

About Fokker-Planck equation with measurable coefficients and applications to the fast diffusion equation

Nadia Belaribi, Francesco Russo

► **To cite this version:**

Nadia Belaribi, Francesco Russo. About Fokker-Planck equation with measurable coefficients and applications to the fast diffusion equation. 2012. <hal-00645483v2>

HAL Id: hal-00645483

<https://hal.inria.fr/hal-00645483v2>

Submitted on 17 Sep 2012

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Uniqueness for Fokker-Planck equations with measurable coefficients and applications to the fast diffusion equation.

Nadia Belaribi* Francesco Russo†

Abstract

The object of this paper is the uniqueness for a d -dimensional Fokker-Planck type equation with inhomogeneous (possibly degenerated) measurable not necessarily bounded coefficients. We provide an application to the probabilistic representation of the so-called Barenblatt's solution of the fast diffusion equation which is the partial differential equation $\partial_t u = \partial_{xx}^2 u^m$ with $m \in]0, 1[$. Together with the mentioned Fokker-Planck equation, we make use of small time density estimates uniformly with respect to the initial condition.

Keywords: Fokker-Planck; fast diffusion; probabilistic representation; non-linear diffusion; stochastic particle algorithm.

AMS 2010 Subject Classification: MSC 2010: 60H30; 60G44; 60J60; 60H07; 35C99; 35K10; 35K55; 35K65; 65C05; 65C35 .

Submitted to EJP on November 28th, 2011, final version accepted on FIXME!.

1 Introduction

The present paper is divided into three parts.

- i) A uniqueness result on a Fokker-Planck type equation with measurable non-negative (possibly degenerated) multidimensional unbounded coefficients.
- ii) An application to the probabilistic representation of a fast diffusion equation.
- iii) Some small time density estimates uniformly with respect to the initial condition.

In the whole paper $T > 0$ will stand for a fixed final time. In a one dimension space, the Fokker-Planck equation is of the type

$$\begin{cases} \partial_t u(t, x) &= \partial_{xx}^2 (a(t, x)u(t, x)) - \partial_x (b(t, x)u(t, x)), \quad t \in]0, T], \quad x \in \mathbb{R}, \\ u(0, \cdot) &= \mu(dx), \end{cases} \quad (1.1)$$

*Université Paris 13 and ENSTA ParisTech, France. E-mail: belaribi@math.univ-paris13.fr

†ENSTA ParisTech, UMA, 828, Boulevard des Maréchaux, F-91120 Palaiseau, France. E-mail: francesco.russo@ensta-paristech.fr

where $a, b : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ are measurable locally bounded coefficients and μ is a finite real Borel measure. The Fokker-Planck equation for measures is a widely studied subject in the literature whether in finite or infinite dimension. Recent work in the case of time-dependent coefficients with some minimal regularity was done by [9, 16, 30] in the case $d \geq 1$. In infinite dimension some interesting work was produced by [8].

In this paper we concentrate on the case of measurable (possibly) degenerate coefficients. Our interest is devoted to the irregularity of the diffusion coefficient, so we will set $b = 0$. A first result in that direction was produced in [7] where a was bounded, possibly degenerated, and the difference of two solutions was supposed to be in $L^2([\kappa, T] \times \mathbb{R})$, for every $\kappa > 0$ (ASSUMPTION (A)). This result was applied to study the probabilistic representation of a porous media type equation with irregular coefficients. We will later come back to this point. We remark that it is not possible to obtain uniqueness without ASSUMPTION (A). In particular [7, Remark 3.11] provides two measure-valued solutions when a is time-homogeneous, continuous, with $\frac{1}{a}$ integrable in a neighborhood of zero.

One natural question is about what happens when a is not bounded and $x \in \mathbb{R}^d$. A partial answer to this question is given in Theorem 3.1 which is probably the most important result of the paper; it is a generalization of [7, Theorem 3.8] where the inhomogeneous function a was bounded. Theorem 3.1 handles the multidimensional case and it allows a to be unbounded.

An application of Theorem 3.1 concerns the parabolic problem:

$$\begin{cases} \partial_t u(t, x) &= \partial_{xx}^2 (u^m(t, x)), \quad t \in]0, T], \quad x \in \mathbb{R}, \\ u(0, \cdot) &= \delta_0, \end{cases} \quad (1.2)$$

where δ_0 is the Dirac measure at zero and u^m denotes $|u|^{m-1}u$. It is well known that, for $m > 1$, there exists an exact solution to (1.2), the so-called *Barenblatt's density*, see [3]. Its explicit formula is recalled for instance in [34, Chapter 4] and more precisely in [4, Section 6.1]. Equation (1.2) is the *classical porous medium equation*.

In this paper, we focus on (1.2) when $m \in]0, 1[$: the *fast diffusion equation*. In fact, an analogous Barenblatt type solution also exists in this case, see [34, Chapter 4] and references therein; it is given by the expression

$$U(t, x) = t^{-\alpha} \left(D + \tilde{k}|x|^2 t^{-2\alpha} \right)^{-\frac{1}{1-m}}, \quad (1.3)$$

where

$$\alpha = \frac{1}{m+1}, \quad \tilde{k} = \frac{1-m}{2(m+1)m}, \quad D = \left(\frac{I}{\sqrt{\tilde{k}}} \right)^{\frac{2(1-m)}{m+1}}, \quad I = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} [\cos(x)]^{\frac{2m}{1-m}} dx. \quad (1.4)$$

Equation (1.2) is a particular case of the so-called *generalized porous media type equation*

$$\begin{cases} \partial_t u(t, x) &= \partial_{xx}^2 \beta(u(t, x)), \quad t \in]0, T], \\ u(0, x) &= u_0(dx), \quad x \in \mathbb{R}, \end{cases} \quad (1.5)$$

where $\beta : \mathbb{R} \rightarrow \mathbb{R}$ is a monotone non-decreasing function such that $\beta(0) = 0$ and u_0 is a finite measure. When $\beta(u) = u^m$, $m \in]0, 1[$ and $u_0 = \delta_0$, two difficulties arise: first, the coefficient β is of singular type since it is not locally Lipschitz, second, the initial condition is a measure. Another type of singular coefficient is $\beta(u) = H(u - u_c)u$, where H is a Heaviside function and $u_c > 0$ is some critical value, see e.g. [2]. Problem (1.2) with $m \in]0, 1[$ was studied by several authors. For a bounded integrable function as initial condition, the equation in (1.2) is well-stated in the sense of distributions, as a by product of the classical papers [10, 6] on (1.5) with general monotonous coefficient β .

When the initial data is locally integrable, existence was proved by [19]. [11] extended the validity of this result when u_0 is a finite Radon measure in a bounded domain, [29] established existence when u_0 is a locally finite measure in the whole space. The Barenblatt's solution is an *extended continuous solution* as defined in [13, 14]; [14, Theorem 5.2] showed uniqueness in that class. [23, Theorem 3.6] showed existence in a bounded domain of solutions to the fast diffusion equation perturbed by a right-hand side source term, being a general finite and positive Borel measure. As far as we know, there is no uniqueness argument in the literature whenever the initial condition is a finite measure in the general sense of distributions. Among recent contributions, [15] investigated the large time behavior of solutions to (1.2).

The present paper provides the probabilistic representation of the (Barenblatt's) solution of (1.2) and exploits this fact in order to approach it via a Monte Carlo simulation with an L^2 error around 10^{-3} . We make use of the probabilistic procedure developed in [4, Section 4] and we compare it to the exact form of the solution \mathcal{U} of (1.2) which is given by the explicit formulae (1.3)-(1.4). The target of [4] was the case $\beta(u) = H(u - u_c)u$; in that paper those techniques were compared with a deterministic numerical analysis recently developed in [12] which was very performing in that target case. At this stage, the implementation of the same deterministic method for the fast diffusion equation does not give satisfying results; this constitutes a further justification for the probabilistic representation.

We define

$$\Phi(u) = |u|^{\frac{m-1}{2}}, \quad u \in \mathbb{R}, \quad m \in]0, 1[.$$

The probabilistic representation of \mathcal{U} consists in finding a suitable stochastic process Y such that the law of Y_t has $\mathcal{U}(t, \cdot)$ as density. Y will be a (weak) solution of the non-linear SDE

$$\begin{cases} Y_t &= \int_0^t \sqrt{2}\Phi(\mathcal{U}(s, Y_s))dW_s, \\ \mathcal{U}(t, \cdot) &= \text{Law density of } Y_t, \quad \forall t \in]0, T], \end{cases} \quad (1.6)$$

where W is a Brownian motion on some suitable filtered probability space $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \geq 0}, P)$.

To the best of our knowledge, the first author who considered a probabilistic representation of a solution of (1.5) was H. P. Jr. McKean ([26]), particularly in relation with the so-called *propagation of chaos*. In his case β was smooth, but the equation also included a first order coefficient. From then on, literature steadily grew and nowadays there is a vast amount of contributions to the subject, especially when the non-linearity is in the first order part, as e.g. in Burgers' equation. We refer the reader to the excellent survey papers [33] and [18]. A probabilistic interpretation of (1.5) when $\beta(u) = u \cdot |u|^{m-1}$, $m > 1$, was provided for instance in [5]. Recent developments related to chaos propagation when $\beta(u) = u^2$ and $\beta(u) = u^m, m > 1$ were proposed in [28] and [17]. The probabilistic representation in the case of possibly discontinuous β was treated in [7] when β is non-degenerate and in [2] when β is degenerate; the latter case includes the case $\beta(u) = H(u - u_c)u$.

As a preamble to the probabilistic representation we make a simple, yet crucial observation. Let W be a standard Brownian motion.

Proposition 1.1. *Let $\beta : \mathbb{R} \rightarrow \mathbb{R}$ such that $\beta(u) = \Phi^2(u) \cdot u$, $\Phi : \mathbb{R} \rightarrow \mathbb{R}_+$ and u_0 be a probability real measure.*

Let Y be a solution to the problem

$$\begin{cases} Y_t &= Y_0 + \int_0^t \sqrt{2}\Phi(u(s, Y_s))dW_s, \\ u(t, \cdot) &= \text{Law density of } Y_t, \quad \forall t \in]0, T], \\ u(0, \cdot) &= u_0(dx). \end{cases} \quad (1.7)$$

Then $u : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ is solution to (1.5).

Proof of the above result is based on the following lemma.

Lemma 1.2. Let $a : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}_+$ be measurable. Let (Y_t) be a process which solves the SDE

$$Y_t = Y_0 + \int_0^t \sqrt{2a(s, Y_s)} dW_s, \quad t \in [0, T].$$

Consider the function $t \mapsto \rho(t, \cdot)$ from $[0, T]$ to the space of finite real measures $\mathcal{M}(\mathbb{R})$, defined as $\rho(t, \cdot)$ being the law of Y_t . Then ρ is a solution, in the sense of distributions (see (2.2)), of

$$\begin{cases} \partial_t u &= \partial_{xx}^2 (au), \quad t \in]0, T], \\ u(0, \cdot) &= \text{Law of } Y_0. \end{cases} \quad (1.8)$$

Proof of Lemma 1.2. This is a classical result, see for instance [32, Chapter 4]. The proof is based on an application of Itô's formula to $\varphi(Y_t)$, $\varphi \in \mathcal{S}(\mathbb{R})$. \square

Proof of Proposition 1.1. We set $a(s, y) = \Phi^2(u(s, y))$. We apply Lemma 1.2 setting $\rho(t, dy) = u(t, y)dy$, $t \in]0, T]$, and $\rho(0, \cdot) = u_0$. \square

When u_0 is the Dirac measure at zero and $\beta(u) = u^m$, with $m \in]\frac{3}{5}, 1[$, Theorem 5.7 states the converse of Proposition 1.1, providing a process Y which is the unique (weak) solution of (1.6). The first step consists in reducing the proof of that Theorem to the proof of Proposition 5.3 where the Dirac measure, as initial condition of (1.2), is replaced by the function $\mathcal{U}(\kappa, \cdot)$, $0 < \kappa \leq T$. This corresponds to the shifted Barenblatt's solution along a time κ , which will be denoted by $\bar{\mathcal{U}}$. Also, in this case Proposition 5.3 provides an unique strong solution of the corresponding non-linear SDE. That reduction is possible through a weak convergence argument of the solutions given by Proposition 5.3 when $\kappa \rightarrow 0$. The idea of the proof of Proposition 5.3 is the following. Let W be a standard Brownian motion and \bar{Y}_0 be a r.v. distributed as $\mathcal{U}(\kappa, \cdot)$; since $\Phi(\bar{\mathcal{U}})$ is Lipschitz, the SDE

$$\bar{Y}_t = \bar{Y}_0 + \int_0^t \Phi(\bar{\mathcal{U}}(s, \bar{Y}_s)) dW_s, \quad t \in]0, T],$$

admits a unique strong solution. The marginal laws of (\bar{Y}_t) and $\bar{\mathcal{U}}$ can be shown to be both solutions to (1.8) for $a(s, y) = (\bar{\mathcal{U}}(s, y))^{m-1}$; that a will be denoted in the sequel by \bar{a} . The leading argument of the proof is carried by Theorem 3.1 which states uniqueness for measure valued solutions of the Fokker-Planck type PDE (1.8) under some **Hypothesis(B)**. More precisely, to conclude that the marginal laws of (\bar{Y}_t) and $\bar{\mathcal{U}}$ coincide via Theorem 3.1, we show that they both verify the so-called **Hypothesis(B2)**. In order to prove that for $\bar{\mathcal{U}}$, we will make use of Lemma 4.2. The verification of **Hypothesis(B2)** for the marginal laws of \bar{Y} is more involved. It makes use of a small time (uniformly with respect to the initial condition) upper bound for the density of an inhomogeneous diffusion flow with linear growth (unbounded) smooth coefficients, even though the diffusion term is non-degenerate and all the derivatives are bounded. This is the object of Proposition 5.1, the proof of which is based on an application of Malliavin calculus. In our opinion this result alone is of interest as we were not able to find it in the literature. When the paper was practically finished we discovered an interesting recent result of M. Pierre, presented in [20, Chapter 6], obtained independently. This result holds in dimension 1 when the coefficients are locally bounded, non-degenerate and the initial condition has a first moment. In this case, the hypothesis of type (B) is not needed. In particular it allows one to establish Proposition 5.3, but not Theorem 5.7 where the coefficients are not locally bounded on $[0, T] \times \mathbb{R}$.

The paper is organized as follows. Section 2 is devoted to basic notations. Section 3 is concentrated on Theorem 3.1 which concerns uniqueness for the deterministic, time inhomogeneous, Fokker-Planck type equation. Section 4 presents some properties of the Barenblatt's solution \mathcal{U} to (1.2). The probabilistic representation of \mathcal{U} is treated in Section 5. Proposition 5.1 performs small time density estimates for time-inhomogeneous diffusions, the proof of which is located in the Appendix. Finally, Section 6 is devoted to numerical experiments.

2 Preliminaries

We start with some basic analytical framework. In the whole paper d will be a strictly positive integer. If $f : \mathbb{R}^d \rightarrow \mathbb{R}$ is a bounded function we will denote $\|f\|_\infty = \sup_{x \in \mathbb{R}^d} |f(x)|$.

By $S(\mathbb{R}^d)$ we denote the space of rapidly decreasing infinitely differentiable functions $\varphi : \mathbb{R}^d \rightarrow \mathbb{R}$, by $S'(\mathbb{R}^d)$ its dual (the space of tempered distributions). We denote by $\mathcal{M}(\mathbb{R}^d)$ the set of finite Borel measures on \mathbb{R}^d . If $x \in \mathbb{R}^d$, $|x|$ will denote the usual Euclidean norm.

For $\varepsilon > 0$, let K_ε be the Green's function of $\varepsilon - \Delta$, that is the kernel of the operator $(\varepsilon - \Delta)^{-1} : L^2(\mathbb{R}^d) \rightarrow H^2(\mathbb{R}^d) \subset L^2(\mathbb{R}^d)$. In particular, for all $\varphi \in L^2(\mathbb{R}^d)$, we have

$$B_\varepsilon \varphi := (\varepsilon - \Delta)^{-1} \varphi(x) = \int_{\mathbb{R}^d} K_\varepsilon(x - y) \varphi(y) dy. \quad (2.1)$$

For more information about the corresponding analysis, the reader can consult [31]. If $\varphi \in C^2(\mathbb{R}^d) \cap S'(\mathbb{R}^d)$, then $(\varepsilon - \Delta)\varphi$ coincides with the classical associated PDE operator evaluated at φ .

Definition 2.1. We will say that a function $\psi : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ is non-degenerate if there is a constant $c_0 > 0$ such that $\psi \geq c_0$.

Definition 2.2. We will say that a function $\psi : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ has linear growth (with respect to the second variable) if there is a constant C such that $|\psi(\cdot, x)| \leq C(1 + |x|)$, $x \in \mathbb{R}$.

Definition 2.3. Let $a : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}_+$ be a Borel function, $z^0 \in \mathcal{M}(\mathbb{R}^d)$. A (weakly measurable) function $z : [0, T] \rightarrow \mathcal{M}(\mathbb{R}^d)$ is said to be a solution in the sense of distributions of

$$\partial_t z = \Delta(az)$$

with initial condition $z(0, \cdot) = z^0$ if, for every $t \in [0, T]$ and $\phi \in \mathcal{S}(\mathbb{R})$, we have

$$\int_{\mathbb{R}^d} \phi(x) z(t, dx) = \int_{\mathbb{R}^d} \phi(x) z^0(dx) + \int_0^t ds \int_{\mathbb{R}^d} \Delta \phi(x) a(s, x) z(s, dx). \quad (2.2)$$

3 Uniqueness for the Fokker-Planck equation

We now state the main result of the paper which concerns uniqueness for the Fokker-Planck type equation with measurable, time-dependent, (possibly degenerated and unbounded) coefficients. It generalizes [7, Theorem 3.8] where the coefficients were bounded and one-dimensional.

The theorem below holds with two classes of hypotheses: **(B1)**, operating in the multidimensional case, and **(B2)**, more specifically in the one-dimensional case.

Theorem 3.1. Let a be a Borel nonnegative function on $[0, T] \times \mathbb{R}^d$. Let $z_i : [0, T] \rightarrow \mathcal{M}(\mathbb{R}^d)$, $i = 1, 2$, be continuous with respect to the weak topology on finite measures on $\mathcal{M}(\mathbb{R}^d)$. Let z^0 be an element of $\mathcal{M}(\mathbb{R}^d)$. Suppose that both z_1 and z_2 solve the problem $\partial_t z = \Delta(az)$ in the sense of distributions with initial condition $z(0, \cdot) = z^0$.

Then $z := (z_1 - z_2)(t, \cdot)$ is identically zero for every t under the following requirement.

Hypothesis (B). There is $\tilde{z} \in L^1_{loc}([0, T] \times \mathbb{R}^d)$ such that $z(t, \cdot)$ admits $\tilde{z}(t, \cdot)$ as density for almost all $t \in [0, T]$; \tilde{z} will still be denoted by z . Moreover, either **(B1)** or **(B2)** below is fulfilled.

$$\mathbf{(B1)} \quad \text{(i)} \quad \int_{[0, T] \times \mathbb{R}^d} |z(t, x)|^2 dt dx < +\infty, \quad \text{(ii)} \quad \int_{[0, T] \times \mathbb{R}^d} |az|^2(t, x) dt dx < +\infty.$$

(B2) We suppose $d = 1$. For every $t_0 > 0$, we have

$$\text{(i)} \quad \int_{[t_0, T] \times \mathbb{R}} |z(t, x)|^2 dt dx < +\infty, \quad \text{(ii)} \quad \int_{[0, T] \times \mathbb{R}} |az|(t, x) dt dx < +\infty, \quad \text{(iii)} \quad \int_{[t_0, T] \times \mathbb{R}} |az|^2(t, x) dt dx < +\infty.$$

Remark 3.2. The weak continuity of $z(t, \cdot)$ and [7, Remark 3.10] imply that $\sup_{t \in [0, T]} \|z(t, \cdot)\|_{var} < +\infty$, where $\|\cdot\|_{var}$ denotes the total variation. In particular $\sup_{0 < t \leq T} \int_{\mathbb{R}^d} |z(t, x)| dx < +\infty$.

Remark 3.3. 1. If a is bounded then the first item of Hypothesis(B1) implies the second one.

2. If a is non-degenerated, assumption (ii) of Hypothesis(B1) implies assumption (i).

Remark 3.4. Let $d = 1$.

1. If a is non-degenerate, the third assumption of Hypothesis(B2) implies the first one.
2. If $z(t, x) \in L^\infty([t_0, T] \times \mathbb{R})$ then the first item of Hypothesis(B2) is always verified.
3. If a is bounded then assumption (ii) of Hypothesis(B2) is always verified by Remark 3.2; the first item of Hypothesis(B2) implies the third one. So Theorem 3.1 is a strict generalization of [7, Theorem 3.8].
4. Let $(z(t, \cdot), t \in [0, T])$ be the marginal law densities of a stochastic process Y solving

$$Y_t = Y_0 + \int_0^t \sqrt{2a(s, Y_s)} dW_s,$$

with Y_0 distributed as z^0 such that $\int_{\mathbb{R}} |x|^2 z^0(dx) < +\infty$.

If \sqrt{a} has linear growth, it is well known that $\sup_{t \leq T} \mathbb{E}(|Y_t|^2) < +\infty$; so

$$\int_{[0, T] \times \mathbb{R}} |a(s, x)z(s, x)| ds dx = \mathbb{E} \left[\int_0^T a(s, Y_s) ds \right] < +\infty.$$

Therefore assumption (ii) in Hypothesis(B2) is always fulfilled.

Proof of Theorem 3.1. Let z_1, z_2 be two solutions of (2.2); we set $z := z_1 - z_2$. We evaluate, for every $t \in [0, T]$, the quantity

$$g_\varepsilon(t) = \|z(t, \cdot)\|_{-1, \varepsilon}^2,$$

where $\|f\|_{-1, \varepsilon} = \|(\varepsilon - \Delta)^{-\frac{1}{2}} f\|_{L^2}$.

Similarly to the first part of the proof of [7, Theorem 3.8], assuming we can show that

$$\lim_{\varepsilon \rightarrow 0} g_\varepsilon(t) = 0, \quad \forall t \in [0, T], \quad (3.1)$$

we are able to prove that $z(t) \equiv 0$ for all $t \in [0, T]$. We explain this fact.

Let $t \in]0, T]$. We recall the notation $B_\varepsilon f = (\varepsilon - \Delta)^{-1} f$, if $f \in L^2(\mathbb{R}^d)$. Since $z(t, \cdot) \in L^2(\mathbb{R}^d)$ then $B_\varepsilon z(t, \cdot) \in H^2(\mathbb{R}^d)$ and so $\nabla B_\varepsilon z(t, \cdot) \in H^1(\mathbb{R}^d)^d \subset L^2(\mathbb{R}^d)^d$. This gives

$$\begin{aligned} g_\varepsilon(t) &= \int_{\mathbb{R}^d} B_\varepsilon z(t, x) z(t, x) dx = \varepsilon \int_{\mathbb{R}^d} (B_\varepsilon z(t, x))^2 dx - \int_{\mathbb{R}^d} B_\varepsilon z(t, x) \Delta B_\varepsilon z(t, x) dx \\ &= \varepsilon \int_{\mathbb{R}^d} (B_\varepsilon z(t, x))^2 dx + \int_{\mathbb{R}^d} |\nabla B_\varepsilon z(t, x)|^2 dx. \end{aligned}$$

Since the two terms of the above sum are non-negative, if (3.1) holds, then $\sqrt{\varepsilon} B_\varepsilon z(t, \cdot) \rightarrow 0$ (resp. $|\nabla B_\varepsilon z(t, \cdot)| \rightarrow 0$) in $L^2(\mathbb{R}^d)$ (resp. in $L^2(\mathbb{R}^d)^d$). So, for all $t \in]0, T]$, $z(t, \cdot) = \varepsilon B_\varepsilon z(t, \cdot) - \Delta B_\varepsilon z(t, \cdot) \rightarrow 0$, in the sense of distributions, as ε goes to zero. Therefore $z \equiv 0$.

We proceed now with the proof of (3.1). We have the following identities in the sense of distributions:

$$z(t, \cdot) = \int_0^t \Delta(az)(s, \cdot) ds = \int_0^t (\Delta - \varepsilon)(az)(s, \cdot) ds + \varepsilon \int_0^t (az)(s, \cdot) ds, \tag{3.2}$$

which implies

$$B_\varepsilon z(t, \cdot) = - \int_0^t (az)(s, \cdot) ds + \varepsilon \int_0^t B_\varepsilon (az)(s, \cdot) ds. \tag{3.3}$$

Let $\delta > 0$ and (ϕ_δ) a sequence of mollifiers converging to the Dirac delta function at zero. We set $z_\delta(t, x) = \int_{\mathbb{R}^d} z(t, y) \phi_\delta(x - y) dy$, observing that $z_\delta \in (L^1 \cap L^\infty)([0, T] \times \mathbb{R}^d)$. Moreover, (3.2) gives

$$z_\delta(t, \cdot) = \int_0^t \Delta(az)_\delta(s, \cdot) ds.$$

We suppose now Hypothesis(B1) (resp. (B2)). Let $t_0 = 0$ (resp. $t_0 > 0$). By assumption (B1)(ii) (resp. (B2)(iii)), we have $\Delta(az)_\delta \in L^2([t_0, T] \times \mathbb{R}^d)$. Thus, z_δ can be seen as a function belonging to $C([t_0, T]; L^2(\mathbb{R}^d))$. Besides, identities (3.2) and (3.3) lead to

$$z_\delta(t, \cdot) = z_\delta(t_0, \cdot) + \int_{t_0}^t (\Delta - \varepsilon)(az)_\delta(s, \cdot) ds + \varepsilon \int_{t_0}^t (az)_\delta(s, \cdot) ds, \tag{3.4}$$

$$B_\varepsilon z_\delta(t, \cdot) = B_\varepsilon z_\delta(t_0, \cdot) - \int_{t_0}^t (az)_\delta(s, \cdot) ds + \varepsilon \int_{t_0}^t B_\varepsilon (az)_\delta(s, \cdot) ds. \tag{3.5}$$

Proceeding through integration by parts with values in $L^2(\mathbb{R}^d)$, we get

$$\begin{aligned} \|z_\delta(t, \cdot)\|_{-1, \varepsilon}^2 - \|z_\delta(t_0, \cdot)\|_{-1, \varepsilon}^2 &= -2 \int_{t_0}^t ds \langle z_\delta(s, \cdot), (az)_\delta(s, \cdot) \rangle_{L^2} \\ &+ 2\varepsilon \int_{t_0}^t ds \langle (az)_\delta(s, \cdot), B_\varepsilon z_\delta(s, \cdot) \rangle_{L^2}. \end{aligned} \tag{3.6}$$

Then, letting δ go to zero, using assumptions (B1)(i)-(ii) (resp. (B2)(i) and (B2)(iii)) and Cauchy-Schwarz inequality, we obtain

$$\begin{aligned} \|z(t, \cdot)\|_{-1, \varepsilon}^2 - \|z(t_0, \cdot)\|_{-1, \varepsilon}^2 &= -2 \int_{t_0}^t ds \int_{\mathbb{R}^d} a(s, x) |z|^2(s, x) dx \\ &+ 2\varepsilon \int_{t_0}^t ds \langle (az)(s, \cdot), B_\varepsilon z(s, \cdot) \rangle_{L^2}. \end{aligned} \tag{3.7}$$

At this stage of the proof, we assume that **Hypothesis(B1)** is satisfied. Since $t_0 = 0$, we have $z(t_0, \cdot) = 0$. Using the inequality $c_1 c_2 \leq \frac{c_1^2 + c_2^2}{2}$, $c_1, c_2 \in \mathbb{R}$ and Cauchy-Schwarz, (3.7) implies

$$\begin{aligned} \|z(t, \cdot)\|_{-1, \varepsilon}^2 &\leq -2 \int_0^t ds \int_{\mathbb{R}^d} (a|z|^2)(s, x) dx + \varepsilon \int_0^t ds \|az(s, \cdot)\|_{L^2}^2 + \varepsilon \int_0^t ds \|B_\varepsilon z(s, \cdot)\|_{L^2}^2 \\ &\leq \varepsilon \int_0^t ds \|az(s, \cdot)\|_{L^2}^2 + \int_0^t ds \|z(s, \cdot)\|_{-1, \varepsilon}^2, \end{aligned} \quad (3.8)$$

because for $f = z(s, \cdot)$, we have

$$\varepsilon \|B_\varepsilon f\|_{L^2}^2 = \varepsilon \int_{\mathbb{R}^d} \frac{(\mathcal{F}(f))^2(\xi)}{(\varepsilon + |\xi|^2)^2} d\xi \leq \int_{\mathbb{R}^d} \frac{(\mathcal{F}(f))^2(\xi)}{\varepsilon + |\xi|^2} d\xi = \|f\|_{-1, \varepsilon}^2.$$

We observe that the first integral of the right-hand side of (3.8) is finite by assumption (B1)(ii). Gronwall's lemma, applied to (3.8), gives

$$\|z(t, \cdot)\|_{-1, \varepsilon}^2 \leq \varepsilon e^T \int_0^T ds \|az(s, \cdot)\|_{L^2}^2.$$

Letting $\varepsilon \rightarrow 0$, it follows that $\|z(t, \cdot)\|_{-1, \varepsilon}^2 = 0$, $\forall t \in [0, T]$. This concludes the first part of the proof.

We now suppose that **Hypothesis(B2)** is satisfied, in particular $d = 1$. By [7, Lemma 2.2] we have

$$\sup_x 2\varepsilon |B_\varepsilon z(s, x)| \leq \sqrt{\varepsilon} \|z(s, \cdot)\|_{var}.$$

Consequently (3.7) gives

$$\|z(t, \cdot)\|_{-1, \varepsilon}^2 - \|z(t_0, \cdot)\|_{-1, \varepsilon}^2 \leq \sqrt{\varepsilon} \sup_{t \leq T} \|z(t, \cdot)\|_{var} \int_{[t_0, T] \times \mathbb{R}} |az|(s, x) ds dx. \quad (3.9)$$

Besides, arguing like in the proof of [7, Theorem 3.8], we obtain that

$$\lim_{t_0 \rightarrow 0} \|z(t_0, \cdot)\|_{-1, \varepsilon}^2 = 0.$$

We first let $t_0 \rightarrow 0$ in (3.9), which implies

$$\|z(t, \cdot)\|_{-1, \varepsilon}^2 \leq \sqrt{\varepsilon} \sup_{t \leq T} \|z(t, \cdot)\|_{var} \int_{[0, T] \times \mathbb{R}} |az|(s, x) ds; \quad (3.10)$$

we remark that the right-hand side of (3.10) is finite by assumption (B2)(ii). Letting ε go to zero, the proof of (3.1) is finally established. \square

4 Basic facts on the fast diffusion equation

We go on providing some properties of the Barenblatt's solution \mathcal{U} to (1.2) when $m \in]0, 1[$ and given by (1.3)-(1.4).

Proposition 4.1.

(i) \mathcal{U} is a solution in the sense of distributions to (1.2). In particular, for every $\varphi \in C_0^\infty(\mathbb{R})$, we have

$$\int_{\mathbb{R}} \varphi(x) \mathcal{U}(t, x) dx = \varphi(0) + \int_0^t ds \int_{\mathbb{R}} \mathcal{U}^m(s, x) \varphi''(x) dx. \quad (4.1)$$

(ii) $\int_{\mathbb{R}} \mathcal{U}(t, x) dx = 1, \quad \forall t > 0$. In particular, for any $t > 0$, $\mathcal{U}(t, \cdot)$ is a probability density.

(iii) The Dirac measure δ_0 is the initial trace of \mathcal{U} , in the sense that

$$\int_{\mathbb{R}} \gamma(x) \mathcal{U}(t, x) dx \rightarrow \gamma(0), \text{ as } t \rightarrow 0, \tag{4.2}$$

for every $\gamma : \mathbb{R} \rightarrow \mathbb{R}$, continuous and bounded.

Proof of Proposition 4.1. (i) This is a well known fact which can be established by inspection.

(ii) For $M \geq 1$, we consider a sequence of smooth functions (φ^M) , such that

$$\varphi^M(x) \begin{cases} = 0, & \text{if } |x| \geq M + 1; \\ \leq 1, & \text{if } |x| \in [M, M + 1]; \\ = 1, & \text{if } |x| \leq M. \end{cases}$$

By (4.1) we have

$$\int_{\mathbb{R}} \varphi^M(x) \mathcal{U}(t, x) dx = 1 + \int_0^t ds \int_{\mathbb{R}} \mathcal{U}^m(s, x) (\varphi^M)''(x) dx. \tag{4.3}$$

Letting $M \rightarrow +\infty$, by Lebesgue's dominated convergence theorem, the left-hand side of (4.3) converges to $\int_{\mathbb{R}} \mathcal{U}(t, x) dx$. The integral on the right-hand side of (4.3) is bounded by

$$C \int_0^t ds \int_M^{M+1} \mathcal{U}^m(s, x) dx \leq C \int_0^t s^{-\alpha m} \left(D + \tilde{k} M^2 s^{-2\alpha} \right)^{\frac{-m}{1-m}} ds \leq \frac{C}{(\tilde{k} M^2)^{\frac{m}{1-m}}} \int_0^T s^{\frac{m}{1-m}} ds.$$

The last integral on the right is finite as $\frac{m}{1-m} > 0$, for every $m \in]0, 1[$. Therefore the integral in the right-hand side of (4.3) goes to zero as $M \rightarrow +\infty$. This concludes the proof of the second item of Proposition 4.1.

(iii) (4.2) follows by elementary changes of variables. □

Note that the second item of Proposition 4.1 determines the explicit expression of the constant D .

Lemma 4.2.

(i) Suppose that $\frac{1}{3} < m < 1$. Then there is $p \geq 2$ and a constant C_p (depending on T) such that for $0 \leq s < \ell \leq T$

$$\int_{]s, \ell[\times \mathbb{R}} dt dx (\mathcal{U}(t, x))^{\frac{p(m-1)}{2} + 1} \leq C_p (\ell - s). \tag{4.4}$$

(ii) In particular, taking $p = 2$ in (4.4), we get

$$\int_{]0, T[\times \mathbb{R}} dt dx (\mathcal{U}(t, x))^m < +\infty, \tag{4.5}$$

again when m belongs to $]\frac{1}{3}, 1[$.

(iii) If $\frac{1}{5} < m < 1$,

$$\int_{]0, T] \times \mathbb{R}} dt dx (\mathcal{U}(t, x))^{2m} < +\infty. \tag{4.6}$$

(iv) If m belongs to $]\frac{3}{5}, 1[$, then

$$\forall \kappa > 0, \int_{\mathbb{R}} |x|^4 \mathcal{U}(\kappa, x) dx < +\infty. \tag{4.7}$$

Proof of Lemma 4.2.

(i) Using (1.3), we have

$$\int_{]s, \ell] \times \mathbb{R}} (\mathcal{U}(t, x))^{\frac{p(m-1)}{2} + 1} dt dx = \int_{]s, \ell] \times \mathbb{R}} t^{-\frac{\alpha p(m-1)}{2} - \alpha} (D + \tilde{k}|x|^2 t^{-2\alpha})^{\frac{p}{2} - \frac{1}{1-m}} dt dx.$$

Then, setting $y = t^{-\alpha} x \sqrt{\frac{\tilde{k}}{D}}$, we get

$$\begin{aligned} \int_{]s, \ell] \times \mathbb{R}} (\mathcal{U}(t, x))^{\frac{p(m-1)}{2} + 1} dt dx &= \frac{D^{\frac{p+1}{2} - \frac{1}{1-m}}}{\sqrt{\tilde{k}}} \int_s^\ell t^{\frac{p}{2} \alpha(1-m)} dt \int_{\mathbb{R}} (1 + y^2)^{\frac{p}{2} - \frac{1}{1-m}} dy \\ &\leq \frac{D^{\frac{p+1}{2} - \frac{1}{1-m}}}{\sqrt{\tilde{k}}} T^{\frac{p}{2} \alpha(1-m)} (\ell - s) \int_{\mathbb{R}} (1 + y^2)^{\frac{p}{2} - \frac{1}{1-m}} dy. \end{aligned}$$

The last integral is finite if $(p + 1)(1 - m) < 2$. This implies (4.4).

(ii) is a particular case of (i) and (iii) follows by similar arguments as for the proof of (i).

(iv) Now we assume that $m \in]\frac{3}{5}, 1[$. For $\kappa > 0$ we have

$$\int_{\mathbb{R}} |x|^4 \mathcal{U}(\kappa, x) dx = \frac{D^{\frac{3-5m}{2(1-m)}}}{\tilde{k}^{5/2}} \kappa^{4\alpha} \int_{\mathbb{R}} |y|^4 (1 + y^2)^{-\frac{1}{1-m}} dy, \tag{4.8}$$

where this last equality was obtained setting $y = \kappa^{-\alpha} x \sqrt{\frac{\tilde{k}}{D}}$. Clearly, since $m \in]\frac{3}{5}, 1[$, the integral in the right-hand side of (4.8) is finite. Therefore (4.7) is fulfilled. \square

Let $\kappa \in]0, T]$. Given $u : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ we associate

$$\bar{u}(t, x) = u(t + \kappa, x), \quad (t, x) \in [0, T - \kappa] \times \mathbb{R}. \tag{4.9}$$

In particular we have

$$\bar{\mathcal{U}}(t, x) = \mathcal{U}(t + \kappa, x). \tag{4.10}$$

Moreover, for every $x \in \mathbb{R}$, we denote

$$u_{0, \kappa}(x) = \mathcal{U}(\kappa, x). \tag{4.11}$$

Remark 4.3. Function $\bar{\mathcal{U}}$ solves the problem

$$\begin{cases} \partial_t u &= \partial_{xx}^2(u^m), \\ u(0, \cdot) &= u_{0, \kappa}. \end{cases} \tag{4.12}$$

5 The probabilistic representation of the fast diffusion equation

We are now interested in a non-linear stochastic differential equation rendering the probabilistic representation related to (1.2) and given by (1.6). Suppose for a moment that Y_0 is a random variable distributed according to δ_0 , so $Y_0 = 0$ a.s. We recall that, if there exists a process Y being a solution in law of (1.6), then Proposition 1.1 implies that u solves (1.2) in the sense of distributions.

In this subsection we shall prove existence and uniqueness of solutions in law for (1.6). In this respect we first state a tool, given by Proposition 5.1 below, concerning the existence of an upper bound for the marginal law densities of the solution Y of an inhomogeneous SDE with unbounded coefficients. This result has an independent interest.

Proposition 5.1. *Let $\sigma, b : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ be continuous (not necessarily bounded) functions such that $\sigma(t, \cdot), b(t, \cdot)$ are smooth with bounded derivatives of orders greater or equal than one. σ is supposed to be non-degenerate.*

Let $x_0 \in \mathbb{R}$ and $Y_t = (Y_t^{x_0})_{t \in [0, T]}$ be the solution of

$$Y_t = x_0 + \int_0^t \sigma(r, Y_r) dW_r + \int_0^t b(r, Y_r) dr. \tag{5.1}$$

Then, for every $s > 0$, the law of Y_s admits a density denoted $p_s(x_0, \cdot)$.

Moreover, we have

$$p_s(x_0, x) \leq \frac{K}{\sqrt{s}} (1 + |x_0|^4), \quad \forall (s, x) \in]0, T] \times \mathbb{R}, \tag{5.2}$$

where K is a constant which depends on $\|\sigma'\|_\infty, \|b'\|_\infty$ and T but not on x_0 .

- Remark 5.2.**
1. The proof of Proposition 5.1 above is given in Appendix 7.1.
 2. If σ and b is bounded, the classical Aronson’s estimates implies that (5.2) holds even without the $|x_0|^4$ multiplicative term. If σ and b are unbounded, [1] provides an adaptation of Aronson’s estimates; unfortunately they first considered time-homogeneous coefficients, and also their result does not imply (5.2).
 3. If σ and b have polynomial growth and are time-homogeneous, various estimates are given in [25]. However the behavior is of type $\mathcal{O}(t^{-\frac{3}{2}})$ instead of $\mathcal{O}(t^{-\frac{1}{2}})$ when $t \rightarrow 0$.

Let Y_κ be a random variable distributed according to $u_{0,\kappa}$. We are interested in the following result.

Proposition 5.3. *Assume that $m \in]\frac{3}{5}, 1[$. Let B be a classical Brownian motion independent of Y_κ . Then there exists a unique (strong) solution $\bar{Y} = (\bar{Y}_t)_{t \in [0, T-\kappa]}$ of*

$$\begin{cases} \bar{Y}_t &= Y_\kappa + \int_0^t \Phi(\bar{U}(s, \bar{Y}_s)) dB_s, \\ \bar{U}(t, \cdot) &= \text{Law density of } \bar{Y}_t, \quad \forall t \in [0, T - \kappa], \\ \bar{U}(0, \cdot) &= u_{0,\kappa}. \end{cases} \tag{5.3}$$

In particular pathwise uniqueness holds.

Corollary 5.4. *Let W be a classical Brownian motion independent of Y_κ . Therefore there is a unique (strong) solution $Y^\kappa = (Y_t^\kappa)_{t \in [\kappa, T]}$ of*

$$\begin{cases} Y_t^\kappa &= Y_\kappa + \int_\kappa^t \Phi(\mathcal{U}(s, Y_s^\kappa)) dW_s, \\ \mathcal{U}(t, \cdot) &= \text{Law density of } Y_t^\kappa, \quad \forall t \in [\kappa, T], \\ \mathcal{U}(\kappa, \cdot) &= u_{0,\kappa}. \end{cases} \tag{5.4}$$

Proof of Corollary 5.4. We start with the proof of uniqueness. Let $\kappa > 0$. We consider two solutions $Y^{\kappa,1}$ and $Y^{\kappa,2}$ of (5.4), we set $\bar{Y}_t^i = Y_{t+\kappa}^{\kappa,i}$, $\forall t \in [0, T - \kappa]$, $i = 1, 2$ and $B_t = W_{t+\kappa} - W_t$, $\forall t \in [0, T - \kappa]$. Clearly \bar{Y}_t^1 and \bar{Y}_t^2 solve (5.3). Therefore, using Proposition 5.3, we deduce uniqueness for problem (5.4). Existence follows by similar arguments. \square

Proof of Proposition 5.3. Let W be a classical Brownian motion on some filtered probability space. Given the function \bar{U} , defined in (4.10), we construct below a unique process \bar{Y} strong solution of

$$\bar{Y}_t = \bar{Y}_0 + \int_0^t \Phi(\bar{U}(s, \bar{Y}_s)) dW_s. \tag{5.5}$$

From (4.10), for every $(s, y) \in [0, T - \kappa] \times \mathbb{R}$, we have

$$\Phi(\bar{U}(s, y)) = \sqrt{2\bar{a}(s, y)},$$

where

$$\bar{a}(s, y) = (s + \kappa)^{\alpha(1-m)}(D + \tilde{k}|y|^2(s + \kappa)^{-2\alpha}). \tag{5.6}$$

In fact, $\Phi(\bar{U})$ is continuous, smooth with respect to the space parameter and all the space derivatives of order greater or equal than one are bounded; in particular $\Phi(\bar{U})$ is Lipschitz and it has linear growth. Therefore (5.5) admits a strong solution.

By Lemma 1.2 the function $t \mapsto \rho(t, \cdot)$ from $[0, T - \kappa]$ to $\mathcal{M}(\mathbb{R})$, where $\rho(t, \cdot)$ is the law of \bar{Y}_t , is a solution to

$$\begin{cases} \partial_t \rho &= \partial_{xx}^2(\bar{a}\rho), \\ \rho(0, \cdot) &= u_{0,\kappa}. \end{cases} \tag{5.7}$$

To conclude it remains to prove that $\bar{U}(t, y)dy$ is the law of \bar{Y}_t , $\forall t \in [0, T - \kappa]$; in particular the law of the r.v. \bar{Y}_t admits a density. For this we will apply Theorem 3.1 for which we need to check the validity of Hypothesis(B2) when $a = \bar{a}$ and for $z := z_1 - z_2$, where $z_1 := \rho$ and $z_2 := \bar{U}$. By additivity this will be of course fulfilled if we prove it separately for $z := \rho$ and $z := \bar{U}$, which are both solutions to (5.7).

Since \bar{a} is non-degenerate, by Remark 3.4(1), we only need to check items (ii) and (iii) of the mentioned Hypothesis(B2). On one hand, since $\bar{a}(s, y) = \bar{U}^{m-1}(s, y)$, $z := \bar{U}$ verifies Hypothesis(B2) because of items (ii) and (iii) of Lemma 4.2. On the other hand, since $\sqrt{\bar{a}}$ has linear growth, by Remark 3.4.(4) ρ fulfills item (ii) of Hypothesis(B2). Moreover, by Lemma 5.5 below, ρ also verifies item (iii) of Hypothesis(B2). Finally Theorem 3.1 implies that $\bar{U} \equiv \rho$. \square

Lemma 5.5. *Let $\psi : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}_+$, continuous (not necessarily bounded) such that $\psi(t, \cdot)$ is smooth with bounded derivatives of orders greater or equal than one. We also suppose ψ to be non-degenerate.*

We consider a stochastic process $X = (X_t)_{t \in [0, T]}$ strong solution of the SDE

$$X_t = X_0 + \int_0^t \psi(s, X_s) dW_s, \tag{5.8}$$

where X_0 is a random variable distributed according to $u_{0,\kappa}$ defined in (4.11) with $m \in]\frac{3}{5}, 1[$.

For $t \in]0, T]$ the law of X_t has a density $\nu(t, \cdot)$ such that $(\psi^2\nu)(t, x)$ belongs to $L^2([t_0, T] \times \mathbb{R})$, for every $t_0 > 0$.

Proof of Lemma 5.5.

If $X_0 = x_0$, where x_0 is a real number, then Proposition 5.1 implies that, for every $t \in]0, T]$, the law of X_t admits a density $p_t(x_0, \cdot)$. Consequently, if the law of X_0 is $u_{0,\kappa}(x)dx$, for every $t \in]0, T]$, the law of X_t has a density given by

$$\nu(t, x) = \int_{\mathbb{R}} u_{0,\kappa}(x_0)p_t(x_0, x)dx_0.$$

By (5.2) in Proposition 5.1 it follows

$$\sup_{(t,x) \in [t_0, T] \times \mathbb{R}} p_t(x_0, x) \leq K_0(1 + |x_0|^4), \quad \text{where } K_0 = \frac{K}{\sqrt{t_0}}. \tag{5.9}$$

Using (5.9) we get

$$K_1 := \sup_{(t,x) \in [t_0, T] \times \mathbb{R}} |\nu(t, x)| \leq K_0 \int_{\mathbb{R}} (1 + |x_0|^4)\mathcal{U}(\kappa, x_0)dx_0 < \infty; \tag{5.10}$$

the latter inequality is valid because of (4.7) in Lemma 4.2. In the sequel of the proof, the constants K_2, K_3, K_4 will only depend on t_0, T and ψ . Furthermore

$$\int_{[t_0, T] \times \mathbb{R}} ((\psi^2\nu)(t, x))^2 dt dx \leq \sup_{(t,x) \in [t_0, T] \times \mathbb{R}} |\nu(t, x)| \mathbb{E} \left[\int_0^T \psi^4(t, X_t) dt \right].$$

Since ψ has linear growth, this expression is bounded by

$$K_1 K_2 \left(1 + \int_0^T \mathbb{E} \left[\sup_{t \in [0, T]} |X_t|^4 \right] dt \right). \tag{5.11}$$

(5.11) follows because of (5.10). Besides, by Burkholder-Davis-Gundy and Jensen’s inequalities, taking into account the linear growth of ψ , it follows that

$$\mathbb{E} \left[\sup_{t \in [0, T]} |X_t|^4 \right] \leq K_3 \left(\mathbb{E} [|X_0|^4] + \int_0^T \mathbb{E} \left[\sup_{s \in [0, T]} |X_s|^4 \right] ds + T \right).$$

Then, by Gronwall’s lemma, there is another constant K_4 such that

$$\mathbb{E} \left[\sup_{t \in [0, T]} |X_t|^4 \right] \leq K_4 \left(1 + \int_{\mathbb{R}} |x_0|^4 \mathcal{U}(\kappa, x_0) dx_0 \right). \tag{5.12}$$

Finally (5.11), (5.12) and (5.10) allow us to conclude the proof. □

We are now ready to provide the probabilistic representation related to function \mathcal{U} which in fact is only a solution in law of (1.6).

Definition 5.6. We say that (1.6) admits a weak (in law) solution if there is a probability space $(\Omega, \mathcal{F}, \mathbb{P})$, a Brownian motion $(W_t)_{t \geq 0}$ and a process $(Y_t)_{t \geq 0}$ such that the system (1.6) holds. (1.6) admits uniqueness in law if, given $(W^1, Y^1), (W^2, Y^2)$ solving (1.6) on some related probability space, it follows that Y^1 and Y^2 have the same law.

Theorem 5.7. Assume that $m \in]\frac{3}{5}, 1[$. Then there is a unique weak solution (in law) Y of problem (1.6).

Remark 5.8. Indeed the assumption on $m \in]\frac{3}{5}, 1[$ is only required for the application of Theorem 3.1. The arguments following the present proof only use $m > \frac{1}{3}$.

Proof of Theorem 5.7. First we start with the existence of a weak solution for (1.6). Let \mathcal{U} be again the (Barenblatt’s) solution of (1.2). We consider the solution $(Y_t^\kappa)_{t \in [0, T]}$ provided by Corollary 5.4 extended to $[0, \kappa]$, setting $Y_t^\kappa = Y_\kappa$, $t \in [0, \kappa]$. We prove that the laws of processes Y^κ are tight. For this we implement the classical Kolmogorov’s criterion, see [22, Section 2.4, Problem 4.11]. We will show the existence of $p > 2$ such that

$$\mathbb{E} [|Y_t^\kappa - Y_s^\kappa|^p] \leq C_p |t - s|^{\frac{p}{2}}, \quad \forall s, t \in [0, T], \tag{5.13}$$

where C_p will stand for a constant (not always the same), depending on p and T but not on κ . Let $s, t \in]0, T]$. Let $p > 2$. By Burkholder-Davis-Gundy inequality we obtain

$$\mathbb{E} [|Y_t^\kappa - Y_s^\kappa|^p] \leq C_p \mathbb{E} \left[\left| \int_s^t \Phi^2(\mathcal{U}(r, Y_r^\kappa)) dr \right|^{\frac{p}{2}} \right].$$

Then, using Jensen’s inequality and the fact that $\mathcal{U}(r, \cdot)$ is the law density of Y_r^κ , $r \geq \kappa$, we get

$$\mathbb{E} [|Y_t^\kappa - Y_s^\kappa|^p] \leq C_p |t - s|^{\frac{p}{2}-1} \int_s^t dr \int_{\mathbb{R}} \Phi^p(\mathcal{U}(r, y)) \mathcal{U}(r, y) dy. \tag{5.14}$$

We have

$$\int_s^t dr \int_{\mathbb{R}} \Phi^p(\mathcal{U}(r, y)) \mathcal{U}(r, y) dy = \int_s^t dr \int_{\mathbb{R}} dy (\mathcal{U}(r, y))^{\frac{p(m-1)}{2}+1},$$

and, by Lemma 4.2 (i), the result follows.

Consequently there is a subsequence $Y^n := Y^{\kappa_n}$ converging in law (as $C([0, T])$ -valued random elements) to some process Y . Let P^n be the corresponding laws on the canonical space $\Omega = C([0, T])$ equipped with the Borel σ -field. Y will denote the canonical process $Y_t(\omega) = \omega(t)$. Let P be the weak limit of (P^n) .

1) We first observe that the marginal laws of Y under P^n converge to the marginal law of Y under P . Let $t \in]0, T]$. If the sequence (κ_n) is lower than t , then the law of Y_t under P^n equals the constant law $\mathcal{U}(t, x) dx$. Consequently, for every $t \in]0, T]$, the law of Y_t under P is $\mathcal{U}(t, x) dx$.

2) We now prove that Y is a (weak) solution of (1.6), under P . By similar arguments as for the classical stochastic differential equations, see [32, Chapter 6], it is enough to prove that Y (under P) fulfills the martingale problem i.e., for every $f \in C_b^2(\mathbb{R})$, the process

$$\mathbf{(MP)} \quad f(Y_t) - f(0) - \frac{1}{2} \int_0^t f''(Y_s) \Phi^2(\mathcal{U}(s, Y_s)) ds,$$

is an (\mathcal{F}_s) -martingale, where (\mathcal{F}_s) is the canonical filtration associated with Y . $C_b^2(\mathbb{R})$ stands for the set $\{f \in C^2(\mathbb{R}) | f, f', f'' \text{ bounded}\}$. Let \mathbb{E} (resp. \mathbb{E}^n) be the expectation operator with respect to P (resp. P^n). Let $s, t \in [0, T]$ with $s < t$ and $R = R(Y_r, r \leq s)$ be an \mathcal{F}_s -measurable, bounded and continuous (on $C([0, T])$) random variable. In order to show the martingale property **(MP)** of Y , we have to prove that

$$\mathbb{E} \left[\left(f(Y_t) - f(Y_s) - \frac{1}{2} \int_s^t f''(Y_r) \Phi^2(\mathcal{U}(r, Y_r)) dr \right) R \right] = 0, \quad f \in C_b^2(\mathbb{R}). \tag{5.15}$$

We first consider the case when $s > 0$. There is $n \geq n_0$, such that $\kappa_n < s$. Let $f \in C_b^2(\mathbb{R})$; since $(Y_s)_{s \geq \kappa_n}$, under P^n , are still martingales we have

$$\mathbb{E}^n \left[\left(f(Y_t) - f(Y_s) - \frac{1}{2} \int_s^t f''(Y_r) \Phi^2(\mathcal{U}(r, Y_r)) dr \right) R \right] = 0. \tag{5.16}$$

We are able to prove that (5.15) follows from (5.16). Let $\varepsilon > 0$ and $N > 0$ such that

$$\int_s^t dr \int_{\{|y| > \frac{N}{C} - 1\}} \mathcal{U}^m(r, y) dy \leq \varepsilon, \tag{5.17}$$

where C is the linear growth constant of $\Phi^2 \circ \mathcal{U}$ in the sense of Definition 2.2. In order to conclude, passing to the limit in (5.16), we will only have to show that

$$\lim_{n \rightarrow +\infty} \mathbb{E}^n [F(Y)] - \mathbb{E} [F(Y)] = 0, \tag{5.18}$$

where $F(\ell) = \int_s^t dr \Phi^2(\mathcal{U}(r, \ell(r))) f''(\ell(r)) R(\ell(\xi), \xi \leq s)$, $F : C([0, T]) \rightarrow \mathbb{R}$ being continuous but not bounded. The left-hand side of (5.18) equals

$$\begin{aligned} & \mathbb{E}^n [F(Y) - F^N(Y)] + \mathbb{E}^n [F^N(Y)] - \mathbb{E} [F^N(Y)] + \mathbb{E} [F^N(Y) - F(Y)] \\ & := \mathcal{E}_1(n, N) + \mathcal{E}_2(n, N) + \mathcal{E}_3(n, N), \end{aligned} \tag{5.19}$$

where

$$F^N(\ell) = \int_s^t dr (\Phi^2(\mathcal{U}(r, \ell(r)) \wedge N) f''(\ell(r)) R(\ell(\xi), \xi \leq s).$$

Since $\kappa_n < s$, for N large enough, we get

$$\begin{aligned} |\mathcal{E}_1(n, N)| & \leq \|R\|_\infty \|f''\|_\infty \int_s^t dr \int_{\{\Phi^2(\mathcal{U}(r, y)) \geq N\}} (\Phi^2(\mathcal{U}(r, y) - N) \mathcal{U}(r, y) dy \\ & \leq \|R\|_\infty \|f''\|_\infty \int_s^t dr \int_{\{|y| > \frac{N}{C} - 1\}} \mathcal{U}^m(r, y) dy \leq \varepsilon \|R\|_\infty \|f''\|_\infty, \end{aligned} \tag{5.20}$$

taking into account (5.17) and the second item of Lemma 4.2. For fixed N , chosen in (5.17), we have $\lim_{n \rightarrow +\infty} \mathcal{E}_2(n, N) = 0$, since F^N is bounded and continuous. Again, since the law density under P of Y_t , $t \geq s$, is $\mathcal{U}(t, \cdot)$, similarly as for (5.20), we obtain $|\mathcal{E}_3(n, N)| \leq \varepsilon \|R\|_\infty \|f''\|_\infty$. Finally, coming back to (5.19), it follows

$$\limsup_{n \rightarrow +\infty} |\mathbb{E}^n [F(Y)] - \mathbb{E} [F(Y)]| \leq 2\varepsilon \|R\|_\infty \|f''\|_\infty;$$

since $\varepsilon > 0$ is arbitrary, (5.18) is established. So (5.15) is verified for $s > 0$.

3) Now, we consider the case when $s = 0$. We first prove that

$$\mathbb{E} \left[\int_0^T \Phi^2(\mathcal{U}(r, Y_r)) dr \right] < +\infty. \tag{5.21}$$

By item 1) of this proof, the law of Y_r , $r > 0$ admits $\mathcal{U}(r, \cdot)$ as density. Consequently, the left-hand side of (5.21) gives

$$\int_{]0, T]} dr \int_{\mathbb{R}} \Phi^2(\mathcal{U}(r, y)) \mathcal{U}(r, y) dy = \int_0^T dr \int_{\mathbb{R}} \mathcal{U}^m(r, y) dy,$$

which is finite by the second item of Lemma 4.2. Coming back to (5.15), we can now let s go to zero. Since Y is continuous and f is bounded, we clearly have $\lim_{s \rightarrow 0} \mathbb{E}[f(Y_s)R] = \mathbb{E}[f(Y_0)R]$. Moreover

$$\lim_{s \rightarrow 0} \mathbb{E} \left[\left(\int_s^t f''(Y_r) \Phi^2(\mathcal{U}(r, Y_r)) dr \right) R \right] = \mathbb{E} \left[\left(\int_0^t f''(Y_r) \Phi^2(\mathcal{U}(r, Y_r)) dr \right) R \right],$$

using Lebesgue's dominated convergence theorem and (5.21). Consequently we obtain

$$\mathbb{E} \left[\left(f(Y_t) - f(Y_0) - \frac{1}{2} \int_0^t f''(Y_r) \Phi^2(\mathcal{U}(r, Y_r)) dr \right) R \right] = 0. \tag{5.22}$$

It remains to show that $Y_0 = 0$ a.s. This follows because $Y_t \rightarrow Y_0$ a.s., and also in law (to δ_0), by the third item of Proposition 4.1. Finally we have shown that the limiting process Y verifies **(MP)**, which proves the existence of solutions to (1.6).

4) We now prove uniqueness. Since \mathcal{U} is fixed, only uniqueness for the first line of equation (1.6) has to be established. Let $(Y_t^i)_{t \in [0, T]}$, $i = 1, 2$, be two solutions. In order to show that the laws of Y^1 and Y^2 are identical, according to [21, Lemma 2.5], we will verify that their finite marginal distributions are the same. For this we consider $0 = t_0 < t_1 < \dots < t_N = T$. Let $0 < \kappa < t_1$. Obviously we have $Y_{t_0}^i = 0$ a.s., in the corresponding probability space, $\forall i \in \{1, 2\}$. Both restrictions $Y^1|_{[\kappa, T]}$ and $Y^2|_{[\kappa, T]}$ verify (5.4). Since that equation admits pathwise uniqueness, it also admits uniqueness in law by Yamada-Watanabe theorem. Consequently $Y^1|_{[\kappa, T]}$ and $Y^2|_{[\kappa, T]}$ have the same law and in conclusion the law of $(Y_{t_1}^1, \dots, Y_{t_N}^1)$ coincides with the law of $(Y_{t_1}^2, \dots, Y_{t_N}^2)$, thus, the law of $(Y_{t_0}^1, \dots, Y_{t_N}^1)$ coincides with the law of $(Y_{t_0}^2, \dots, Y_{t_N}^2)$. \square

6 Numerical experiments

In order to avoid singularity problems due to the initial condition being a Dirac delta function, we will consider a time translation of \mathcal{U} , denoted v , and defined by

$$v(t, \cdot) = \mathcal{U}(t + 1, \cdot), \quad \forall t \in [0, T].$$

v still solves equation (1.2), for $m \in]0, 1[$, but with now a smooth initial data given by

$$v_0(x) = \mathcal{U}(1, x), \quad \forall x \in \mathbb{R}. \tag{6.1}$$

Indeed, we have the following formula

$$v(t, x) = (t + 1)^{-\alpha} \left(D + \tilde{k}|x|^2(t + 1)^{-2\alpha} \right)^{-\frac{1}{1-m}}, \tag{6.2}$$

where α , \tilde{k} and D are still given by (1.4).

We now wish to compare the exact solution of problem (1.2) to a numerical probabilistic solution. In fact, in order to perform such approximated solutions, we use the algorithm described in Sections 4 of [4] (implemented in Matlab). We focus on the case

$$m = \frac{1}{2}.$$

Simulation experiments: we compute the numerical solution over the time-space grid $[0, 1.5] \times [-15, 15]$. We use $n = 50000$ particles and a time step $\Delta t = 2 \times 10^{-4}$. Figures 1.(a)-(b)-(c)-(d), display the exact and the numerical solutions at times $t = 0$, $t = 0.5$, $t = 1$ and $t = T = 1.5$, respectively. The exact solution for the fast diffusion equation (1.2), given in (6.2), is depicted by solid lines. Besides, Figure 1.(e) describes the time evolution of the discrete L^2 error on the time interval $[0, 1.5]$.

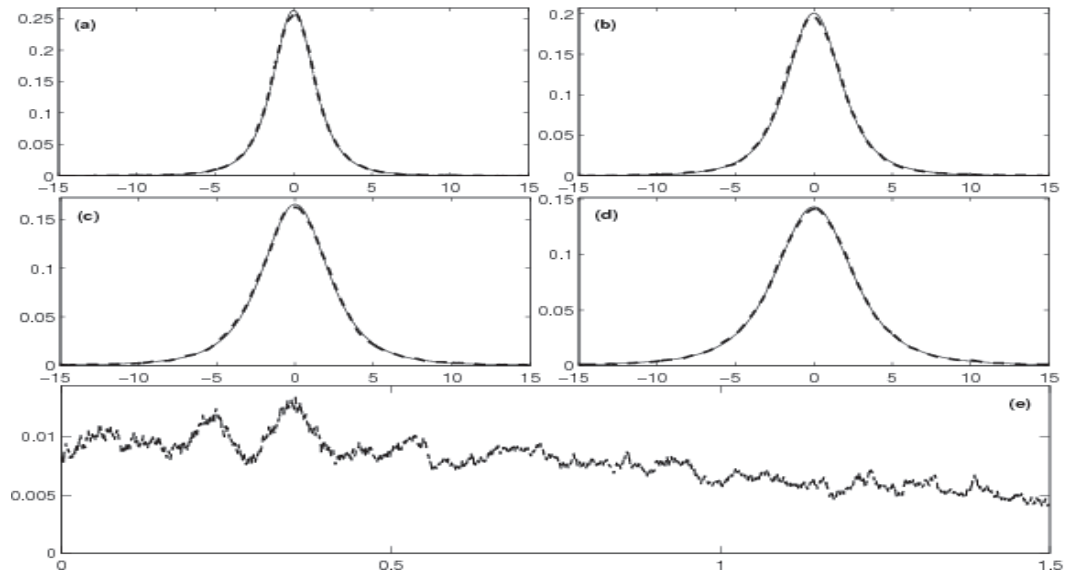


Figure 1: Numerical (dashed line) and exact solutions (solid line) values at $t=0$ (a), $t=0.5$ (b), $t=1$ (c) and $t=1.5$ (d). The evolution of the L^2 error over the time interval $[0, 1.5]$ (e).

7 Appendix

7.1 Proof of Proposition 5.1

We start with some notations for the Malliavin calculus. The set \mathbb{D}^∞ represents the classical Sobolev-Malliavin space of smooth test random variables. $\mathbb{D}^{1,2}$ is defined in the lines after [27, Lemma 1.2.2] and $\mathbb{L}^{1,2}$ is introduced in [27, Definition 1.3.2]. See also [24] for a complete monograph on Malliavin calculus. We state a preliminary result.

Proposition 7.1. *Let N be a non-negative random variable. Suppose, for every $p \geq 1$, the existence of constants $C(p)$ and $\epsilon_0(p)$ such that*

$$P(N \leq \epsilon) \leq C(p) \cdot \epsilon^{p+1}, \quad \forall \epsilon \in]0, \epsilon_0(p)]. \tag{7.1}$$

Then, for every $p \geq 1$,

$$\mathbb{E}(N^{-p}) \leq \epsilon_0(p) \cdot C(p+1) + \epsilon_0(p)^{-p} P(N > \epsilon_0(p)). \tag{7.2}$$

Proof of Proposition 7.1. Let $p \geq 1$ and $\epsilon_0(p) > 0$. Setting $F(x) = P(N \leq x)$, $x \in \mathbb{R}_+$, we have

$$\mathbb{E}(N^{-p}) = I_1 + I_2, \tag{7.3}$$

where

$$I_1 = \int_0^{\epsilon_0(p)} x^{-p} dF(x) \text{ and } I_2 = \int_{\epsilon_0(p)}^{+\infty} x^{-p} dF(x).$$

(7.1) implies that I_1 and I_2 are well-defined. Indeed, on one hand, applying integration by parts on I_1 , we get

$$I_1 = [x^{-p}F(x)]_0^{\epsilon_0(p)} + p \int_0^{\epsilon_0(p)} x^{-p-1}F(x)dx;$$

moreover, there is a constant $C(p)$ such that

$$I_1 \leq (p+1)\epsilon_0(p)C(p). \quad (7.4)$$

On the other hand, again (7.1) says that

$$I_2 \leq \epsilon_0(p)^{-p}(1 - F(\epsilon_0(p))). \quad (7.5)$$

Consequently, using (7.4) and (7.5) and coming back to (7.3), (7.2) is established. \square

Proof of Proposition 5.1. In this proof σ' (resp. b') stands for $\partial_x\sigma$ (resp. $\partial_x b$). Let $Y = (Y_t^{x_0})_{t \in [0, T]}$, be the solution of (5.1). According to [27, Theorem 2.2.2] we have $Y_s \in \mathbb{D}^\infty$, $\forall s \in [0, T]$. Let $s > 0$. Since σ is non-degenerate, by [27, Theorem 2.3.1], the law of Y_s admits a density that we denote by $p_s(x_0, \cdot)$.

The second step consists in a re-scaling, transforming the time s into a noise multiplicative parameter λ ; we set $\lambda = \sqrt{s}$. Indeed, (Y_t) is distributed as (Y_t^λ) , where

$$Y_t^\lambda = x_0 + \lambda \int_0^t \sigma(r\lambda^2, Y_r^\lambda) dW_r + \lambda^2 \int_0^t b(r\lambda^2, Y_r^\lambda) dr.$$

In particular, $Y_s \sim Y_1^\lambda$. By previous arguments, for every $t > 0$, $Y_t^\lambda \in \mathbb{D}^\infty$ and its law admits a density denoted by $p_t^\lambda(x_0, \cdot)$. Our aim consists in showing the existence of a constant K such that

$$p_1^\lambda(x_0, y) \leq \frac{K}{\lambda}(1 + |x_0|^4), \quad \forall y \in \mathbb{R}, \lambda \in]0, \sqrt{T}], \quad (7.6)$$

where K is a constant which does not depend on x_0 and λ . In fact, we will prove that, for every $\lambda \in]0, \sqrt{T}]$,

$$\sup_{y \in \mathbb{R}, t \in]0, 1]} p_t^\lambda(x_0, y) \leq \frac{K}{\lambda}(1 + |x_0|^4). \quad (7.7)$$

We set $Z_t^\lambda = \frac{Y_t^\lambda - x_0}{\lambda}$, $t \in [0, 1]$, so that the density q_t^λ of Z_t^λ fulfills $q_t^\lambda(z) = \lambda p_t^\lambda(x_0, \lambda z + x_0)$, $(t, z) \in [0, 1] \times \mathbb{R}$. In fact, we will have attained (7.7), if we show

$$\sup_{z \in \mathbb{R}, \lambda \in]0, \sqrt{T}]} q_t^\lambda(z) \leq K(1 + |x_0|^4), \quad t \in]0, 1]. \quad (7.8)$$

We express the equation fulfilled by Z ; it yields

$$Z_t^\lambda = \int_0^t \sigma^\lambda(r, Z_r^\lambda) dW_r + \int_0^t b^\lambda(r, Z_r^\lambda) dr, \quad (7.9)$$

where, for every $(r, z) \in [0, 1] \times \mathbb{R}$, we set

$$\sigma^\lambda(r, z) = \sigma(r\lambda^2, \lambda z + x_0), \quad \text{and } b^\lambda(r, z) = \lambda b(r\lambda^2, \lambda z + x_0).$$

At this stage we state the following lemma.

Lemma 7.2. For $\lambda \in]0, 1]$, we shorten by $Z := (Z_t^\lambda)_{t \in [0, T]}$, the solution of (7.9). For every $\gamma \geq 1$, we have

$$\sup_{\lambda \in]0, \sqrt{T}] } \mathbb{E} \left[\sup_{t \in [0, 1]} |Z_t|^\gamma \right] \leq C(1 + |x_0|^\gamma), \tag{7.10}$$

where C is a constant depending on $\|\sigma'\|_\infty, \|b'\|_\infty$ and T , but not on x_0 .

Remark 7.3.

1. For simplicity, in the whole proof of Proposition 5.1, we will set $T = 1$.
2. Since there is no more ambiguity, we will use again the letter s in the considered integrals.

Proof of Lemma 7.2. Let $\lambda \in]0, 1]$ and $\gamma \geq 1$. In the proof C_1 is a constant depending on T , and C_2, C_3 depend on $T, \|\sigma'\|_\infty$ and $\|b'\|_\infty$. Using Burkholder-Davis-Gundy and Jensen's inequalities, we get

$$\mathbb{E} \left[\sup_{\rho \in [0, t]} |Z_\rho|^\gamma \right] \leq C_1 \left(\int_0^t \mathbb{E} [|\sigma(s\lambda^2, \lambda Z_s + x_0)|^\gamma] ds + \int_0^t \mathbb{E} [|\lambda b(s\lambda^2, \lambda Z_s + x_0)|^\gamma] ds \right).$$

Since σ' and b' are bounded, σ and b have linear growth. Therefore, previous expression is bounded by

$$C_2(1 + \lambda^\gamma) \left(1 + |x_0|^\gamma + \lambda^\gamma \int_0^t \mathbb{E} \left[\sup_{\rho \in [0, s]} |Z_\rho|^\gamma \right] ds \right).$$

Since $\lambda \in]0, 1]$ we obtain

$$\mathbb{E} \left[\sup_{\rho \in [0, t]} |Z_\rho|^\gamma \right] \leq C_3 \left(1 + |x_0|^\gamma + \int_0^t \mathbb{E} \left[\sup_{\rho \in [0, s]} |Z_\rho|^\gamma \right] ds \right),$$

Consequently, using Gronwall's lemma, the result follows. □

Now, in order to perform (7.8), we make use of Malliavin calculus for deriving expression (7.9). Omitting λ in the notation Z_t^λ , we get

$$D_r Z_t = \sigma(r\lambda^2, \lambda Z_r + x_0) \mathbb{1}_{[r, 1]}(t) + \lambda \int_r^t \sigma'(s\lambda^2, \lambda Z_s + x_0) D_r Z_s dW_s + \lambda^2 \int_r^t b'(s\lambda^2, \lambda Z_s + x_0) D_r Z_s ds.$$

Consequently

$$D_r Z_t = \sigma(r\lambda^2, \lambda Z_r + x_0) \mathcal{E} \left(\lambda \int_r^t \sigma'(s\lambda^2, \lambda Z_s + x_0) dW_s + \lambda^2 \int_r^t b'(s\lambda^2, \lambda Z_s + x_0) ds \right), \quad r < t,$$

where $\mathcal{E}(S)$ denotes the Doléans exponential of the continuous semi-martingale

$$S_t = \lambda \int_r^t \sigma'(s\lambda^2, \lambda Z_s + x_0) dW_s + \lambda^2 \int_r^t b'(s\lambda^2, \lambda Z_s + x_0) ds, \quad t \in [0, 1].$$

We recall that, $\mathcal{E}_t(S) = \exp(S_t - \frac{1}{2}[S]_t)$. Consequently, for fixed $t \in]0, 1]$, we have

$$\langle DZ_t, DZ_t \rangle = \int_0^t \sigma^2(r\lambda^2, \lambda Z_r + x_0) \mathcal{E}^2(\lambda; r, t) dr, \tag{7.11}$$

where

$$\mathcal{E}(\lambda; r, t) = \exp \left(\lambda \int_r^t \sigma'(s\lambda^2, \lambda Z_s + x_0) dW_s + \lambda^2 \int_r^t \left(b' - \frac{(\sigma')^2}{2} \right) (s\lambda^2, \lambda Z_s + x_0) ds \right). \quad (7.12)$$

We set, for every $s \leq t$, $G(\lambda, s) = \frac{D_s Z_t}{\langle DZ_t, DZ_t \rangle}$. In view of the application of [27, Proposition 2.1.1], which implies the useful expression (7.25) for the density of Z_t , we will need to show that $G(\lambda, \cdot)$ belongs to the domain of the divergence operator δ , denoted by $Dom \delta$. It will be the case if $G(\lambda, \cdot) \in \mathbb{L}^{1,2}(H)$ with $H = L^2([0, T])$. In fact, by the lines after [27, Definition 1.3.2], we know that $\mathbb{L}^{1,2} \subset Dom \delta$. Since $Z_t \in \mathbb{D}^\infty$, we can deduce that $\frac{1}{\langle DZ_t, DZ_t \rangle}$ belongs to \mathbb{D}^∞ , provided that we prove

$$\frac{1}{\langle DZ_t, DZ_t \rangle} \in L^p(\Omega), \quad \forall p \geq 1, \quad (7.13)$$

see [27, Lemma 2.1.6]. Since \mathbb{D}^∞ is an algebra, $G(\lambda, s) \in \mathbb{D}^\infty$, for $s \in]0, T]$, and so $G(\lambda, s) \in \mathbb{D}^{1,2}$. (7.13) will be the object of Proposition 7.5. According to [27, Definition 1.3.2], to affirm that $G(\lambda, \cdot)$ belongs to $\mathbb{L}^{1,2}$ it remains to show the existence of a measurable version of $D_{s_1} G(\lambda, s)$, $(s_1, s) \in [0, t]^2$, such that

$$\mathbb{E} \left[\int_{[0,t]^2} (D_{s_1} G(\lambda, s))^2 ds_1 ds \right] < +\infty. \quad (7.14)$$

We first state the following Lemma.

Lemma 7.4. *For every $q > 1$, there exists a constant $C_0(q)$ such that*

$$\sup_{0 < r \leq t \leq 1, \lambda \in]0, 1]} \mathbb{E} [(\mathcal{E}(\lambda; r, t))^q] \leq C_0(q). \quad (7.15)$$

Proof of Lemma 7.4. Let $\lambda \in]0, 1]$ and $q > 1$. For fixed $0 < r \leq t \leq 1$, (7.12) gives

$$\mathcal{E}^q(\lambda; r, t) = M(\lambda; r, t, q) \exp \left(\frac{\lambda^2(q^2 - q)}{2} \int_r^t (\sigma')^2(\rho\lambda^2, \lambda Z_\rho + x_0) d\rho + q\lambda^2 \int_r^t b'(\rho\lambda^2, \lambda Z_\rho + x_0) d\rho \right),$$

where

$$M(\lambda; r, t, q) = \exp \left(\int_r^t \lambda q \sigma'(\rho\lambda^2, \lambda Z_\rho + x_0) dW_\rho - \frac{1}{2} \int_r^t (q\lambda \sigma')^2(\rho\lambda^2, \lambda Z_\rho + x_0) d\rho \right). \quad (7.16)$$

In fact, since σ' is bounded, the stochastic exponential $M(\lambda; r, t, q)$ verifies Novikov's condition; therefore it is a martingale. So $\mathbb{E}(M(\lambda; r, t, q)) = 1$. In addition, since b' is also bounded and $\lambda \in]0, 1]$, we get $\mathbb{E} [(\mathcal{E}(\lambda; r, t))^q] \leq C_0(q)$, where $C_0(q) = \exp(2(q^2 - q)\|\sigma'\|_\infty^2 + 2q\|b'\|_\infty)$. Consequently (7.15) is established. \square

Proposition 7.5. *There is a constant \mathcal{C} (not depending on x_0) such that*

$$\sup_{(t, \lambda) \in]0, 1]^2} \mathbb{E}[(\langle DZ_t, DZ_t \rangle)^{-p}] \leq \mathcal{C}, \quad \forall p \geq 1.$$

Proof of Proposition 7.5. Let $t \in]0, 1]$ fixed, $\epsilon_0 = \frac{c_0 t}{8}$, where c_0 is a non-degeneracy constant of σ^2 in the sense of Definition 2.1. Consider $\epsilon \in]0, \epsilon_0[$, we set $N := N^\lambda = \langle$

$DZ_t^\lambda, DZ_t^\lambda >$, where we recall that $\langle DZ_t, DZ_t \rangle$ appears in (7.11) and (7.12). According to Proposition 7.1 we have to evaluate $P(N \leq \epsilon)$. Taking into account the lower bound of σ we have

$$\begin{aligned} P(N \leq \epsilon) &\leq P\left(\int_0^t dr \mathcal{E}^2(\lambda; r, t) \leq \frac{\epsilon}{c_0}\right) \\ &\leq P\left(\left(\int_{t-\frac{4\epsilon}{c_0}}^t dr \mathcal{E}^2(\lambda; r, t)\right)^{\frac{1}{2}} \leq \sqrt{\frac{\epsilon}{c_0}}\right) \\ &\leq P\left(\left(\int_{t-\frac{4\epsilon}{c_0}}^t dr\right)^{\frac{1}{2}} - \left(\int_{t-\frac{4\epsilon}{c_0}}^t \mathcal{E}^2(\lambda; r, t) dr\right)^{\frac{1}{2}} \geq \sqrt{\frac{\epsilon}{c_0}}\right). \end{aligned} \tag{7.17}$$

By the inverse triangle inequality of the $L^2([t - \frac{4\epsilon}{c_0}, t])$ -norm we get

$$P(N \leq \epsilon) \leq P\left(\int_{t-\frac{4\epsilon}{c_0}}^t (1 - \mathcal{E}(\lambda; r, t))^2 dr \geq \frac{\epsilon}{c_0}\right).$$

Let $p \geq 1$. By Chebyshev’s inequality this is lower than

$$\left(\frac{c_0}{\epsilon}\right)^{p+1} \mathbb{E}\left[\left(\int_{t-\frac{4\epsilon}{c_0}}^t (1 - \mathcal{E}(\lambda; r, t))^2 dr\right)^{p+1}\right].$$

Then, using Jensen’s inequality, we get

$$P(N \leq \epsilon) \leq 4^p \left(\frac{\epsilon}{c_0}\right)^{-1} \int_{t-\frac{4\epsilon}{c_0}}^t \mathbb{E}\left[(1 - \mathcal{E}(\lambda; r, t))^{2(p+1)}\right] dr. \tag{7.18}$$

Furthermore (7.12) implies that $\mathcal{E}(\lambda; r, t)$ solves

$$\mathcal{E}(\lambda; r, t) = 1 + \lambda \int_r^t \mathcal{E}(\lambda; r, s) \sigma'(s\lambda^2, \lambda Z_s + x_0) dW_s + \lambda^2 \int_r^t \mathcal{E}(\lambda; r, s) b'(s\lambda^2, \lambda Z_s + x_0) ds.$$

Thus

$$\begin{aligned} \mathbb{E}\left[(\mathcal{E}(\lambda; r, t) - 1)^{2(p+1)}\right] &\leq 2^{2(p+1)} \mathbb{E}\left[\left|\lambda \int_r^t \mathcal{E}(\lambda; r, s) \sigma'(s\lambda^2, \lambda Z_s + x_0) dW_s\right|^{2(p+1)}\right] \\ &\quad + 2^{2(p+1)} \mathbb{E}\left[\left|\lambda^2 \int_r^t \mathcal{E}(\lambda; r, s) b'(s\lambda^2, \lambda Z_s + x_0) ds\right|^{2(p+1)}\right]. \end{aligned} \tag{7.19}$$

On one hand, using Jensen’s inequality and $\lambda \in [0, 1]$, we obtain

$$\mathbb{E}\left[\left|\lambda^2 \int_r^t \mathcal{E}(\lambda; r, s) b'(s\lambda^2, \lambda Z_s + x_0) ds\right|^{2(p+1)}\right] \leq \|b'\|_\infty |t - r|^{2p+1} \int_r^t \mathbb{E}\left[(\mathcal{E}(\lambda; r, s))^{2(p+1)}\right] ds. \tag{7.20}$$

On the other hand, by Burkholder-Davis-Gundy inequality, we get

$$\mathbb{E}\left[\left|\lambda \int_r^t \mathcal{E}(\lambda; r, s) \sigma'(s\lambda^2, \lambda Z_s + x_0) dW_s\right|^{2(p+1)}\right] \leq \mathbb{E}\left[\left|\lambda^2 \int_r^t \mathcal{E}^2(\lambda; r, s) (\sigma')^2(s\lambda^2, \lambda Z_s + x_0) ds\right|^{p+1}\right].$$

Applying again Jensen's inequality gives

$$\mathbb{E} \left[\left| \lambda \int_r^t \mathcal{E}(\lambda; r, s) \sigma'(s\lambda^2, \lambda Z_s + x_0) dW_s \right|^{2(p+1)} \right] \leq \|\sigma'\|_\infty |t-r|^p \int_r^t \mathbb{E} [(\mathcal{E}(\lambda; r, s))^{2(p+1)}] ds. \quad (7.21)$$

Therefore (7.20), (7.21) and (7.19) lead to

$$\mathbb{E} [(\mathcal{E}(\lambda; r, t) - 1)^{2(p+1)}] \leq C(T, \|\sigma'\|_\infty, \|b'\|_\infty) |t-r|^p \int_r^t \mathbb{E} [(\mathcal{E}(\lambda; r, s))^{2(p+1)}] ds. \quad (7.22)$$

By Lemma 7.4 and (7.22), there is a constant $C_0(2(p+1))$ such that

$$\mathbb{E} [(\mathcal{E}(\lambda; r, t) - 1)^{2(p+1)}] \leq C_1(T, \|\sigma'\|_\infty, \|b'\|_\infty) C_0(2(p+1)). \quad (7.23)$$

Then, coming back to (7.18) and using (7.23), we obtain

$$\forall \epsilon \in]0, \epsilon_0], \quad P(N \leq \epsilon) \leq C(p) \epsilon^{p+1}, \quad (7.24)$$

where $C(p) = \frac{4^{2(p+1)} C_0(2(p+1)) C_1(T, \|\sigma'\|_\infty, \|b'\|_\infty)}{p+1}$. Finally, using Proposition 7.1, the result follows. \square

We go on with the proof of Proposition 5.1 taking into account the considerations before Lemma 7.4. In fact [27, Proposition 2.1.1] allows us to express, for fixed $t \in]0, 1]$,

$$q_t^\lambda(z) = \mathbb{E} [\mathbf{1}_{\{Z_t > z\}} \delta(G(\lambda, \cdot))]; \quad (7.25)$$

using Cauchy-Schwarz inequality, it implies that

$$q_t^\lambda(z) \leq \sqrt{\mathbb{E} [|\delta(G(\lambda, \cdot))|^2]}. \quad (7.26)$$

According to (1.48) in [27], (7.26) implies

$$q_t^\lambda(z) \leq \left(\mathbb{E} \left[\int_0^t G^2(\lambda, s) ds \right] + \mathbb{E} \left[\int_{[0,t]^2} (D_{s_1} G(\lambda, s))^2 ds_1 ds \right] \right)^{\frac{1}{2}}. \quad (7.27)$$

Now we state a result that estimates the two terms in the right-hand side of (7.27). Indeed, we have the following.

Proposition 7.6. *For every $\lambda \in]0, 1]$, $G(\lambda, \cdot) \in \mathbb{L}^{1,2}$. Moreover, the following statements hold: (i) $\mathbb{E} \left[\int_0^t G^2(\lambda, s) ds \right] \leq C_1 (1 + |x_0|^2)$, (ii) $\mathbb{E} \left[\int_{[0,t]^2} (D_{s_1} G(\lambda, s))^2 ds_1 ds \right] \leq C_2 (1 + |x_0|^8)$, where C_1 and C_2 depend on T , $\|\sigma'\|_\infty$ and $\|b'\|_\infty$, but not on x_0 .*

Proof of Proposition 7.6. (i) First we set $I_1 = \mathbb{E} \left[\int_0^t G^2(\lambda, s) ds \right]$, recalling that

$$G(\lambda, s) = \frac{\sigma(s\lambda^2, \lambda Z_s + x_0) \mathcal{E}(\lambda; s, t)}{G_{den}}, \quad \text{where } G_{den} = \langle DZ_t, DZ_t \rangle.$$

By Cauchy-Schwarz inequality, we have

$$I_1 \leq \left(\mathbb{E} \left[\int_0^t \sigma^4(s\lambda^2, \lambda Z_s + x_0) ds \right] \mathbb{E} \left[\int_0^t \frac{\mathcal{E}^4(\lambda; s, t)}{G_{den}^4} ds \right] \right)^{\frac{1}{2}}. \quad (7.28)$$

Since σ has linear growth, by Lemma 7.2 and using again Cauchy-Schwarz inequality, there is a constant C_1 such that

$$I_1 \leq C_1 (1 + |x_0|^2) \left(\mathbb{E} [G_{den}^{-s}] \int_0^t \mathbb{E} [\mathcal{E}^s(\lambda; s, t)] ds \right)^{\frac{1}{4}}. \tag{7.29}$$

Consequently, using Proposition 7.5 and Lemma 7.4, the first item of Proposition 7.6 is established.

(ii) We set $I_2 = \mathbb{E} \left[\int_{[0,t]^2} (D_{s_1} G(\lambda, s))^2 ds_1 ds \right]$.

On one hand, by usual Malliavin differentiation rules, we obtain

$$\begin{aligned} D_{s_1} G(\lambda, s) = & \lambda \sigma'(s\lambda^2, \lambda Z_s + x_0) \sigma(s_1\lambda^2, \lambda Z_{s_1} + x_0) \mathcal{E}(\lambda; s_1, s) \frac{\mathcal{E}(\lambda; s, t)}{G_{den}} \mathbb{1}_{[0,s]}(s_1) \\ & + \sigma(s\lambda^2, \lambda Z_s + x_0) \frac{D_{s_1} \mathcal{E}(\lambda; s, t)}{G_{den}} - \sigma(s\lambda^2, \lambda Z_s + x_0) \mathcal{E}(\lambda; s, t) \frac{D_{s_1} G_{den}}{G_{den}^2}. \end{aligned} \tag{7.30}$$

The right-hand side being measurable with respect to $\Omega \times [0, T]^2$, $G(\lambda, \cdot)$ will belong to $\mathbb{L}^{1,2}$ if (ii) is established. At this point we need to evaluate $D_{s_1} \mathcal{E}(\lambda; s, t)$. From now on, for the sake of simplicity, we will only expose the calculations in the case when $b \equiv 0$. In fact, we have

$$\begin{aligned} D_{s_1} \mathcal{E}(\lambda; s, t) = & \mathcal{E}(\lambda; s, t) D_{s_1} \left(\lambda \int_s^t \sigma'(\ell\lambda^2, \lambda Z_\ell + x_0) dW_\ell - \frac{\lambda^2}{2} \int_s^t (\sigma')^2(\ell\lambda^2, \lambda Z_\ell + x_0) d\ell \right) \\ = & \mathcal{E}(\lambda; s, t) \left(\lambda \sigma'(s_1\lambda^2, \lambda Z_{s_1} + x_0) \mathbb{1}_{[s,t]}(s_1) \right. \\ & + \lambda^2 \sigma(s_1\lambda^2, \lambda Z_{s_1} + x_0) \int_s^t \mathbb{1}_{[s_1,t]}(\ell) \sigma''(\ell\lambda^2, \lambda Z_\ell + x_0) \mathcal{E}(\lambda; s_1, \ell) dW_\ell \\ & \left. - \lambda^3 \sigma(s_1\lambda^2, \lambda Z_{s_1} + x_0) \int_s^t \mathbb{1}_{[s_1,t]}(\ell) (\sigma' \sigma'')(\ell\lambda^2, \lambda Z_\ell + x_0) \mathcal{E}(\lambda; s_1, \ell) d\ell \right). \end{aligned} \tag{7.31}$$

On the other hand, we get

$$\begin{aligned} D_{s_1} G_{den} = & 2\lambda \int_0^t \sigma \sigma'(\xi\lambda^2, \lambda Z_\xi + x_0) \mathbb{1}_{[s_1,t]}(\xi) \sigma(s_1\lambda^2, \lambda Z_{s_1} + x_0) \mathcal{E}(\lambda; s_1, \xi) \mathcal{E}^2(\lambda; \xi, t) d\xi \\ & + 2 \int_0^t \sigma^2(\xi\lambda^2, \lambda Z_\xi + x_0) \mathcal{E}(\lambda; \xi, t) D_{s_1} \mathcal{E}(\lambda; \xi, t) d\xi. \end{aligned} \tag{7.32}$$

Therefore, coming back to (7.30) and using (7.31), we obtain that

$$I_2 \leq 4 [J_1 + J_2 + J_3], \tag{7.33}$$

where

$$\begin{aligned}
 J_1 &= \mathbb{E} \left[\int_0^t ds \int_0^s ds_1 \left| \sigma'(s\lambda^2, \lambda Z_s + x_0) \sigma(s_1\lambda^2, \lambda Z_{s_1} + x_0) \mathcal{E}(\lambda; s_1, s) \frac{\mathcal{E}(\lambda; s, t)}{G_{den}} \right|^2 \right], \\
 J_2 &= \mathbb{E} \left[\int_{[0,t]^2} ds_1 ds \sigma^2(s\lambda^2, \lambda Z_s + x_0) \frac{\mathcal{E}^2(\lambda; s, t)}{G_{den}^2} \left(\mathbb{1}_{[s,t]}(s_1) \sigma'(s_1\lambda^2, \lambda Z_{s_1} + x_0) \right. \right. \\
 &\quad \left. \left. + \sigma(s_1\lambda^2, \lambda Z_{s_1} + x_0) \int_s^t \mathbb{1}_{[s_1,t]}(\ell) \sigma''(\ell\lambda^2, \lambda Z_\ell + x_0) \mathcal{E}(\lambda; s_1, \ell) \left[dW_\ell - \sigma' \sigma''(\ell\lambda^2, \lambda Z_\ell + x_0) d\ell \right] \right)^2 \right], \\
 J_3 &= \mathbb{E} \left[\int_{[0,t]^2} ds_1 ds \frac{\sigma^2(s\lambda^2, \lambda Z_s + x_0)}{G_{den}^4} \mathcal{E}^2(\lambda; s, t) (D_{s_1} G_{den})^2 \right],
 \end{aligned}$$

with $D_{s_1} G_{den}$ given in (7.32). In the sequel we will enumerate constants K_1 to K_{20} ; all those will not depend on x_0 or t , but eventually on T , σ and b . We start estimating J_1 . Since σ' is bounded, by Cauchy-Schwarz inequality, we have

$$J_1 \leq K_1 \left(\mathbb{E} \left[\int_0^t ds \int_0^s ds_1 \sigma^4(s_1\lambda^2, \lambda Z_{s_1} + x_0) \right] \mathbb{E} \left[\int_0^t ds \int_0^s ds_1 \mathcal{E}^4(\lambda; s, t) \frac{\mathcal{E}^4(\lambda; s_1, s)}{G_{den}^4} \right] \right)^{\frac{1}{2}}.$$

Since σ has linear growth, Lemma 7.2 and a further use of Cauchy-Schwarz inequality imply that, J_1 is bounded by

$$K_2 (1 + |x_0|^2) (\mathbb{E} [G_{den}^{-8}])^{\frac{1}{4}} \left(\mathbb{E} \left[\int_0^t ds \int_0^s ds_1 \mathcal{E}^{16}(\lambda; s, t) \right] \mathbb{E} \left[\int_0^t ds \int_0^s ds_1 \mathcal{E}^{16}(\lambda; s_1, s) \right] \right)^{\frac{1}{8}}.$$

Therefore, by Proposition 7.5 and Lemma 7.4, we obtain that

$$J_1 \leq K_3 (1 + |x_0|^2). \tag{7.34}$$

We go on with the analysis of J_2 . Since σ' , σ'' are bounded, we have

$$J_2 \leq K_4 \mathbb{E} \left[\int_{[0,t]^2} ds ds_1 \sigma^2(s\lambda^2, \lambda Z_s + x_0) \frac{\mathcal{E}^2(\lambda; s, t)}{G_{den}^2} \left(1 + \sigma^2(s_1\lambda^2, \lambda Z_{s_1} + x_0) M^2(s, s_1; t) \right) \right],$$

where $M(s, s_1; t) = \int_{s \vee s_1}^t \sigma''(\ell\lambda^2, \lambda Z_\ell + x_0) \mathcal{E}(\lambda; s_1, \ell) dW_\ell$, $t \geq s \vee s_1$, is a martingale having all moments because of Lemma 7.4. Since σ has linear growth, using Cauchy-Schwarz inequality and Lemma 7.2, we get

$$J_2 \leq K_5 (1 + |x_0|^4) \left(\mathbb{E} \left[\int_0^t ds \mathcal{E}^{16}(\lambda; s, t) \right] \mathbb{E} [G_{den}^{-16}] \right)^{\frac{1}{8}} \left(\int_{[0,t]^2} ds ds_1 \mathbb{E} [M^8(s, s_1; t)] \right)^{\frac{1}{4}}.$$

Then Burkholder-Davis-Gundy inequality, Lemma 7.4 and Proposition 7.5 imply

$$J_2 \leq K_6 (1 + |x_0|^4). \tag{7.35}$$

Finally we treat J_3 . Applying Cauchy-Schwarz inequality, we have

$$J_3 \leq \left(\mathbb{E} \left[\int_{[0,t]^2} ds ds_1 \frac{\sigma^4(s\lambda^2, \lambda Z_s + x_0)}{G_{den}^s} \mathcal{E}^4(\lambda; s, t) \right] \mathbb{E} \left[\int_{[0,t]^2} ds ds_1 (D_{s_1} G_{den})^4 \right] \right)^{\frac{1}{2}}.$$

Since σ has linear growth, again by Cauchy-Schwarz inequality and Lemma 7.2, we get

$$J_3 \leq K_7(1 + |x_0|^2) \left(\mathbb{E} [G_{den}^{-32}] \int_0^t ds \mathbb{E} [\mathcal{E}^{16}(\lambda; s, t)] \right)^{\frac{1}{8}} \left(\mathbb{E} \left[\int_0^t ds_1 (D_{s_1} G_{den})^4 \right] \right)^{\frac{1}{2}};$$

by Lemma 7.4 and Proposition 7.5 it follows

$$J_3 \leq K_8(1 + |x_0|^2) \left(\mathbb{E} \left[\int_0^t ds_1 (D_{s_1} G_{den})^4 \right] \right)^{\frac{1}{2}}. \tag{7.36}$$

Since σ' is bounded, (7.32) and Jensen's inequality give

$$\mathbb{E} \left[\int_0^t ds_1 (D_{s_1} G_{den})^4 \right] \leq K_9 (A_1 + A_2), \tag{7.37}$$

where

$$A_1 = \mathbb{E} \left[\int_0^t ds_1 \int_{s_1}^t d\xi \sigma^4(s_1\lambda^2, \lambda Z_{s_1} + x_0) \sigma^4(\xi\lambda^2, \lambda Z_\xi + x_0) \mathcal{E}^4(\lambda; s_1, \xi) \mathcal{E}^8(\lambda; \xi, t) \right],$$

$$A_2 = \mathbb{E} \left[\int_0^t ds_1 \int_0^t d\xi \sigma^8(\xi\lambda^2, \lambda Z_\xi + x_0) \mathcal{E}^4(\lambda; \xi, t) (D_{s_1} \mathcal{E}(\lambda; \xi, t))^4 \right].$$

Since σ has linear growth, Cauchy-Schwarz inequality and Lemma 7.2 imply that A_1 is bounded by

$$K_{10} (1 + |x_0|^8) \left(\mathbb{E} \left[\int_0^t ds_1 \int_{s_1}^t d\xi \mathcal{E}^8(\lambda; s_1, \xi) \mathcal{E}^{16}(\lambda; \xi, t) \right] \right)^{\frac{1}{2}}.$$

Again, by Cauchy-Schwarz inequality and Lemma 7.4, we obtain

$$A_1 \leq K_{11} (1 + |x_0|^8). \tag{7.38}$$

We proceed estimating A_2 . Using Cauchy-Schwarz inequality, A_2 is bounded by

$$K_{12} \left(\mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi \sigma^{16}(\xi\lambda^2, \lambda Z_\xi + x_0) \right] \mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi \mathcal{E}^8(\lambda; \xi, t) (D_{s_1} \mathcal{E}(\lambda; \xi, t))^8 \right] \right)^{\frac{1}{2}}.$$

Since σ has linear growth, Cauchy-Schwarz inequality and Lemma 7.2 lead to

$$A_2 \leq K_{13} (1 + |x_0|^8) \left(\mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi \mathcal{E}^{16}(\lambda; \xi, t) \right] \mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi (D_{s_1} \mathcal{E}(\lambda; \xi, t))^{16} \right] \right)^{\frac{1}{4}};$$

Lemma 7.4 implies

$$A_2 \leq K_{14} (1 + |x_0|^8) \left(\mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi (D_{s_1} \mathcal{E}(\lambda; \xi, t))^{16} \right] \right)^{\frac{1}{4}}. \quad (7.39)$$

Since σ' and σ'' are bounded, using (7.31) and Jensen's inequality, it follows that

$$\mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi (D_{s_1} \mathcal{E}(\lambda; \xi, t))^{16} \right] \leq K_{15} (R_1 + R_2), \quad (7.40)$$

where

$$R_1 = \mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi \mathcal{E}^{16}(\lambda; \xi, t) \right],$$

$$R_2 = \mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi \sigma^{16}(s_1 \lambda^2, \lambda Z_{s_1} + x_0) \mathcal{E}^{16}(\lambda; \xi, t) \left(M^{16}(s_1, \xi; t) + \int_{\xi \vee s_1}^t \mathcal{E}^{16}(\lambda; s_1, \rho) d\rho \right) \right],$$

and $M(s_1, \xi; t) = \int_{\xi \vee s_1}^t \sigma''(\rho \lambda^2, \lambda Z_\rho + x_0) \mathcal{E}(\lambda; s_1, \rho) dW_\rho$, $t \geq \xi \vee s_1$, is a square integrable martingale, taking into account Lemma 7.4. Again by Lemma 7.4, R_1 is uniformly bounded in t and x_0 . On the other hand, using Cauchy-Schwarz and Jensen's inequalities, R_2 is bounded by

$$K_{16} \left(\mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi \sigma^{32}(s_1 \lambda^2, \lambda Z_{s_1} + x_0) \mathcal{E}^{32}(\lambda; \xi, t) \right] \mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi \left(M^{32}(s_1, \xi; t) + \int_{\xi \vee s_1}^t \mathcal{E}^{32}(\lambda; s_1, \rho) d\rho \right) \right] \right)^{\frac{1}{2}}.$$

Since σ has linear growth, again by Cauchy-Schwarz inequality, Lemma 7.2 and Lemma 7.4, we get

$$R_2 \leq K_{17} (1 + |x_0|^{16}) \mathbb{E} \left[\int_{[0,t]^2} ds_1 d\xi M^{32}(s_1, \xi; t) \right].$$

Burkholder-Davis-Gundy inequality gives

$$R_2 \leq K_{18} (1 + |x_0|^{16}). \quad (7.41)$$

Coming back to (7.39), using (7.41) and (7.40), we obtain

$$A_2 \leq K_{19} (1 + |x_0|^{12}); \quad (7.42)$$

thus, replacing (7.38) and (7.42) in (7.37) and coming back to (7.36), imply

$$J_3 \leq K_{20} (1 + |x_0|^8). \quad (7.43)$$

Consequently, substituting (7.34), (7.35) and (7.43) in (7.33), item (ii) of Proposition 7.6 is established. \square

Returning to the proof of Proposition 5.1 and substituting in (7.27) the right-hand side of the first and the second item of Proposition 7.6, the inequality (5.2) is verified. Finally this concludes the proof of Proposition 5.1. \square

References

- [1] M. Arnaudon, A. Thalmaier, and F. Wang, *Gradient estimates and Harnack inequalities on non-compact Riemannian manifolds*, Stochastic Process. Appl. **119** (2009), no. 10, 3653–3670.
- [2] V. Barbu, M. Röckner, and F. Russo, *Probabilistic representation for solutions of an irregular porous media type equation: the irregular degenerate case*, Probab. Theory Related Fields **151** (2011), no. 1-2, 1–43.
- [3] G. I. Barenblatt, *On some unsteady motions of a liquid and gas in a porous medium*, Akad. Nauk SSSR. Prikl. Mat. Meh. **16** (1952), 67–78.
- [4] N. Belaribi, F. Couvelier, and F. Russo, *A probabilistic algorithm approximating solutions of a singular pde of porous media type*, Monte Carlo Methods and Applications **17** (2011), no. 4, 317–369.
- [5] S. Benachour, P. Chassaing, B. Roynette, and P. Vallois, *Processus associés à l'équation des milieux poreux*, Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4) **23** (1996), no. 4, 793–832.
- [6] P. Benilan and M. G. Crandall, *The continuous dependence on φ of solutions of $u_t - \Delta\varphi(u) = 0$* , Indiana Univ. Math. J. **30** (1981), no. 2, 161–177.
- [7] P. Blanchard, M. Röckner, and F. Russo, *Probabilistic representation for solutions of an irregular porous media type equation*, Ann. Probab. **38** (2010), no. 5, 1870–1900.
- [8] V. I. Bogachev, G. Da Prato, and M. Röckner, *Infinite-dimensional Kolmogorov operators with time-dependent drift coefficients*, Dokl. Akad. Nauk **419** (2008), no. 5, 587–591.
- [9] V. I. Bogachev, G. Da Prato, M. Röckner, and W. Stannat, *Uniqueness of solutions to weak parabolic equations for measures*, Bull. Lond. Math. Soc. **39** (2007), no. 4, 631–640.
- [10] H. Brezis and M. G. Crandall, *Uniqueness of solutions of the initial-value problem for $u_t - \Delta\varphi(u) = 0$* , J. Math. Pures Appl. (9) **58** (1979), no. 2, 153–163.
- [11] H. Brezis and A. Friedman, *Non linear parabolic equations involving measures as initial conditions.*, J. Math. Pures Appl. (9) **62** (1983), no. 1, 73–97.
- [12] F. Cavalli, G. Naldi, G. Puppo, and M. Semplice, *High-order relaxation schemes for nonlinear degenerate diffusion problems*, SIAM J. Numer. Anal. **45** (2007), no. 5, 2098–2119 (electronic).
- [13] E. Chasseigne and J. L. Vazquez, *Theory of extended solutions for fast diffusion equations in optimal classes of data. Radiation from singularities*, Arch. Ration. Mech. Anal. **164** (2002), no. 2, 133–187.
- [14] ———, *Extended solutions for general fast diffusion equations with optimal measure data*, Adv. Differential Equations **11** (2006), no. 6, 627–646.
- [15] J. Dolbeault and G. Toscani, *Fast diffusion equations: matching large time asymptotics by relative entropy methods*, Kinet. Relat. Models **4** (2011), no. 3, 701–716.
- [16] A. Figalli, *Existence and uniqueness of martingale solutions for SDEs with rough or degenerate coefficients*, J. Funct. Anal. **254** (2008), no. 1, 109–153.
- [17] A. Figalli and R. Philipowski, *Convergence to the viscous porous medium equation and propagation of chaos*, ALEA Lat. Am. J. Probab. Math. Stat. **4** (2008), 185–203.
- [18] C. Graham, Th. G. Kurtz, S. Méléard, S. Ph. Protter, M. Pulvirenti, and D. Talay, *Probabilistic models for nonlinear partial differential equations*, Lectures given at the 1st Session and Summer School held in Montecatini Terme, Lecture Notes in Mathematics (1995), 22–30.
- [19] M. A. Herrero and M. Pierre, *The Cauchy problem for $u_t = \Delta u^m$ when $0 < m < 1$* , Trans. Amer. Math. Soc. **291** (1985), no. 1, 145–158.
- [20] F. Hirsch, C. Profeta, B. Roynette, and M. Yor, *Peacocks and associated martingales, with explicit constructions*, Bocconi & Springer Series, vol. 3, Springer, Milan, 2011.
- [21] J. Jacod, *Théorèmes limite pour les processus*, École d'été de probabilités de Saint-Flour, XIII—1983, Lecture Notes in Math., vol. 1117, Springer, Berlin, 1985, pp. 298–409.
- [22] I. Karatzas and S. E. Shreve, *Brownian motion and stochastic calculus*, second ed., Graduate Texts in Mathematics, vol. 113, Springer-Verlag, New York, 1991.
- [23] T. Lukkari, *The fast diffusion equation with measure data*, NoDEA: Nonlinear Differential Equations and Appl. **19** (2012), no. 3, 329–6343.

- [24] P. Malliavin, *Stochastic analysis*, Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences], vol. 313, Springer-Verlag, Berlin, 1997.
- [25] S. De Marco, *On probability distributions of diffusions and financial models with non-globally smooth coefficients*, Ph. D. Thesis, Université Paris-Est Marne-la-Vallée and Scuola Normale Superiore di Pisa, 2010.
- [26] H. P. Jr. McKean, *Propagation of chaos for a class of non-linear parabolic equations.*, Stochastic Differential Equations (Lecture Series in Differential Equations, Session 7, Catholic Univ., 1967), Air Force Office Sci. Res., Arlington, Va., 1967, pp. 41–57.
- [27] D. Nualart, *The Malliavin calculus and related topics*, second ed., Probability and its Applications (New York), Springer-Verlag, Berlin, 2006.
- [28] R. Philipowski, *Interacting diffusions approximating the porous medium equation and propagation of chaos*, Stochastic Process. Appl. **117** (2007), no. 4, 526–538.
- [29] M. Pierre, *Nonlinear fast diffusion with measures as data*, Pitman Res. Notes Math. Ser. **149** (1987), 179–188.
- [30] M. Röckner and X. Zhang, *Weak uniqueness of Fokker-Planck equations with degenerate and bounded coefficients*, C. R. Math. Acad. Sci. Paris **348** (2010), no. 7-8, 435–438.
- [31] E. M. Stein, *Singular integrals and estimates for the Cauchy-Riemann equations*, Bull. Amer. Math. Soc. **79** (1973), 440–445.
- [32] D. W. Stroock and S. R. S. Varadhan, *Multidimensional diffusion processes*, Classics in Mathematics, Springer-Verlag, Berlin, 2006, Reprint of the 1997 edition.
- [33] A. S. Sznitman, *Topics in propagation of chaos*, École d’Été de Probabilités de Saint-Flour XIX—1989, Lecture Notes in Math., vol. 1464, Springer, Berlin, 1991, pp. 165–251.
- [34] J. L. Vazquez, *The porous medium equation*, Oxford Mathematical Monographs, The Clarendon Press Oxford University Press, Oxford, 2007, Mathematical theory.

Acknowledgments. The authors were partially supported by the ANR Project MASTERIE 2010 BLAN 0121 01. Part of the work was done during the stay of the second named author at Bielefeld University, SFB 701 (Mathematik). The authors acknowledge the stimulating remarks of an anonymous Referee and of the Editors. They are also grateful to Dr. Juliet Ryan for her precious help in correcting several language mistakes.