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Coordinated observation of field line resonance in the mid-tail

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Abstract. Standing Alfvén waves of 1.1 mHz (~15 min in period) were observed by the Cluster satellites in the mid-tail during 06:00–07:00 UT on 8 August 2003. Pulsations with the same frequency were also observed at several ground stations near Cluster's footprint. The standing wave properties were determined from the electric and magnetic field measurements of Cluster. Data from the ground magnetometers indicated a latitudinal amplitude and phase structure consistent with the driven field line resonance (FLR) at 1.1 mHz. Simultaneously, quasi-periodic oscillations at different frequencies were observed in the post-midnight/early morning sector by GOES 12 ($l_0 \approx 8.7$), Polar ($l_0 \approx 11$ –14) and Geotail ($l_0 \approx 9.8$). The 8 August 2003 event yields rare and interesting datasets. It provides, for the first time, coordinated in situ and ground-based observations of a very low frequency FLR in the mid-tail on stretched field lines.

Keywords. Magnetospheric physics (Magnetotail; MHD waves and instabilities; Plasma waves and instabilities; Solar wind-magnetosphere interactions)

1 Introduction

Research on field-line resonance (FLR) has a long and rich history from theoretical, observational and numerical simulation perspectives. In 1954, Dungey (1954) considered the possibility that ultra-low frequency (ULF) waves might constitute standing Alfvén waves along geomagnetic field lines. Samson et al. (1971) reported the first observational evidence

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of FLR and found the local amplitude maximum and polarization change across the resonance. Chen and Hasegawa (1974) and Southwood (1974) both independently provided a theoretical description on how compressional MHD waves couple to shear Alfvén waves, giving rise to FLRs. By the 1980s, the concept of a cavity mode wave was introduced by Kivelson and Southwood (1985; 1986) to explain the quantization of compressional wave spectrum. More theoretical work was carried out using different geometries during the same decade (e.g. Allan et al., 1986; Lee and Lysak, 1989) as different geometries (boundary conditions) allow different wave mode structures. Theories on FLR were further developed and refined during the 1990s, as seen in the proposal of waveguide mode (Samson et al., 1992a), numerical studies including non-linear effects (e.g. Rankin et al., 1993), two-fluid studies (e.g. Streltsov and Lotko, 1996), waveguide excitation by magnetosheath flows (e.g. Mann et al., 1999), tail waveguide theory (e.g. Allan and Wright, 1998) and so on. There have also been numerous FLR studies based on observations that either use ground data or satellite measurements or all the data combined (e.g. Singer and Kivelson, 1979; Ruohoniemi et al., 1991; Walker et al., 1992; Takahashi et al., 1996; Laakso et al., 1998; Ohtani et al., 1999; Mann et al., 1998; 2002; Keiling et al., 2003).

FLRs are usually perceived as predominantly a dayside phenomenon with a preferential occurrence in the dawn-side. However, HF radar, ground magnetometer and optical emission observations have also shown that ULF pulsations resulting from FLRs in the millihertz range (especially the 1–4 mHz range) are a common feature of data in the nightside sector (e.g. Samson et al., 1991, 1992b and references therein). The link between FLRs and certain classes

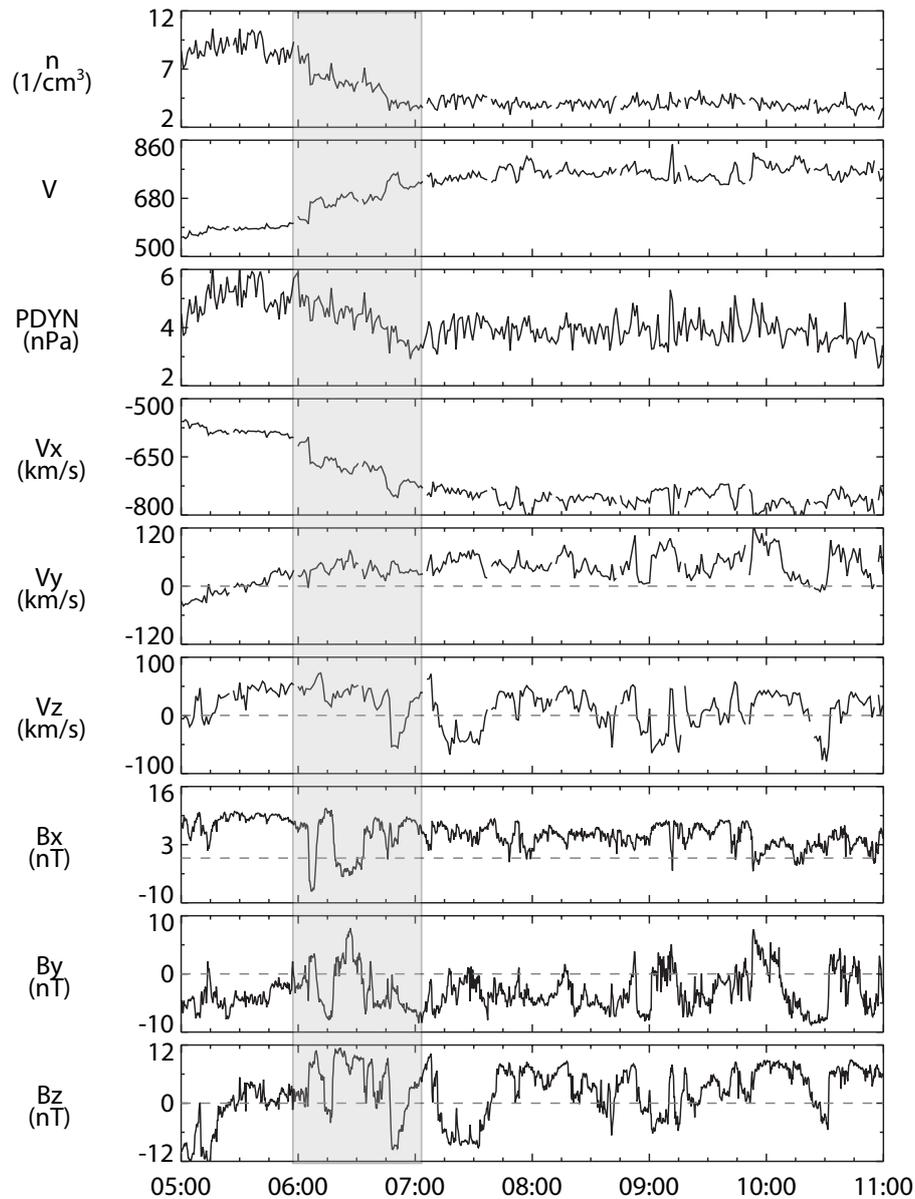


Fig. 1. Solar wind and IMF conditions measured by ACE spacecraft (it has been time-shifted by 38 min) during the interval 05:00–11:00 UT. From top to bottom are: density in cm^{-3} , solar wind speed in km/s, solar wind dynamic pressure in nPa, three components of solar wind velocity in km/s, and three components of the IMF in nT. The highlighted region in Fig. 1 indicates the interval during which the ULF pulsations were seen. GSM coordinates are used here.

of auroral arc formation has been reported and is gaining more support both observationally (e.g. Samson et al., 1992b; Lotko et al., 1998) and theoretically (e.g. Lui and Murphree, 1998; Streltsov and Lotko, 1999; Rankin et al. 1999).

FLRs can provide a self-consistent and comprehensive explanation of energy accumulation near magnetic surfaces where the frequency of global scale compressional waves matches the local field line eigenfrequency. FLRs can also account for density cavities in the low-altitude auroral acceleration region by nonlinear ponderomotive forces (Rankin et

al., 1995) and the narrow structuring of auroral arcs through mode conversion to inertial Alfvén waves. They play an important role in understanding ULF pulsations in the magnetosphere and also serve as an important diagnostic tool in monitoring plasma mass density. For reviews on FLR research, see Takahashi (1998) and Glassmeier et al. (1999).

Among the numerous publications, few studies of FLRs have been reported using data in the magnetotail. In this paper we present, for the first time, convincing in situ evidence of FLRs of a very low frequency (1.1 mHz) in the

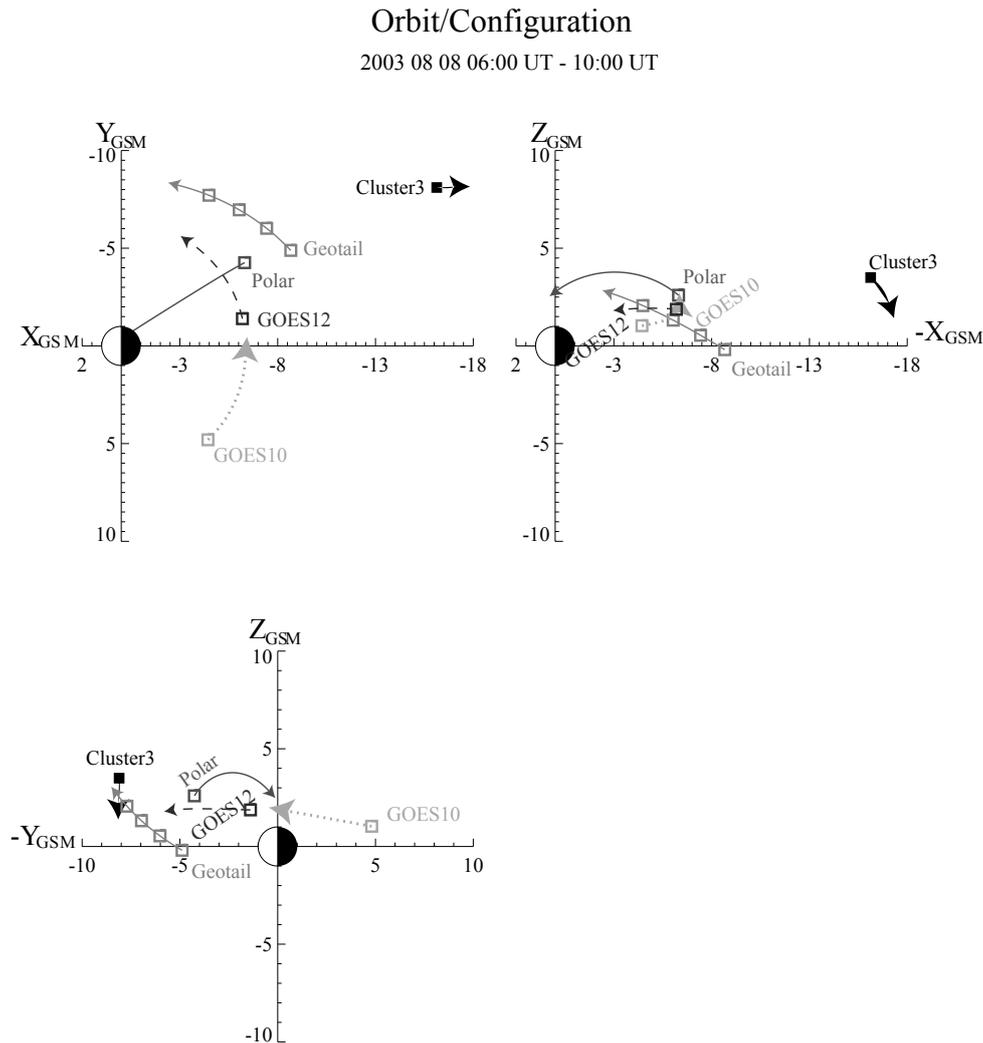


Fig. 2. The orbit configuration in three different planes (x-y, x-z and y-z) in GSM during 06:00–10:00 UT for all the satellites involved in this study, with the arrows indicating the motion.

magnetotail at $l_0 \approx 21 R_E$ (equatorial radial distance using Tsyganenko 96 model) observed by Cluster satellites. This event took place during the interval 06:00–07:00 UT on 8 August 2003. The expected latitudinal FLR amplitude and phase profile were also observed by ground magnetometer stations in the CANOPUS (now CARISMA) array near Cluster’s footpoint, which provides further support to our interpretation of the event. The 1.1-mHz pulsations were also observed by some of the ground stations of the Greenland Chains. Simultaneously, quasi-periodic oscillations at higher frequencies were registered by GOES 12, Polar and Geotail.

The organization of the paper is as follows: Sect. 2 provides the solar wind condition measured by ACE spacecraft. Detailed observations of ULF pulsations at different frequencies are given in Sect. 3, including data from several satellites and ground magnetometers. Discussion is provided in Sect. 4 and a summary is given in Sect. 5.

2 Solar wind condition

The magnetosphere was subject to a high-speed solar wind stream with an average speed of about 680 km/s, an average density of 4 cm^{-3} and dynamic pressure of 4 nPa during the event. The IMF (interplanetary magnetic field) B_z had an average value of 3.8 nT. Both the IMF and solar wind velocity were oscillatory. Figure 1 shows the detailed measurements from the ACE spacecraft (it has been time-shifted by 38 min to the Cluster location in the mid-tail) during the interval 05:00–11:00 UT. From top to bottom are: density in cm^{-3} , solar wind speed in km/s, solar wind dynamic pressure in nPa, three components of solar wind velocity in km/s, and three components of the IMF in nT. The highlighted region in Fig. 1 indicates the interval during which the ULF pulsations were seen. The Geocentric Solar Magnetospheric (GSM) coordinate system is used here.

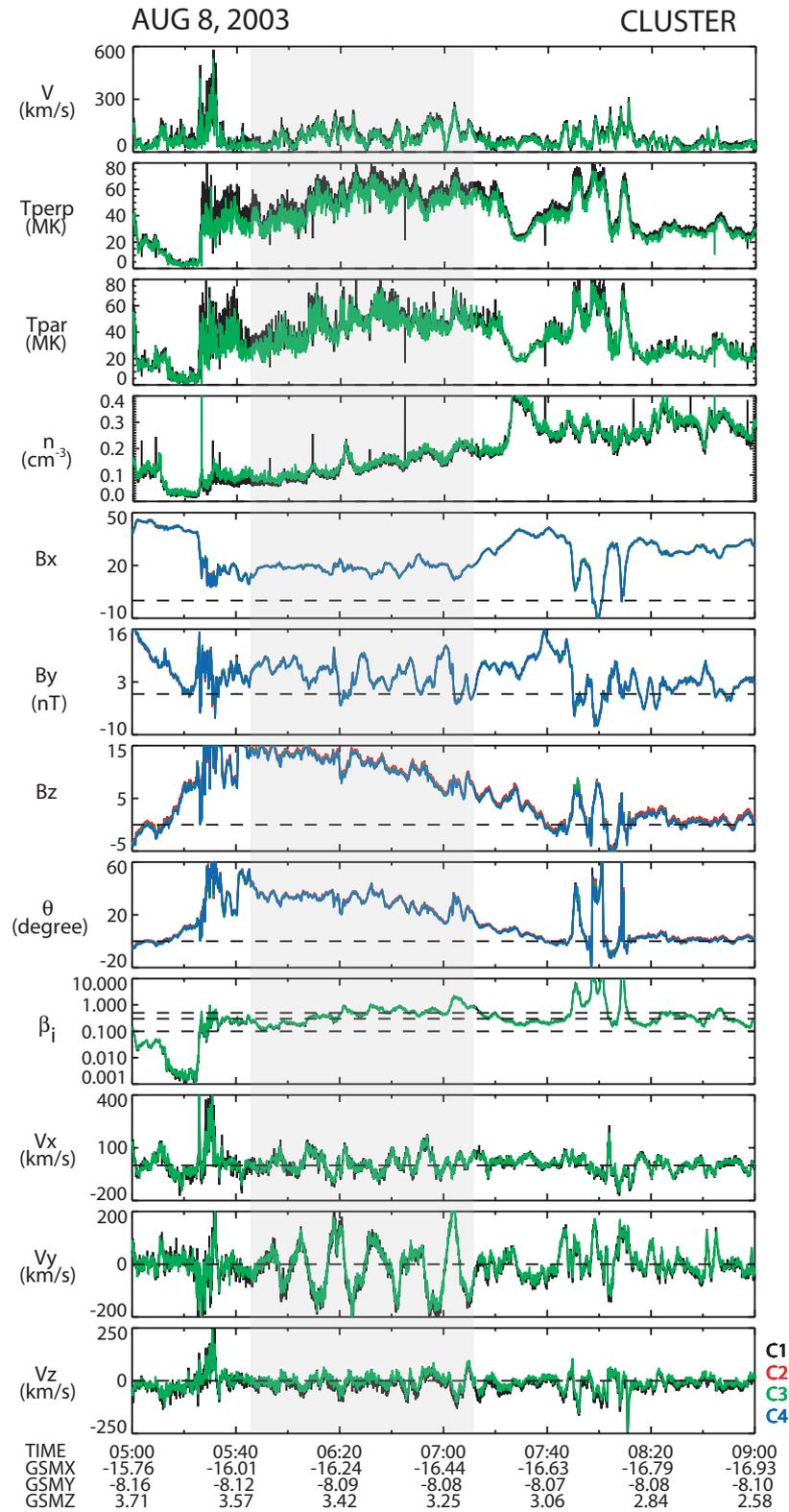


Fig. 3. Cluster data during 05:00–09:00 UT, including both ion moment data and magnetic field data. The highlighted region is the interval during which quasi-periodic oscillations were seen. Figure 3 serves as an observational context for the oscillations.

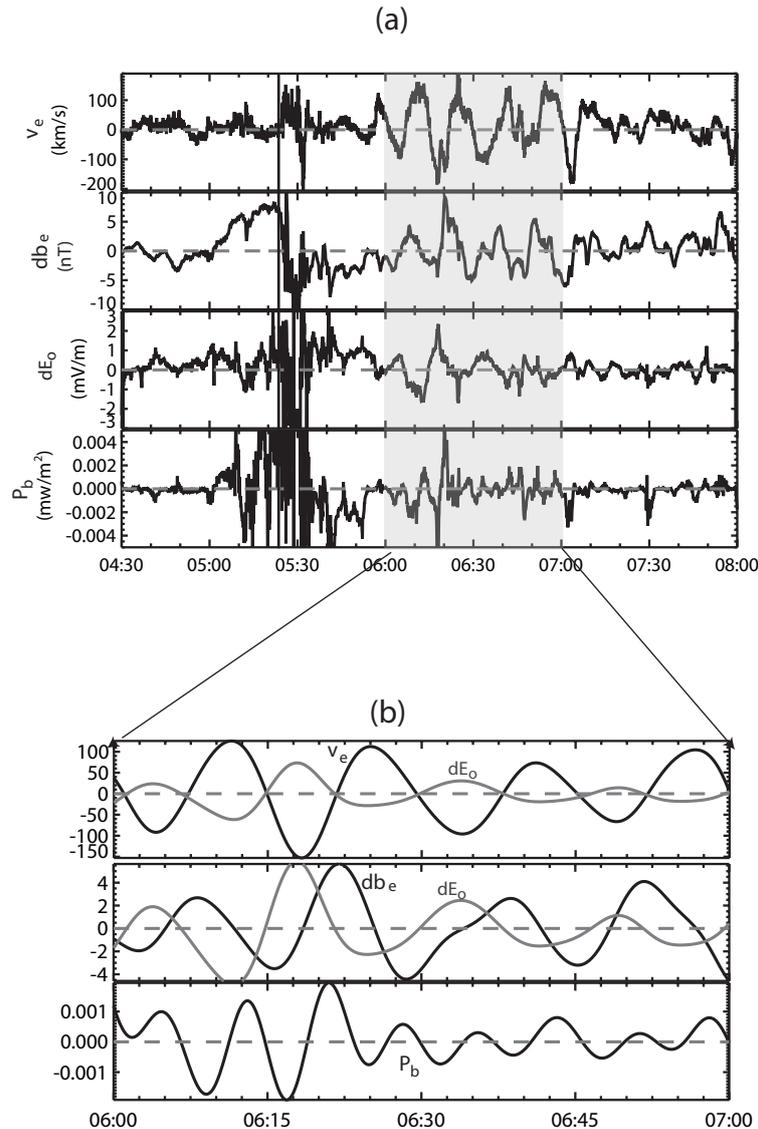


Fig. 4. Parameters of oscillations seen on 8 August 2003 by Cluster 3 in magnetic field-aligned coordinates (see text for details) with (a) showing the unfiltered data and (b) showing an expanded view of the event with a band-pass filter [0.4 mHz, 2.5 mHz] applied.

3 Observations

3.1 Orbit configuration

Figure 2 shows the orbit configuration in three different planes (x-y, x-z and y-z) in GSM during 06:00–10:00 UT for all the satellites involved in this study, with the arrows indicating the moving direction. We can see that Cluster was located further down in the tail, while GOES 10, GOES 12, Polar and Geotail were closer to the Earth. It is worth noting that only GOES 10 was located in the pre-midnight, while all the other satellites were in the post-midnight.

3.2 Oscillations at 1.1 mHz

The 1.1-mHz quasi-periodic oscillations were observed by Cluster in the mid-tail and some of the ground magnetometer stations in the CANOPUS and Greenland chains around Cluster’s footprint.

3.2.1 Cluster in-situ measurements

Cluster was located in the plasma sheet near 02:00 magnetic local time (MLT). Due to the small separation (<300 km) between the four Cluster spacecraft, the measurements from the four spacecraft are almost identical. Here we use Cluster 3 (C3) as the reference spacecraft. Figure 3 shows Cluster data during 05:00–0900 UT, including both ion moment data

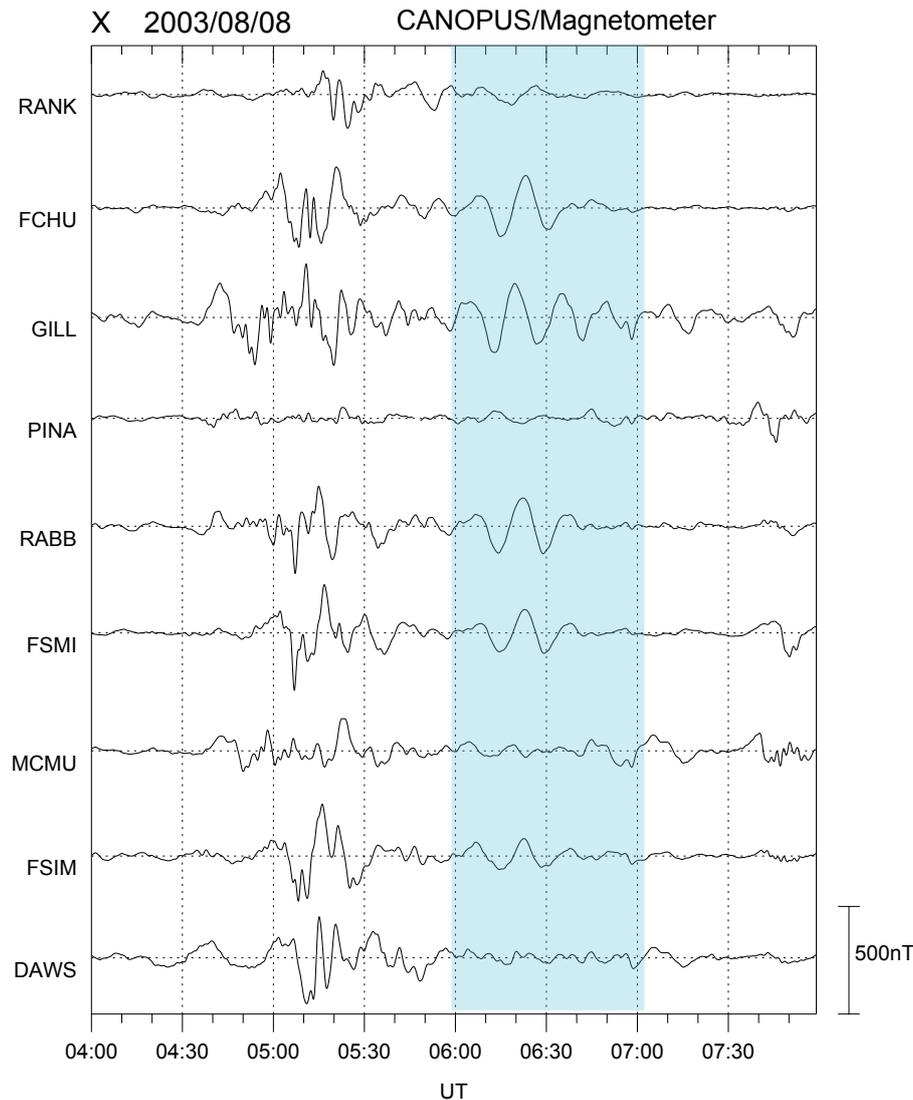


Fig. 5. The band-pass (with cut-off periods of 100–2000 s) filtered data in the X (geographic north) component from the CANOPUS chains (Here X is very close to H, local magnetic north, as the angle between them is only about 3°). Very monochromatic 1.1-mHz oscillations with large amplitudes were detected by several CANOPUS stations.

from the HIA (hot ion analyzer) component of the CIS (Cluster Ion Spectrometry) (Rème et al., 2001) instrument and magnetic field data from the FGM (FluxGate Magnetometer) measurements (Balogh et al., 2001). The highlighted region is the interval during which quasi-periodic oscillations were seen. Figure 3 serves as an observational context for the oscillations. Cluster traversed several different plasma regions during the four-hour interval mainly because of the dynamical changes in the magnetotail: the lobe region (05:14–05:24 UT, due to plasma sheet thinning), plasma sheet boundary layer (05:24–05:45 UT), plasma sheet region (05:45–08:00 UT), the region of neutral sheet crossing (\sim 07:50–08:20 UT). A dipolarization event took place around 05:25 UT accompanied by relatively large flows.

Quasi-periodic oscillations were most noticeable in V_y and B_y with a period of 15 min and frequency at 1.1 mHz during 06:00–07:15 UT. C3 started from $(-16.13, -8.11, 3.50) R_E$ in GSM coordinates at 06:00 UT, going towards the equator. During the 1-h interval, Cluster's position did not change much, with C3 at $(-16.44, -8.08, 3.25) R_E$ at 07:00 UT. The average plasma ion beta during the one-hour interval was about 0.5, ion density was $0.1\text{--}0.2\text{ cm}^{-3}$, the average ion temperature was about 5 keV and the average angle between the magnetic field direction and the equatorial plane is $\sim 30^\circ$.

Figure 4 shows the oscillating physical quantities in field-aligned coordinates, with (a) containing unfiltered measurements at 4-s resolution and (b) presenting expanded views of the event with a band-pass filter between 400 s and 2400 s

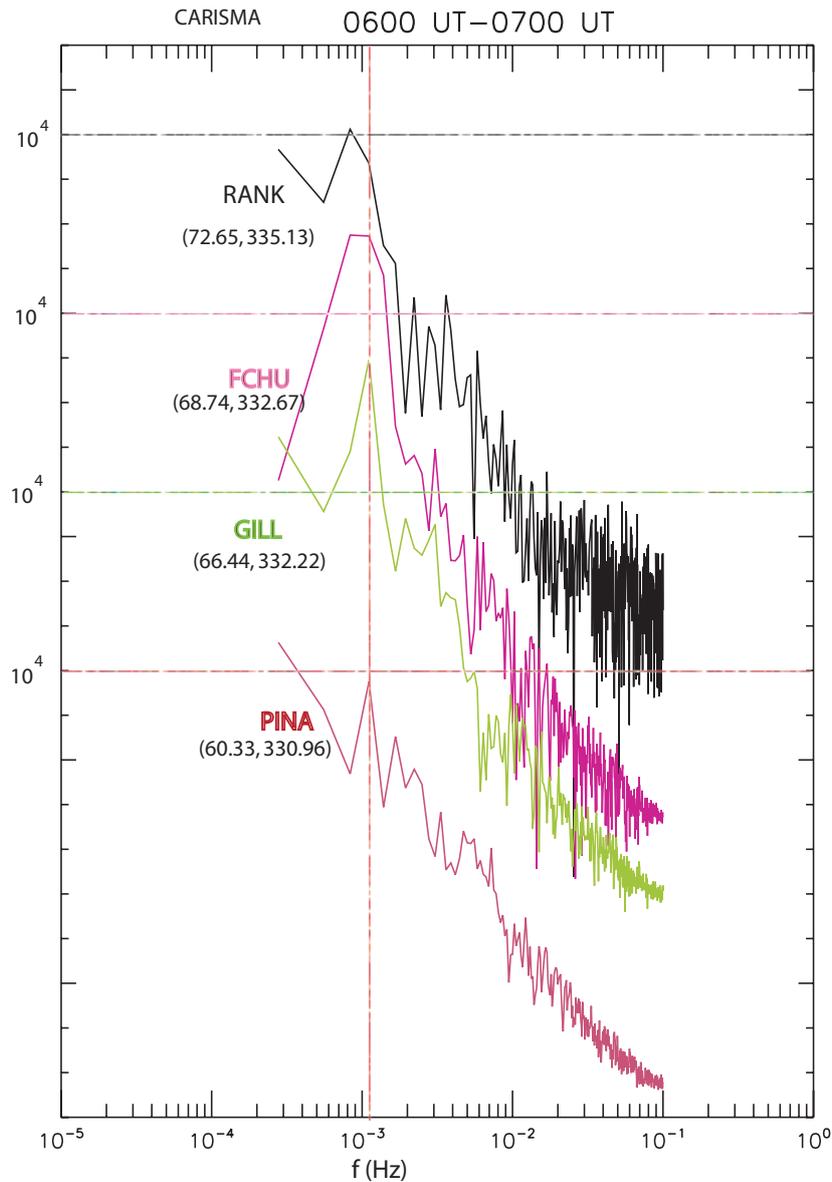


Fig. 6. The power spectra along the north-south Churchill line using the data within the one-hour interval 06:00–0700 UT.

(0.4–2.5 mHz). In the field-aligned coordinates, “o” is in $-\mathbf{B}_0 \times (\mathbf{B}_0 \times \mathbf{r})$ direction – approximately radially outward, “e” is in the direction of $(\mathbf{B}_0 \times \mathbf{r})$ – approximately eastward and ‘b’ is in the direction of the magnetic field \mathbf{B}_0 , where \mathbf{B}_0 is the 1-h running average of the measured magnetic field and \mathbf{r} represents the spacecraft’s location.

The top panel of Fig. 4a shows the azimuthal ion velocity V_e in km/s with its varying amplitude up to 150 km/s, the second panel is the perturbed azimuthal magnetic field db_e in nT with its amplitude up to 10 nT, the third panel shows the perturbed radial electric field dE_o in mV/m (an amplitude of 3 mV/m) and the bottom panel is the magnetic field-aligned Poynting vector P_b (from the two components dE_o and db_e ,

the contribution from $-d\mathbf{E}_e \times d\mathbf{b}_0$ being negligible). The electric field data are obtained from the EFW (Electric Fields and Waves) (Gustafsson et al., 1997) instrument. Since the EFW instrument only has two axes, $\mathbf{E} \cdot \mathbf{B} = 0$ is assumed in obtaining the electric field. The field-aligned Poynting vector oscillates at about twice of the FLR frequency with its amplitude up to $4 \mu\text{W}/\text{m}^2$. The shaded region is to show the interval when FLR oscillations were seen.

The top panel of Fig. 4b shows filtered V_e (in black) and dE_o (in grey) with dE_o being multiplied by a factor of 50. The second panel is the overlapping plot of db_e (black) and dE_o (grey), where dE_o was multiplied by 4. The last panel shows the field-aligned Poynting vector resulted from the

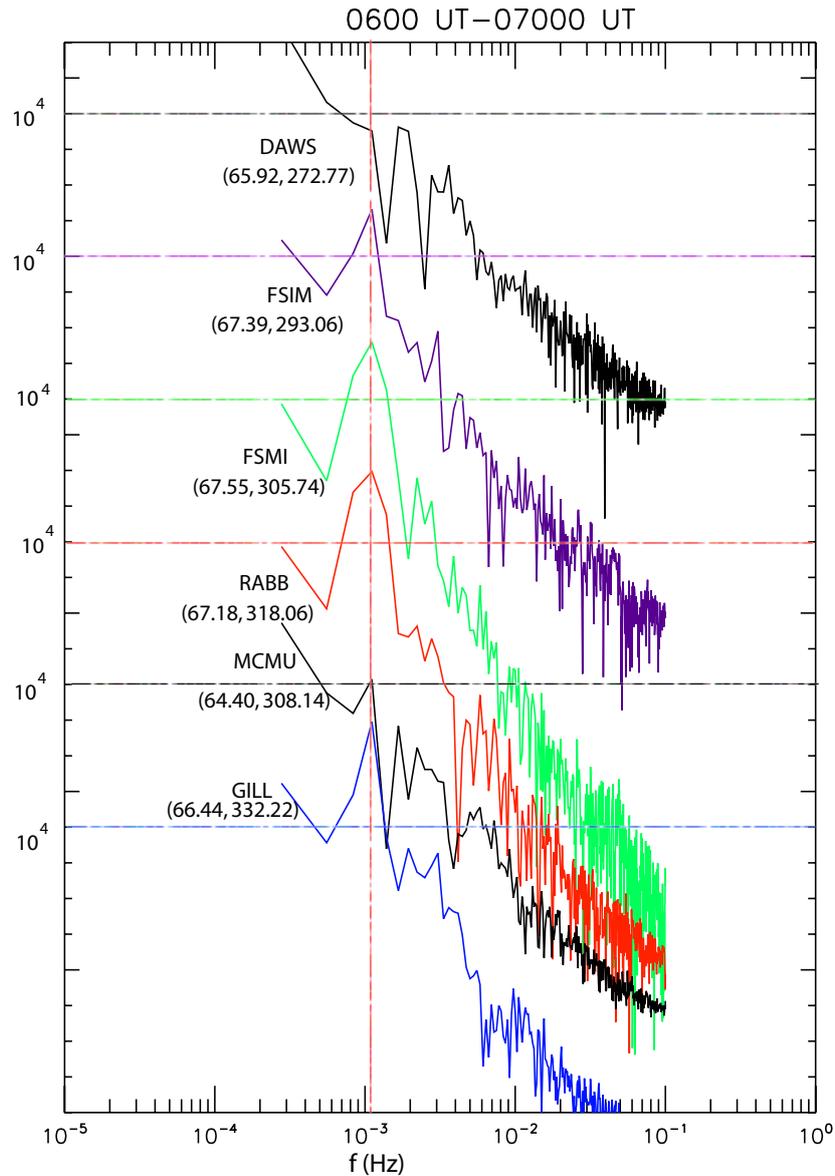


Fig. 7. The power spectra for all the stations along the west-east direction with the same format as Fig. 6.

filtered fields in the top two panels. That V_e and dE_0 are 180° out of phase is to be expected. The fact that db_e and dE_0 are 90° out of phase and the field-aligned Poynting flux oscillates at twice of the FLR frequency confirms that these oscillations are standing Alfvén waves (Fenrich et al., 1995). The oscillations belong to the toroidal mode as the oscillations are mainly in V_e , db_e and dE_0 . The cross phase between db_e and V_e is 90° , indicating the observed wave is a fundamental mode and the spacecraft was located above (North to) the equatorial node of B. Therefore quasi-periodic oscillations at 1.1 mHz observed by Cluster are proven to be toroidal standing Alfvén waves.

3.2.2 The 1.1 mHz ground-based observations

The 1.1-mHz oscillations were also seen on the ground across most of the CANOPUS and Greenland magnetometer arrays.

Figure 5 shows the band-pass filtered (with cut-off periods: 100–2000 s) data in the X (geographic north) component from the CANOPUS chains (Here X is very close to H, local magnetic north, as the angle between them is only about 3°). Very monochromatic 1.1-mHz oscillations (were also seen in the raw data) with large amplitudes were detected by several CANOPUS stations, such as Gillam (GILL), Fort Churchill (FCHU), Rabbit Lake (RABB), Fort Smith (FSMI) and Fort Simpson (FSIM), with GILL having the

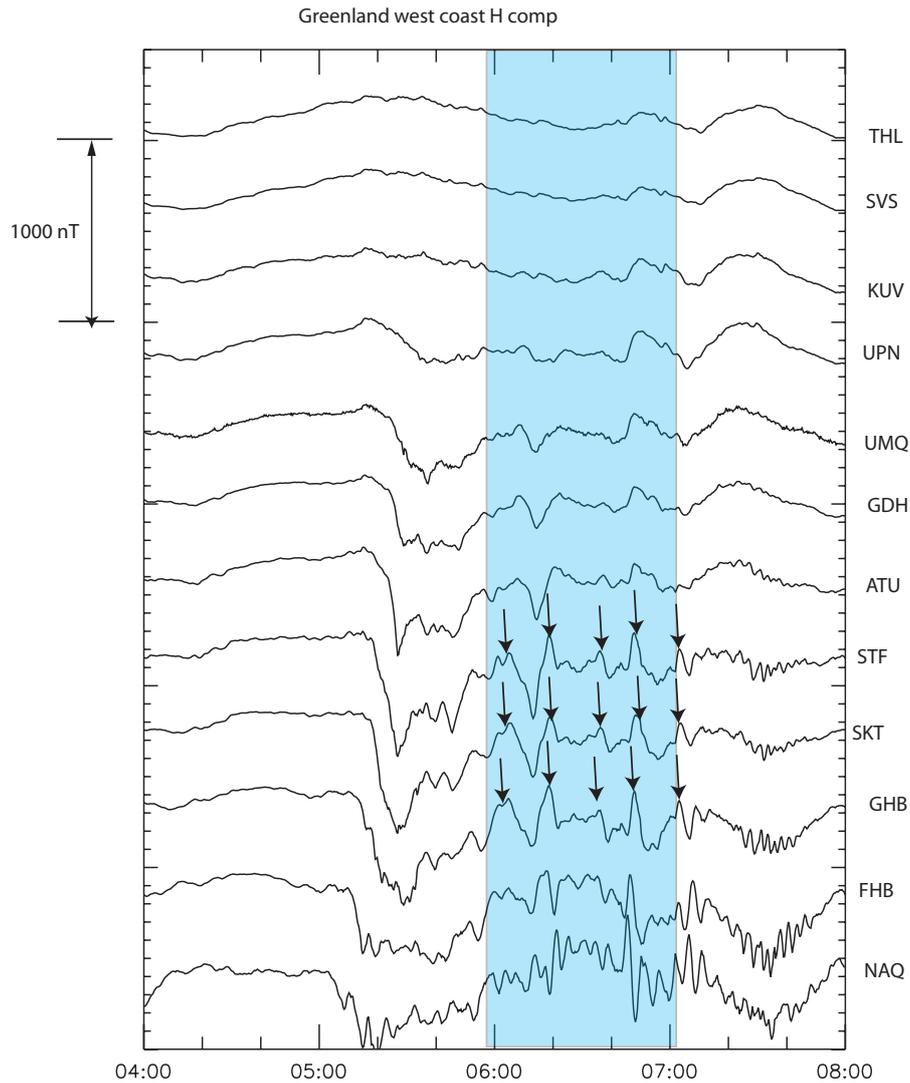


Fig. 8. The stacked H component data from the Greenland west coast magnetometers. It is in the order of decreasing latitude from top to bottom.

largest amplitude (peak to peak amplitude of ~ 400 nT) and longest duration. These oscillations occurred right after a substorm. Figure 6 shows the power spectra resulting from the FFT (Fast Fourier Transform) analysis along the north-south Churchill line. All power spectra plots here use the data within the one-hour interval 06:00–07:00 UT. The CGM (corrected geomagnetic) longitude and latitude pair is indicated for each station. The reference power of 10^4 nT² Hz is also shown for easy comparison. Figure 7 is the power spectra for all the stations along the west-east direction with the same format as Fig. 6. 1.1-mHz peaks are obvious in Figs. 6 and 7.

The stacked H component data from the Greenland west coast magnetometers is shown in Fig. 8 in the order of decreasing latitude. Figure 9a shows the power spectra for the northern stations while Fig. 9b shows those of the southern

stations in Greenland's west coast, in the order of increasing latitude. The power spectra for the Greenland east coast stations are shown in Fig. 10. Besides the 1.1-mHz spectra peak, there are also other spectra peaks: such as the 1.6 mHz and ~ 3.0 mHz ones in the Greenland west coast stations and the ~ 2.1 mHz in Greenland's east coast. The latitude independent 1.1-mHz spectra peak at almost all of the CANOPUS and Greenland stations is apparent, indicating a driven field line resonance at this discrete frequency. The 1.1-mHz pulsations at Kangerlussuaq (STF), Maniitsoq (SKT) and Nuuk (GHB) in Greenland's west coast are less coherent, compared to those at CANOPUS stations. Notice that the pulsations took place after a substorm at CANOPUS locations while the pulsations were still inside the substorm at Greenland. The substorm may have had some effects on the coherence of the waves.

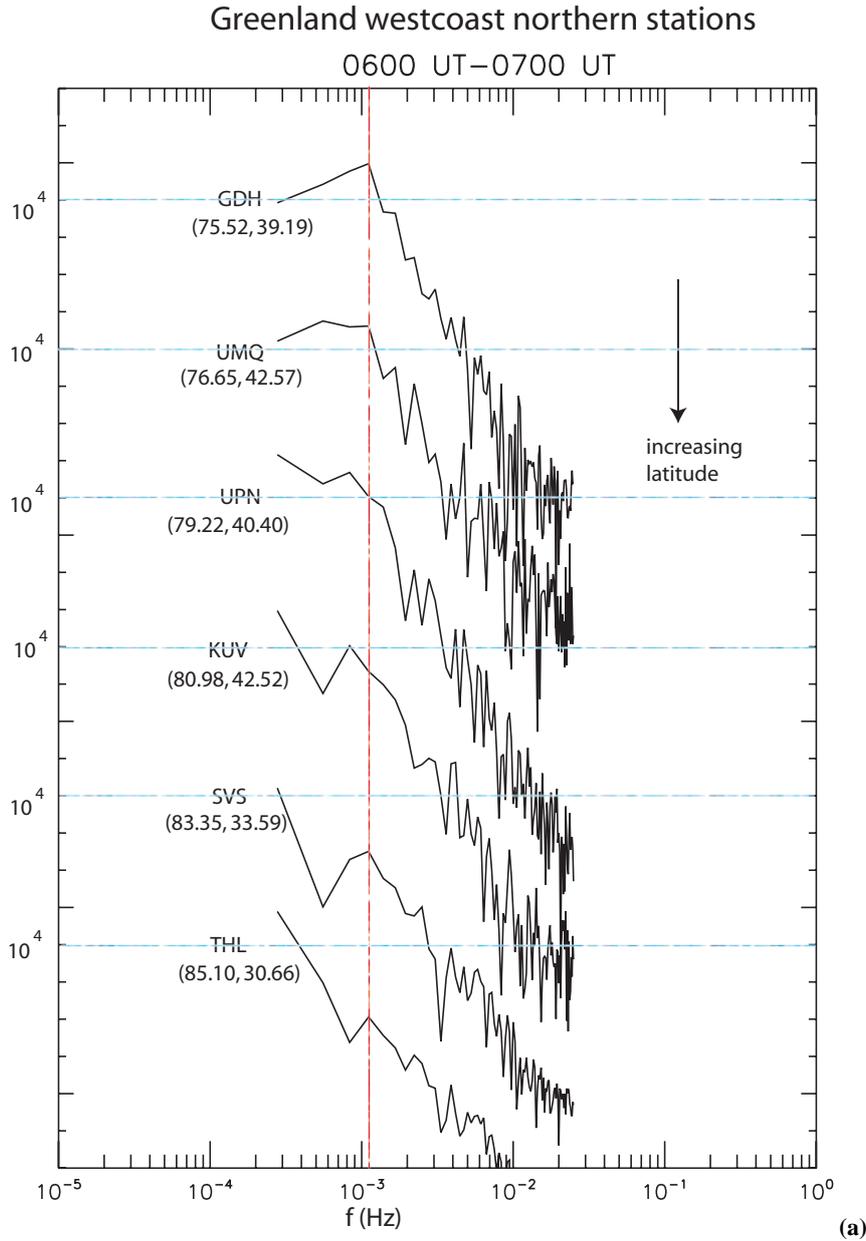


Fig. 9. (a) The power spectra for the northern stations and (b) the power spectra for the southern stations in Greenland's west coast in the order of increasing latitude;

The H component data from both the CANOPUS and Greenland chains shows the clear poleward phase propagation which is expected for a driven FLR (e.g. Wright and Allan, 1996). The complex demodulation analysis (e.g. Mathie and Mann, 2000) of the 1.1-mHz wave from the Churchill line, shown in Fig. 11, further confirms that. Figure 11a shows the amplitude versus CGM latitude for both the H and D component and Fig. 11b shows the phase versus CGM latitude for the two components. The complex demodulation is a time series analysis technique which allows an estimation of the temporal variation of the amplitude and phase

of particular frequency components of a time series. The FLR characteristic is verified by the amplitude maximum between GILL and FCHU and 180° latitudinal (CGM latitude is used here) phase change in the H-component. The dominance of H component on the ground is expected for a dominantly toroidal wave in space observed by Cluster due to ionospheric rotational effects.

Figure 12 shows the 1.1-mHz ground observations in longitudinal order, from stations in the premidnight sector (CANOPUS: top panel) to those in the post-midnight/early morning sector (Greenland: bottom panel). The 1.1-mHz

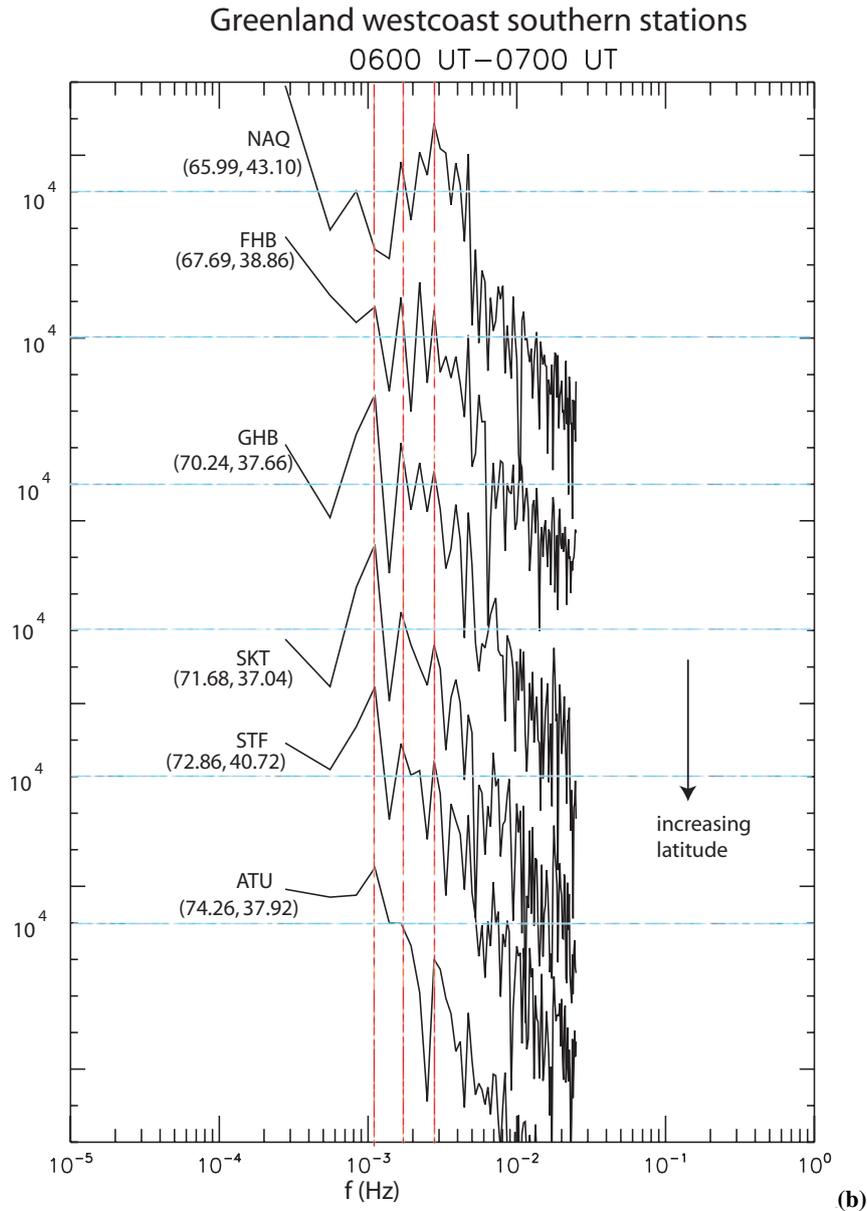


Fig. 9. Continued.

wave spanned over a very large longitudinal range and had a westward propagation as the easternmost station Daneborg (DNB) saw it first. The excited FLR seemed to get damped while going westward as stations in the east (post-midnight) saw more waveforms than those in the west (pre-midnight). By comparing the instantaneous values of the D component phase between the longitudinally spaced CANOPUS station pair (FSMI-GILL) of 1.1-mHz FLR, an estimation of the azimuthal wave number m (e.g. Mann et al., 2002) was obtained, with $m = -1.2 \pm 0.1$ (the H-component gives $m = -1.6 \pm 0.1$). The D component is chosen as the more reliable indicator of the azimuthal propagation as it is not as susceptible as the H component to large phase changes in the

vicinity of a resonance. The negative value of m is consistent with the westward propagation in Fig. 12 and the low m value is expected for the efficient coupling of compressional energy to shear Alfvén FLRs. The negative m value is also consistent with the tailward and westward propagation of the wave front by utilizing the four Cluster spacecraft. Using the relationship $V_{ph} = 2\pi R_E \cos \lambda / m$, we found ionospheric azimuthal phase speed ($V_{ph} = 12.8$ km/s) for the 1.1-mHz FLR, which is in the range of $9 \leq V_{ph} \leq 17$ km/s reported by Ziesolleck and McDiarmid (1994).

The summary of 1.1-mHz wave observations is shown in Fig. 13 in terms of the MLT (magnetic local time)–Mlat (magnetic latitude) map. Strong magnetic pulsations with

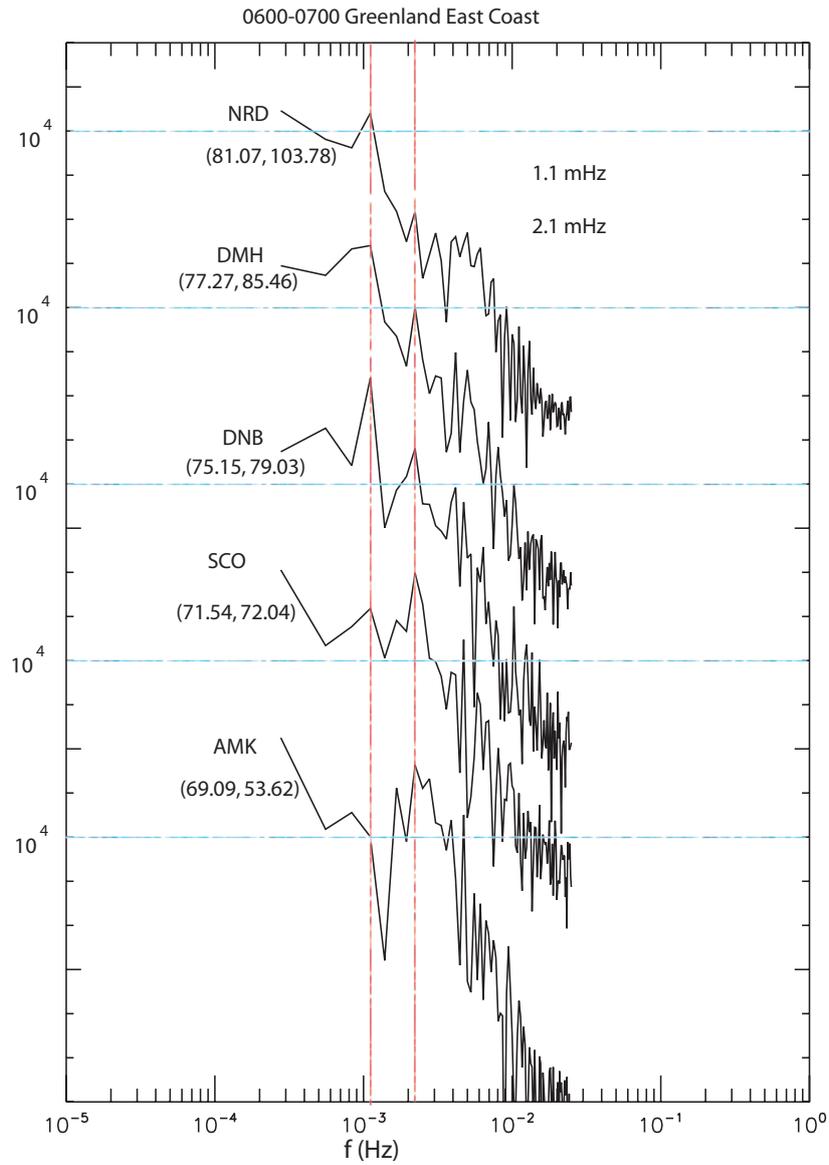


Fig. 10. The power spectra for stations in Greenland's east coast.

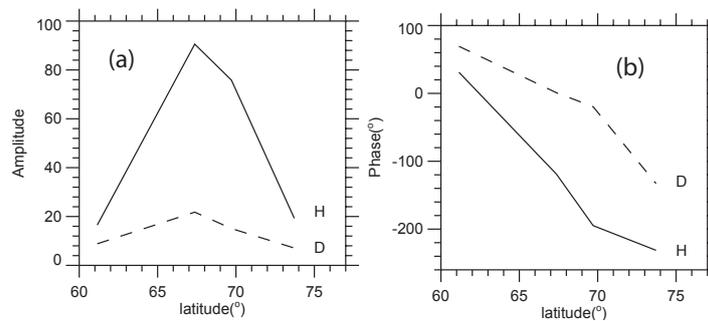


Fig. 11. The complex demodulation analysis of the 1.1-mHz wave from the CANOPUS Churchill line. The FLR characteristic is verified by the amplitude maximum between GILL and FCHU and the 180° latitudinal (CGM latitude is used here) phase change in the H-component.

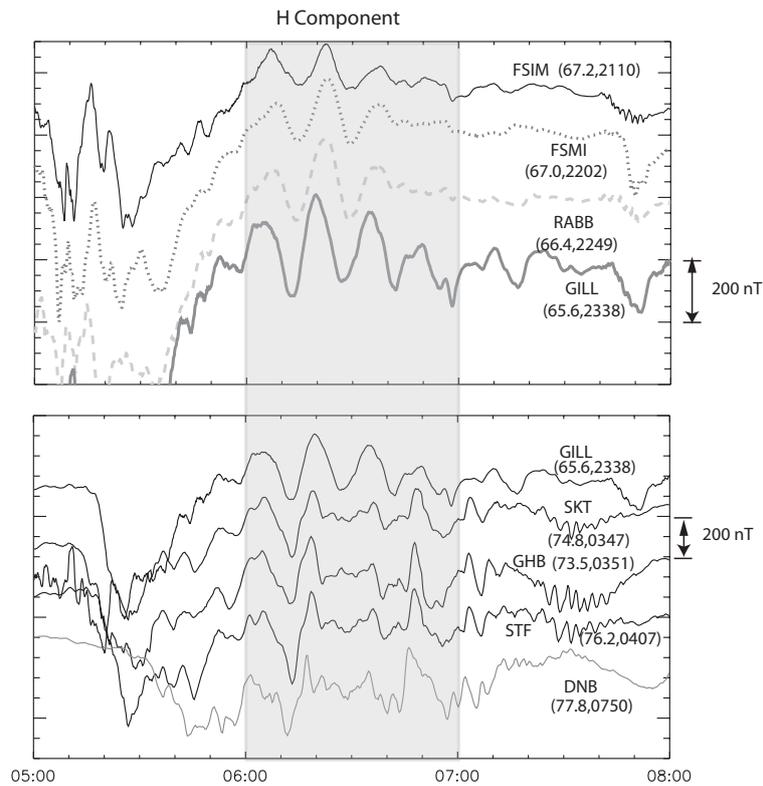


Fig. 12. Part of the ground-based 1.1-mHz pulsations in longitudinal order. Notice the westward propagation of the 1.1-mHz waves.

spectral amplitude greater than $10^6 \text{ nT}^2/\text{Hz}$ were found at GILL and RABB of the CANOPUS chain, SKT, GHB and STF in Greenland’s west coast, and DNB in Greenland’s east coast, which are indicated by a black “+”. Stations with less wave power are indicated by a green “+”. The map shows that the occurrence of the 1.1-mHz pulsation moved to higher latitude with location further into the post-midnight local time. The Tsyganenko magnetic field model T96 (Tsyganenko, 1995) with solar wind dynamics pressure of 4.3 nPa, Dst of -53 nT , IMF B_y of -2.1 nT and B_z of 3.8 nT puts Cluster’s footpoint (70.3° , 01:53, $21.2 R_E$ – this triplet represents MLAT-MLT- equatorial crossing distance l_0) in between GILL (66.4° , 23:38, $21.9 R_E$) and GHB (70.2° , 03:51, $16.8 R_E$). The parameters for other stations are: RABB (67.1° , 22:49, $18.9 R_E$), SKT (71.7° , 03:47, $20.1 R_E$), STF (72.9° , 04:07, $22.2 R_E$), and DNB (75.2° , 07:00, $18.5 R_E$). We can see that the large amplitude 1.1-mHz wave occurred on similar field lines with the average l_0 about $20 R_E$.

No one would likely doubt that Cluster could be located on a field line with an equatorial crossing distance l_0 of $21.9 R_E$ as its location ($\sim 19 R_E$ in radial distance and the off-equator position) in the mid-tail suggested so. However, it is hard to believe intuitively that Gillam could be on a similar stretched field line. There are always controversies over magnetic field line mapping using empirical models such as Tsyganenko models. However, the observational evidence

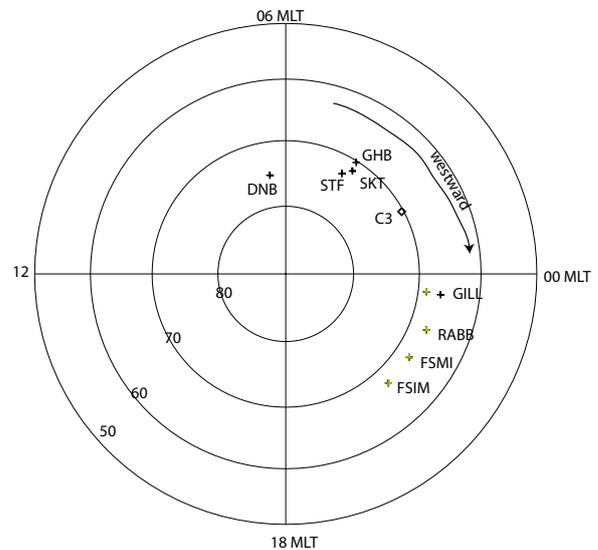


Fig. 13. The summary of 1.1-mHz wave observations in terms of the MLT-Mlat map.

of a driven field line resonance at the discrete frequency of 1.1 mHz with the resonance near Gillam from CANOPUS data, together with the simultaneous observation of the same frequency standing toroidal Alfvén wave by Cluster in space

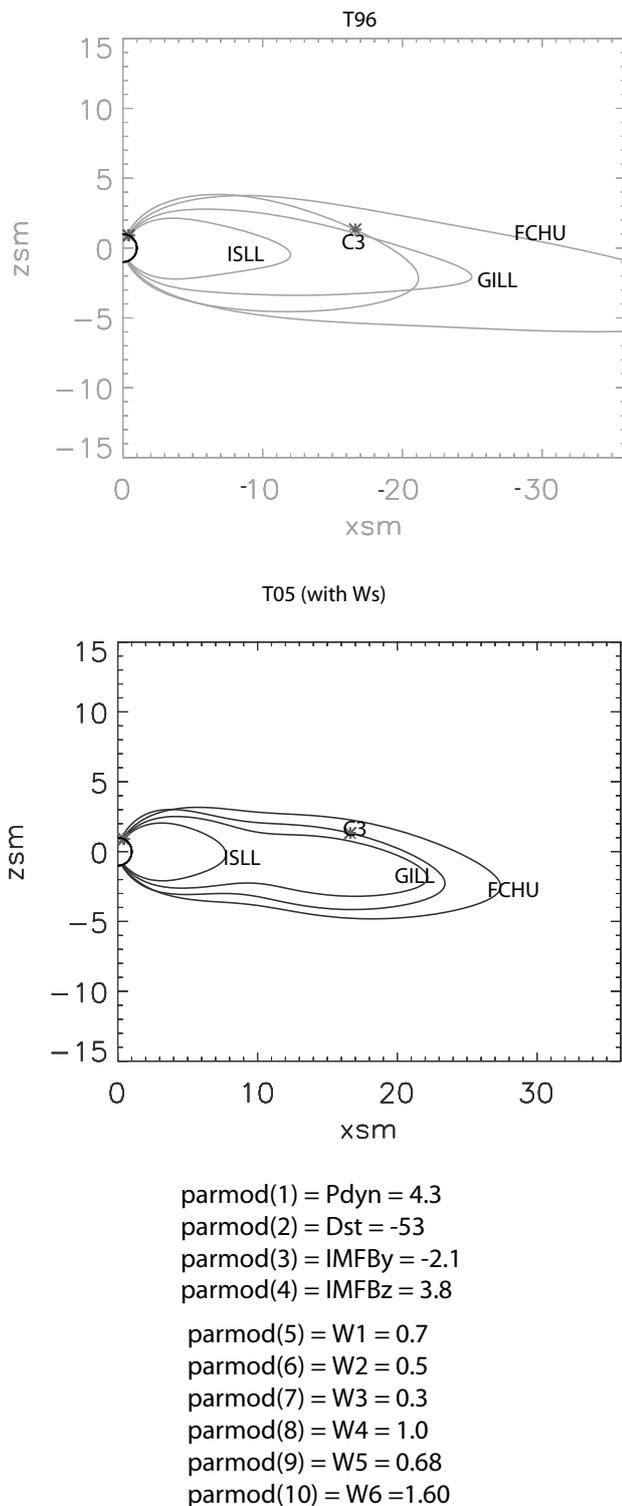


Fig. 14. Modeled magnetic field lines in the x - y plane in SM coordinates with (a) from the T96 model and (b) from the T05 model. The magnetotail is rather stretched from both models.

argue strongly that Cluster and Gillam could be on similar field lines. Here we use Tsyganenko models as a reference in interpreting the data. With the same set of solar wind, IMF and Dst parameters as mentioned above, Tsyganenko 2005 model – T05 (Tsyganenko and Sitnov, 2005) also gives the equatorial crossing distances similar to those of T96 even though the l_0 values are relatively smaller (see bottom panel of Fig. 14). l_0 at Gillam is still around $20 R_E$ using T05. The history of solar wind and IMF conditions and the different response time of magnetospheric current systems are taken into account in T05. The corresponding W parameters ($W1=0.7$, $W2=0.5$, $W3=0.3$, $W4=1.0$, $W5=0.68$, $W6=1.6$) based on ACE data for the event are also shown in Fig. 14. Figure 14 shows the field line tracing results projected in the X - Z plane (in SM coordinates) with top panel from the T96 model and bottom panel from the T05 model. Remember that Gillam and Cluster 3 were in different meridional plane (different local time) and their field lines using T96 don't cross in the 3-D picture. Even though the average IMF B_z was positive, the large solar wind dynamic pressure (4.3 nPa) and a Dst value of -53 nT (rendering relatively strong ring current) are both likely to create a stretched magnetotail.

3.3 Pulsations at other higher frequencies

Closer to Earth compared to Cluster, Polar (66.1° – 66.8° , $01:46$, 11 – $14 R_E$), Geotail (65.9° , $02:00$, $9.8 R_E$) and GOES 12 (63.3° , $00:55$, $8.7 R_E$) also observed quasi-periodic oscillations during the event, noticeable in db_e for Polar, in V_x and V_y (GSM) for Geotail and in B_y (GSM) for GOES 12 (see Fig. 15). Therefore azimuthal oscillations are an obvious feature in all three satellites. The fact that oscillations were found in V_y but not in B_y by Geotail indicates the observed wave had a node of \mathbf{B} at the equator as Geotail was very close to the equatorial region. The observed different discrete frequencies are 1.4 mHz (Polar: after $06:25 \text{ UT}$), 1.9 mHz (Polar: prior to $06:25 \text{ UT}$) and 2.5 mHz (Geotail and GOES 12). However, there were no clear oscillations seen by GOES10, which was located in the pre-midnight.

The quasi-periodic oscillations observed at the multiple spacecraft in space and on the ground are believed to be manifestations of the field-line resonances, as the peak frequency is found to have L-shell dependence that is in agreement with the FLR theory (Fig. 16).

4 Discussion

The apparent westward propagation displayed in Fig. 12 and the negative m value in the near midnight/early morning sector indicate that the 1.1-mHz FLR was possibly driven by a wave source at the discrete frequency (1.1 mHz) resulted from the flank MHD cavity/waveguide (Samson et al., 1992a) mode, propagating tailward. Research results have shown that the compressional waveguide/cavity modes

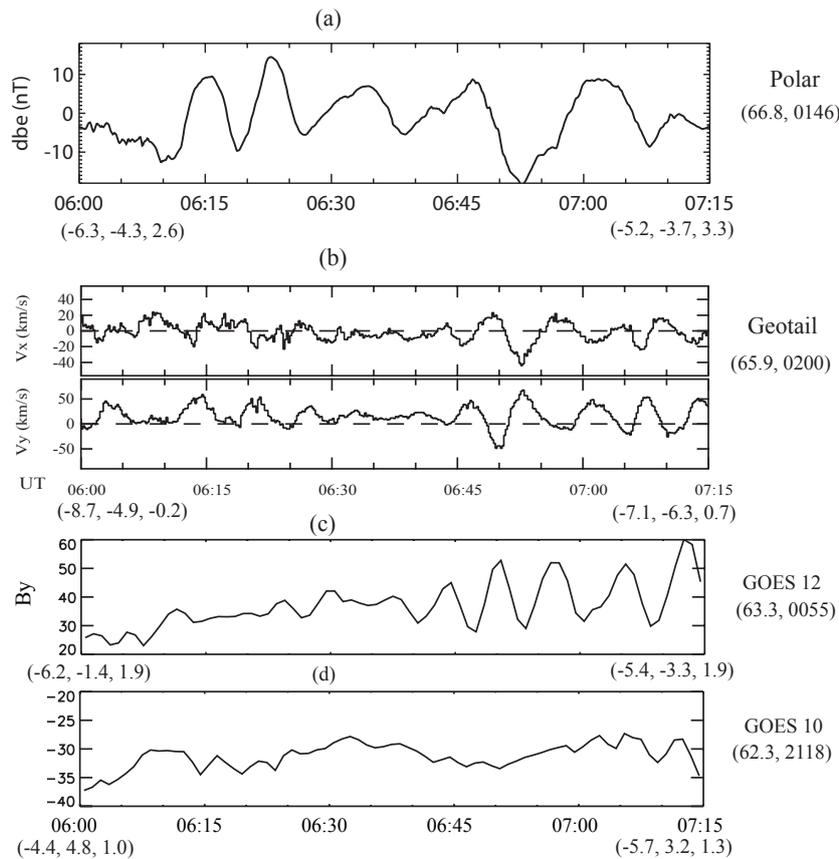


Fig. 15. Oscillations at different frequencies observed by Polar in (a), Geotail in (b), GOES 12 in (c), and GOES 10 in (d).

can be excited by a variety of sources including solar wind and IMF variability, transient dayside reconnection, Kelvin Helmholtz instabilities (KHIs) in the low-latitude boundary layer (LLBL) and substorm associated instabilities. In this case, substorm related instabilities are not a possible source as ground magnetometer data of CANOPUS and Greenland showed that the substorm had an eastward propagation, opposite to the wave propagation. KHI was a possible candidate for energizing the cavity/waveguide modes which drive the 1.1 mHz large amplitude FLRs, as the negative m value indicates a wave source at the dawn flank and the high-speed solar wind stream is a favorable condition for exciting KHIs. However, we cannot exclude other possibilities.

As for the higher frequency (>1.1 mHz) pulsations observed by Polar, Geotail and GOES 12, the source could be a wide band compressional wave coupling to standing waves almost everywhere as suggested by Hasegawa et al. (1983), as the wave packets at higher frequencies look different from that of 1.1 mHz. The change of wave properties at $\sim 06:15$ UT and $\sim 06:45$ UT at NAQ in Fig. 8 is coincident with the North-South-North turning of the IMF B_z at the corresponding times. It seems that IMF has more direct control over the wave properties at higher frequencies.

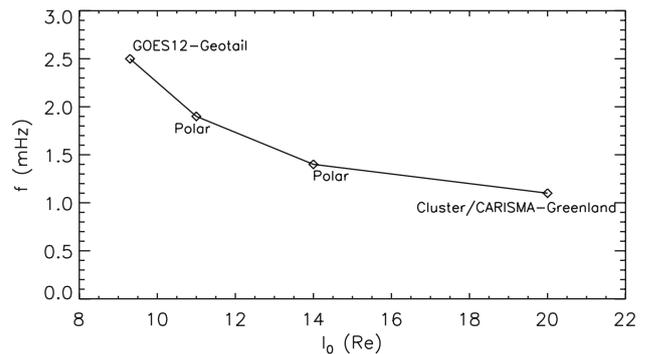


Fig. 16. The frequency (f) versus l shell relationship based on the observations.

The unusually low frequency waves observed at Gillam latitude, such as the 1.3-mHz FLR reported by Lotko et al. (1998), have caused some concern as the frequency cannot be explained using dipole field lines. This study not only shows that an even lower frequency FLR (1.1 mHz) at Gillam is possible, but also provides an explanation for the unusually low frequency oscillations: Gillam can be located on very stretched field lines. It has been demonstrated

theoretically that a stretched magnetic topology can explain anomalously low-frequency FLRs (Lui and Cheng, 2001; Wanliss et al., 2002). Wanliss et al. (2002) have shown theoretically that it requires $n_H=2\text{ cm}^{-3}$ with inclusion of oxygen (see the paper for details) in order to get $f=1.4\text{ mHz}$ at $l_0=13 R_E$. From the CODIF (Composition and Distribution Function) measurements of Cluster, the average density for all the measured species is: $n_H=0.34\text{ cm}^{-3}$, $n_O=0.031\text{ cm}^{-3}$, $n_{He^+}=0.003\text{ cm}^{-3}$, $n_{He^{++}}=0.036\text{ cm}^{-3}$. We can see that there was a substantial presence of O^+ and He^{++} during the event, probably due to a Dst value of -53 nT . This gives a mass density of $\sim 1\text{ amucm}^{-3}$. Using the same mass density of 1 amucm^{-3} (constant along magnetic field) and a self consistently calculated equilibrium magnetosphere, Lui and Cheng (2001) provided the FLR frequency for odd mode shear Alfvén waves at $l_0=19.9$, with the value being 2.46 mHz . Even though this value is higher than the 1.1 mHz waves observed by Cluster, the non-uniform distribution of the plasma mass density along the field line and longer field line length ($l_0=21.2$) at Cluster may help explain the lower frequency (1.1 mHz) standing Alfvén wave at Cluster location. From comparisons with both theoretical results, we can see that stretched field line configuration is required to explain the “anomalously” low frequency oscillations at Cluster and Gillam. Mid-tail extremely low frequency oscillations (sometimes with a frequency even smaller than 1.1 mHz) have also been observed most noticeably in V_y (GSM) polarization on other occasions (e.g. the $\sim 12\text{-min}$ oscillation during 03:30–05:30 UT on 27 July 2001 and during 01:00–03:00 UT on 13 August 2003; the $\sim 21\text{-min}$ oscillation taking place during 10:40–12:40 UT on 11 September 2002) when Cluster was in the tail plasma sheet region, suggesting Pc5/even lower frequency pulsations may be common on stretched field lines in the tail.

This study has important implications in that there may be other means beside the re-reconnection in the tail for the solar wind energy to be transferred to the magnetospheric system. KHI can produce a series of wave modes at all latitudes but the perturbations in the LLBL and subsequently in the plasma sheet are the most important ones as LLBLs are most unstable to KHI and the perturbations produced carry most of the energy and momentum. FLRs through coupling of the cavity/waveguide mode (which can be excited by KHI) to local shear Alfvén waves in the mid-tail plasma sheet region can be an important means for the solar wind energy to be transferred to the tail and subsequently affect the state of the magnetosphere and ionosphere system.

5 Summary

This study constitutes the first coordinated observation of a driven discrete FLR in the mid-magnetotail at a very low frequency (1.1 mHz) by Cluster spacecraft in space and an

array of magnetometer stations on the ground near Cluster’s footpoint, which exhibits classical FLR features. Evidence of FLRs is further provided by simultaneous quasi-periodic oscillations observed by Polar, Geotail and GOES 12 that were located closer to Earth. The observed discrete frequency and L shell relation is in accord with FLR theory. This event shows that the magnetosphere was in oscillation over a large azimuthal extent, but over a narrow latitudinal range.

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