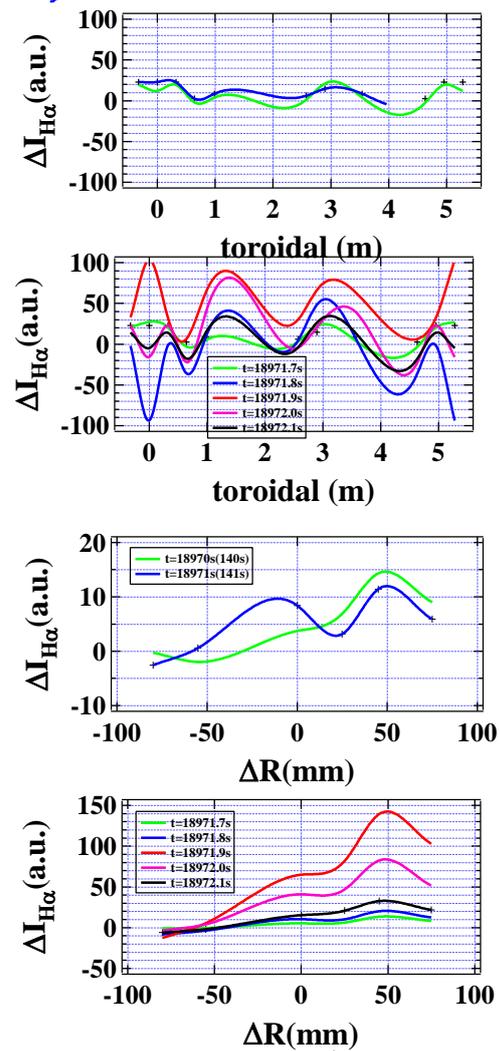


Fig.7 $\Delta H_\alpha(Z)$ (b) and $\Delta H_\alpha^9(\Delta R)$ (d) at the termination phase. For reference, profiles during the slow variation phase are shown (a) and (c).