

On a nonlocal, nonlinear Schrödinger equation occurring in Plasma Physics

Thierry Colin

► **To cite this version:**

Thierry Colin. On a nonlocal, nonlinear Schrödinger equation occurring in Plasma Physics. V. Makhankov, I. Puzynin and O. Pashaev. Nonlinear evolution equations and dynamical systems, NEED's 92, 1992, Dubna, Russia. World Scientific, 1993. <hal-00848075>

HAL Id: hal-00848075

<https://hal.inria.fr/hal-00848075>

Submitted on 25 Jul 2013

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

On a nonlocal, nonlinear Schrödinger equation occurring in Plasma Physics

T. Colin*

CMLA, Ecole Normale Supérieure de Cachan et CNRS,
61 av. du Pdt Wilson, 94235 Cachan Cedex, France.

ABSTRACT

In this work, our goal is to study the Cauchy problem for a nonlocal, nonlinear Schrödinger equation occurring in Plasma Physics. We shall construct solutions in the energy space; we give some sufficient conditions on the initial data which ensure that the solutions are global and we show that in some cases, finite time blow up occurs. We prove the existence of standing waves, solution to these equations and provide some stability results.

1 Introduction

The goal of this work is to study the following system :

$$\begin{cases} i\phi_t + \Delta\phi = -\operatorname{div}(|\nabla\psi|^\sigma\nabla\psi), \\ \Delta\psi = \phi, \\ \phi(x, 0) = \phi_0(x), \end{cases} \quad (1)$$

for $t \geq 0$ and $x \in R^3$.

This system is equivalent to :

$$\begin{cases} i(\nabla\psi)_t + \Delta(\nabla\psi) = \nabla(-\Delta)^{-1}\operatorname{div}(|\nabla\psi|^\sigma\nabla\psi), \\ \nabla\psi(x, 0) = \nabla\psi_0(x). \end{cases} \quad (2)$$

The physical case occurs for $\sigma = 2$ and was introduced by Musher, Rubenchick and Zakharov⁵ as a model for a nonlinear Plasma.

The plan of the work is the following:

- 2 Existence and finite time blow-up results.
- 3 Standing waves.
- 4 Radial standing waves.
- 5 Return to the finite time blow-up.

*This work was partially supported by a grant from the Ministère de la Recherche et de la Technologie, France, Contrat no 90S0315.

6 Stability for $\sigma < 4/3$.

7 Instability for $\sigma > 4/3$.

We will not give complete proofs, we refer the reader to Colin² for more details.

2 Existence and finite time blow-up results.

If ϕ is a solution to Eq. 1, one can show that the following quantities are invariants of the motion:

$$m(t) = \int_{R^3} |\nabla\psi|^2 dx = m(0), \quad (3)$$

$$E(t) = \int_{R^3} \frac{1}{2} |\Delta\psi|^2 - \frac{1}{\sigma+2} |\nabla\psi|^{\sigma+2} = E(0). \quad (4)$$

Therefore, we define the energy space H as the completion of the set of C^∞ functions ψ with compact support endowed with the norm $|\nabla\psi|_{H^1}$ (H^1 denoting the standard Sobolev space on R^3). Because of the Sobolev imbedding, H turns out to be equal to $\{\psi \in L^6, \nabla\psi \in H^1\}$ normed by $|\nabla\psi|_{H^1}$.

With these notations, we have:

Theorem 1 *Let $0 < \sigma < 4$ and $\psi_0 \in H$, then there exists a unique maximal solution to Eq. 2 on $[0, T(\psi_0)[$ in $\mathcal{C}([0, T(\psi_0)[, H)$. The equalities Eq. 3 and Eq. 4 hold. Moreover, if $T(\psi_0) < \infty$ then*

$$\lim_{t \rightarrow T(\psi_0)} |\nabla\psi(t)|_{H^1} = +\infty.$$

If $0 < \sigma < 4/3$, then $T(\psi_0) = \infty$.

If $4/3 \leq \sigma < 4$ and if $|\psi_0|_H$ is sufficiently small, then $T(\psi_0) = \infty$.

Proof: For the proof of this Theorem, we remark that the operator $\nabla(\Delta)^{-1}div$ is a Calderón-Zygmund operator, hence it maps L^p into L^p for all $1 < p < \infty$. Therefore, we can use the standard methods for Schrödinger type equations (see Kato³ for example). We first prove that Eq. 2 is equivalent to the following integral equation:

$$\nabla\psi = \mathcal{T}(\psi), \quad (5)$$

where

$$\mathcal{T}(\psi) \equiv S(t)\nabla\psi_0 - i \int_0^t S(t-s)\nabla(-\Delta)^{-1}div(|\nabla\psi|^\sigma \nabla\psi)(s)ds, \quad (6)$$

$S(t)$ denoting the free propagator of the equation

$$i\nabla\psi_t + \Delta\nabla\psi = 0.$$

Using the $L^p - L^q$ estimates (see Kato³), one can prove that \mathcal{T} is a contraction in the convenient spaces, hence there exists a local solution to Eq. 2.

If $\sigma < 4/3$, using Eq. 3, 4 and a Gagliardo-Nirenberg inequality, we obtain a bound of ψ in the space H , and hence the solution is global.

If $\sigma \geq 4/3$ and if $|\nabla\psi_0|_{H^1}$ is small, one can show, using a Gagliardo-Nirenberg inequality, that $E(t)$ controls the L^2 norm of $\Delta\psi$, therefore the solution is global in time, see Colin² for details. \blacksquare

For the case where $|\nabla\psi_0|_H$ is not small, we have the following result:

Theorem 2 *i) Let $\psi_0 \in H$ be such that*

$$\int |x|^2 |\nabla\psi_0|^2 < \infty,$$

then the solution ψ to Eq. 2 satisfies:

$$|x||\nabla\psi| \in L^\infty(0, t, L^2),$$

for all $t < T(\psi_0)$.

ii) If $\sigma \geq 4/3$, there exists some radial initial data such that the corresponding solutions blow up in finite time.

Sketch of proof:

i) In order to prove that the quantity $|x||\nabla\psi|$ $|_{L^2}$ persists, we work on the integral equation Eq. 5. We have to estimate the commutator of x and $\nabla(-\Delta)^{-1}div$. This can be done using the Fourier transform and some estimates on the Riesz potential (see Stein⁷). Next, we show that if $|x||\nabla\psi_0| \in L^2(R^3)$, then there exists $T_0 > 0$ such that $|x||\nabla\psi| \in L^\infty(0, t, L^2)$ for all $t < T_0$. Then it is sufficient to prove that $T(\psi_0)$ and T_0 are equal. This is done by supposing the contrary and working on Eq. 5.

ii) To prove the second part of Theorem 2, we make two elementary remarks:

- 1) $P = \nabla(-\Delta)^{-1}div$ is the projector on the set of the gradients in $(L^2(R^3))^3$.
- 2) Every radial vector of $(L^2(R^3))^3$ is a gradient.

These remarks prove that the restriction of Eq. 2 to the radial functions is:

$$i(\nabla\psi)_t + \Delta\nabla\psi = |\nabla\psi|^\sigma \nabla\psi.$$

This equation satisfies the Virial identity:

$$\frac{d^2}{dt^2} \int |x|^2 |\nabla\psi|^2 dx = 16E(\psi_0) + \frac{16}{\sigma+2} \left(1 - \frac{3\sigma}{4}\right) \int |\nabla\psi|^{\sigma+2}, \quad (7)$$

where

$$E(\psi_0) = \int \frac{1}{2} |\Delta\psi_0|^2 - \frac{1}{\sigma+2} |\nabla\psi_0|^{\sigma+2}.$$

Hence, in the case where $\sigma \geq \frac{4}{3}$ and $E(\psi_0) < 0$, the solution has to blow up in finite time. \blacksquare

3 Standing waves.

The goal of this section is to find $\psi(x)$ and $\omega \in R$ such that $e^{i\omega t}\psi(x)$ is a solution to Eq. 2 .

Theorem 3 *For all $0 < \sigma < 4$ and all $\omega > 0$, there exists a function $\psi_\omega \in H$ satisfying:*

$$\begin{cases} -\omega\phi_\omega + \Delta\phi_\omega = -div(|\nabla\psi_\omega|^\sigma\nabla\psi_\omega), \\ \Delta\psi_\omega = \phi_\omega, \end{cases}$$

i.e. $e^{i\omega t}\psi_\omega(x)$ is a solution to Eq. 2.

Proof: We solve the following minimization problem:

$$inf\{-\int |\nabla\psi|^{\sigma+2}, \psi \in H, \int |\nabla\psi|^2 = \lambda, \int |\Delta\psi|^2 = \mu\},$$

where λ and μ are convenient parameters. In order to solve this problem we take a minimizing sequence and thanks to the concentration-compactness method of P.L. Lions⁴, we prove that this sequence is compact up to translations, providing a minimum ψ . Writing the Euler-Lagrange equation satisfied by ϕ , there exists two Lagrange multipliers α and β such that:

$$\alpha\Delta\psi + \beta\Delta^2\psi = -div(|\nabla\psi|^\sigma\nabla\psi). \quad (8)$$

We now make a scaling on ψ , and using the so-called Pohozahev identity (i.e. multiplying Eq. 8 by $x.\nabla\psi$ and integrating over all R^3), we find that ω has to be positive. Hence $e^{i\omega t}\psi(x)$ is a standing wave solution to Eq. 1. ■

4 Radial Standing waves.

The goal of this section is to prove that there exists some radial standing waves solution to the system 1. Our result in this direction reads as follows.

Theorem 4 *For all $\omega > 0$ and $0 < \sigma < 4$, there exists radial functions ψ_ω satisfying*

$$\begin{cases} -\omega\phi_\omega + \Delta\phi_\omega = -div(|\nabla\psi_\omega|^\sigma\nabla\psi_\omega), \\ \Delta\psi_\omega = \phi_\omega. \end{cases}$$

Method:

The system 1 in the radial case is equivalent to the following ordinary differential equation:

$$f'' + \frac{2}{r}f' - (\omega + \frac{2}{r^2})f = -|f|^\sigma f,$$

where $f = \psi_r$.

We find a solution to this equation by considering the following minimization problem:

$$\inf \left\{ - \int f^{\sigma+2} r^2 dr, f \in H_r^1(R^3) \text{ such that } \int f'^2 r^2 dr + 2 \int f^2 dr = \mu, \int f^2 dr = \lambda \right\},$$

where $H_r^1(R^3)$ denotes the set of all radial functions which are in the Sobolev space $H^1(R^3)$. Using the compactness of $H_r^1(R^3) \subset L^{\sigma+2}$, due to Strauss⁸, we are able to prove that the minimizing sequences are compact in $H^1(R^3)$, providing therefore a minimum. We conclude as in the previous section. \blacksquare

There is an open question: are the two family of standing waves the same?

5 Return to the finite time blow up.

Thanks to the previous section, we prove:

Proposition 1 *Let $\sigma = 4/3$, then for all $t_0 > 0$, there exists a radial function ψ_0 such that the corresponding solution to Eq. 2 blows up exactly at $t = t_0$.*

Proof: The restriction of Eq. 2 to the radial functions is:

$$iu_t + \Delta u - \frac{2}{|x|^2} u = -|u|^{4/3} u, \quad (9)$$

where $u = \psi_r$.

This equation satisfies the pseudo-conformal transformation law, *i.e.* if u is a solution to Eq. 9 then

$$v(x, t) = \frac{1}{t^{3/2}} e^{\frac{i|x|^2}{4t}} \bar{u}\left(\frac{x}{t}, \frac{1}{t}\right)$$

is also a solution to Eq. 9.

We now apply this transformation with $u =$ any radial standing wave, and the proposition is proved. \blacksquare

6 Stability for $\sigma < 4/3$.

Theorem 5 *For $\sigma < 4/3$, the set S_ω of standing waves defined by the minimization problem of section 3 is stable:*

for all $\epsilon > 0$, there exists a $\delta > 0$ such that

$$\begin{aligned} \forall \psi_0 \in H, \quad \inf_{\psi_\omega \in S_\omega} \|\psi_0 - \psi_\omega\|_H &< \delta \\ \Rightarrow \inf_{\psi_\omega \in S_\omega} \|\psi(\cdot, t) - \psi_\omega\|_H &\leq \epsilon. \end{aligned}$$

Sketch of proof:

We first proved that the problems

$$(P1) \quad \inf\left\{-\int |\nabla\psi|^{\sigma+2}, \psi \in H, \int |\nabla\psi|^2 = \int |\nabla\psi_\omega|^2, \right. \\ \left. \int |\Delta\psi|^2 = \int |\Delta\psi_\omega|^2\right\}$$

and

$$(P2) \quad \inf\left\{\frac{1}{2}\int |\Delta\psi|^2 - \frac{1}{\sigma+2}\int |\nabla\psi|^{\sigma+2} + \frac{\omega}{2}\int |\nabla\psi|^2, \right. \\ \left. \psi \in H, \int |\nabla\psi|^2 = \int |\nabla\psi_\omega|^2\right\},$$

where ψ_ω is any standing wave that we found in section 3, are equivalent.

In order to do that, we use the concentration-compactness lemma to prove that (P2) has a solution, and a careful scaling analysis shows that this solution is also a solution to (P1). This method gives us the compactness of the minimizing sequences, which implies the stability according to the method of Cazenave and Lions¹. \blacksquare

7 Instability for $\sigma > 4/3$.

In this case, we prove:

Theorem 6 *For $\sigma > 4/3$, the orbit $e^{i\theta}\psi_\omega$ is unstable, i.e. there exists $\epsilon_0 > 0$ and $\psi_n^0 \rightarrow \psi_\omega$ in H such that*

$$\sup_{t>0} \inf_{\theta \in \mathbb{R}} \|\psi_n(\cdot, t) - e^{i\theta}\psi_\omega\|_H \geq \epsilon_0.$$

Sketch of proof:

The problems

$$(P1) \quad \inf\left\{-\int |\nabla\psi|^{\sigma+2}, \psi \in H, \int |\nabla\psi|^2 = \int |\nabla\psi_\omega|^2, \right. \\ \left. \int |\Delta\psi|^2 = \int |\Delta\psi_\omega|^2\right\}$$

and

$$(P3) \quad \inf\left\{\frac{1}{2}\int |\Delta\psi|^2 - \frac{1}{\sigma+2}\int |\nabla\psi|^{\sigma+2} + \frac{\omega}{2}\int |\nabla\psi|^2, \right. \\ \left. \psi \in H, \int |\Delta\psi|^2 = \int |\Delta\psi_\omega|^2\right\},$$

where ψ_ω is any standing wave, are equivalent (we apply the same method than in the previous section).

We then define

$$d(\omega) = \inf \left\{ \frac{1}{2} \int |\Delta \psi|^2 - \frac{1}{\sigma + 2} \int |\nabla \psi|^{\sigma+2} + \frac{\omega}{2} \int |\nabla \psi|^2, \right. \\ \left. \int |\Delta \psi|^2 = \int |\Delta \psi_\omega|^2 \right\}.$$

A simple scaling leads to

$$d(\omega) = \frac{\omega^{\frac{2}{\sigma}-\frac{1}{2}}}{3} \int |\Delta \psi_1|^2,$$

so that d is convex.

With these notations and results we are in the framework of Shatah and Strauss⁶, which leads to Theorem 6. ■

References:

1. T. Cazenave, P.L. Lions, Orbital stability of standing waves for some nonlinear Schrödinger equations. *Comm. Math. Phys.* **85**, 1982, p549-561.
2. T.Colin, On a nonlocal, nonlinear Schrödinger equation occuring in plasma Physics. *Prépublications du CMLA n° 9208, ENS de Cachan. To appear.*
3. T.Kato, On nonlinear Schrödinger equations. *Ann. Inst. Henri Poincaré, Physique théorique*, **46**, n° 1, 1987, p113-129.
4. P.L. Lions, The concentration-compactness principle in the calculus of variations. part1: *Ann. Inst. Henri Poincaré, analyse non linéaire*, **vol1**, n° 2, 1984 p109-145.
part2: *Ann. Inst. Henri Poincaré, analyse non linéaire*, **vol1**, n° 4, 1984 p223-283.
5. S.L. Musher, A.M. Rubenchick, V.E. Zakharov, Hamiltonian approach to the description of nonlinear plasma phenomena. *Physics report* **129**, n° 5, 1985, p285-366.
6. J.Shatah, W.Strauss, Instability of nonlinear bound states. *Comm. Math. Phys.* **100**, 1985, p173-190.
7. E.M. Stein, Singular integrals and differentiability properties of functions. *Princeton University Press.*
8. W.Strauss, Existence of solitary waves in higher dimensions. *Comm. Math. Phys.* **55**, 1977, p149-162.