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Modelling thermal adaptation in microalgae : an adaptive dynamics point of view

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Abstract: Controlling microalgae outdoor cultures is a growing challenge as these organisms could be used at the industrial scale to produce biofuels. In this context, understanding how temperature affects microalgae is a major field. Because of their high division rate, microalgae rapidly adapt to their environment through a process of selection - mutation so that it is not possible to neglect the evolution part. We present here a simple Monod-like model to account for temperature effect on microalgae. We then use it in an evolutionary perspective thanks to the Adaptive Dynamics theory to understand how temperature drives evolution. We analyze the model for a constant temperature, showing that the optimal temperature trait tends to equal the environment temperature. We then study the case where the temperature is periodically fluctuating, and we find a stable equilibrium. Finally, we simulate the model to show that evolutionary branching in two distinct morphs can occur under this condition, which could be a first step to find a criterion for strain selection.

Keywords: temperature, evolution, adaptive dynamics theory, microalgae

1. INTRODUCTION

Autotrophic micro-organisms called ‘microalgae’ brought together a variety of unicellular photosynthetic eukaryotes. Some of them are able to accumulate an important amount of neutral lipids (Mata et al., 2010), which can be turned into biofuel, others can store important quantities of carbohydrates, used for methane production. Thus, they are a promising biotechnology resource for the future (Wijffels et al., 2013). Their growth is mostly influenced by light, but temperature is the second determining factor because microalgae are ectotherms (*i.e* they can’t regulate their internal temperature). The temperature effect on growth is often neglected in the modelling of microalgae. In regard to the increasing interest for outdoor production of microalgae, temperature effect on growth has to be investigated (Ras et al., 2013).

Because of their high division rate, microalgae are able to rapidly adapt to their environment through a process of selection - mutation. In outdoor cultures, the daily range of temperature fluctuation can be more than 20°C (Bechet et al., 2010). It is thus important to understand how temperature drives evolution. To date, there is only one model which predicts microalgae adaptation to temperature (Thomas et al., 2012), by simulations based on the Adaptive Dynamics theory. However, this model implies several assumptions that have been put in doubt (Boyd et al., 2013).

We propose here a simple Monod-like model that represents the effect of a constant temperature on growth. We employ it to model microalgae adaptation to temperature. In line with Thomas et al. (2012), we use the Adaptive

Dynamics theory. We keep the model as simple as possible to study it analytically. In a second part, we use the model under periodic temperature to account for more realistic conditions. Finally, we simulate strain separation through evolutionary branching under fluctuating temperature.

2. MODELLING TEMPERATURE EFFECT ON MICROALGAE GROWTH

The influence of temperature on all ectotherms growth is a well-known asymmetric curve, called thermal reaction norm (Fig. 1 A). In microalgae, two models have been used to describe temperature effect on growth (Norberg, 2004; Bernard and Rémond, 2012). We use the model from Bernard and Rémond (2012) based on the Cardinal Temperature Model with Inflection (CTMI) of Rosso et al. (1993) because its parameters can be directly estimated on the available data:

$$\mu(T) = \begin{cases} 0 & \text{for } T < T_{\min} \\ \mu_{opt} \frac{\lambda(T)}{\beta(T)} & \text{for } T_{\min} < T < T_{\max} \\ 0 & \text{for } T > T_{\max} \end{cases} \quad (1)$$

with

$$\lambda(T) = (T - T_{max})(T - T_{min})^2$$
$$\beta(T) = (T_{opt} - T_{min}) \left[(T_{opt} - T_{min})(T - T_{opt}) - (T_{opt} - T_{max})(T_{opt} + T_{min} - 2T) \right] \quad (2)$$

and

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$$T_{opt} > \frac{T_{min} + T_{max}}{2} \quad (3)$$

T corresponds to the temperature of the environment, $\mu(T)$ is the growth rate of a given species, T_{min} and T_{max} are the minimum and maximum temperatures that sustain growth, T_{opt} is the optimal temperature for growth and μ_{opt} is the theoretical maximum growth rate. Contrary to Norberg (2004) used in Thomas et al. (2012), this model does not imply any assumption concerning the structural link between the cardinal temperatures and the maximal growth rate (e.g. if T_{opt} is hotter, μ_{opt} will be higher).

In the following, we use a Monod-type model including the temperature effect on microalgae growth in a chemostat:

$$M : \begin{cases} \dot{S} = D(S_{in} - S) - \mu(T)\rho(S)X \\ \dot{X} = \mu(T)\rho(S)X - DX \end{cases} \quad (4)$$

where S is the nutrient concentration in the chemostat, X is the algal biomass concentration, D is the dilution rate and with:

$$\rho(S) = \frac{S}{K + S} \quad (5)$$

K is a half-saturation coefficient.

It is possible to calculate the non zero equilibrium (S^*, X^*) of system (M):

$$S^* = \frac{KD}{\mu(T) - D} \quad (6)$$

$$X^* = (S_{in} - S^*)$$

with the hypothesis $\mu(T)\rho(S_{in}) - D > 0$.

In the following, we use a Lyapunov function taken from Harisson (1979) to prove that (6) is globally asymptotically stable.

Lemma 1. The Lyapunov candidate function is given:

$$V(S, X) = \int_{S^*}^S \frac{\mu(T)\rho(w) - \mu(T)\rho(S^*)}{\mu(T)\rho(w)} dw + \int_{X^*}^X \frac{w - X^*}{w} dw \quad (7)$$

with $V : B \rightarrow \mathbb{R}^2$ where B is an open containing (S^*, X^*) . $V(S, X)$ is zero at (S^*, X^*) , positive at all other points (S, X) , defined and monotone increasing when $|X - X^*|$ or $|S - S^*|$ increases. The time derivative of V is:

$$\dot{V}(S, X) = \frac{1}{\mu(T)\rho(S)} (D - \mu(T)\rho(S)) [\mu(T)\rho(S) (S_{in} - S^*) - D(S_{in} - S)] \quad (8)$$

Proof: It is obvious that $V(S^*, X^*) = 0$. Moreover, since the integrands are of the same sign as $X - X^*$ and $S - S^*$, the integrals are positive and increasing as $|X - X^*|$ and $|S - S^*|$ increase.

Proposition 2. If $D < \mu(T)\rho(S_{in})$, System (M) admits one non zero equilibrium which is globally asymptotically stable.

Proof: It is sufficient to prove that $(8) < 0 \forall (S, X) \in B, (S, X) \neq (S^*, X^*)$.

If $S > S^*$ (resp. $S < S^*$), then $D - \mu(T)\rho(S) < 0$ (resp. > 0) because $\mu(T)\rho(S^*) = D$, and $[\mu(T)\rho(S)(S_{in} - S^*) - D(S_{in} - S)] > 0$ (resp. < 0) because $S_{in} - S < S_{in} - S^*$ (resp. $S_{in} - S > S_{in} - S^*$). Thus, $(8) < 0$ is verified, and (S^*, X^*) is globally asymptotically stable.

3. EVOLUTIONARY MODEL

3.1 General case

We now study the system (M) in the context of adaptive dynamics, introducing a mutant X_{mut} with an adaptive trait a_{mut} (i.e. a quantifiable trait that is likely to evolve) different from the resident trait a , with $\mu_{mut}(a_{mut}, T)$ and $\mu(a, T)$ (Dieckmann and Law, 1996). For further introduction to adaptive dynamics, see Dieckmann (2004). To define the canonical equation of adaptive dynamics, i.e. the equation of the evolution of trait a (Dieckmann and Law, 1996), we need to find the mutant growth rate (*per capita*) in the resident population at equilibrium, called invasion fitness $f(a_{mut}, a)$.

Proposition 3. The invasion fitness for System (4) is given by:

$$f(a_{mut}, a) = D \left(\frac{\lambda_{mut}}{\lambda} \frac{\beta}{\beta_{mut}} - 1 \right) \quad (9)$$

Proof: This results from the steady-state condition of the resident system, and from the hypothesis that the mutant is initially rare. We can thus replace S by its equilibrium value S^* :

$$f(a_{mut}, a) = \mu_{mut}(a_{mut}, T)\rho(S^*) - D \quad (10)$$

3.2 Modelling the evolution of the optimal temperature trait

We choose to study the adaptive trait $a = T_{opt}$, assuming that temperature will mainly affect the optimal conditions for growth. Because of the constraint (3) on T_{opt} , we choose to study the case:

$$T > \frac{T_{min} + T_{max}}{2} \quad (11)$$

We calculate the selection gradient, which gives the direction of the selection, using (9):

$$\left. \frac{\partial f(a_{mut}, a)}{\partial a_{mut}} \right|_{a_{mut}=a} = -\frac{\beta'(a)}{\beta(a)} D \quad (12)$$

with

$$\begin{aligned} \beta'(a) = & -6a^2 + (6T + 2T_{max} + 4T_{min})a \\ & + (-2T_{max} - 4T_{min})T \end{aligned} \quad (13)$$

We can deduce the canonical equation of the adaptive trait, that describes the evolution of T_{opt} at the evolutionary time scale θ :

$$\frac{da}{d\theta} = -M_p \sigma^2 X^* D \frac{\beta'(a)}{\beta(a)} \quad (14)$$

where M_p is the probability to be a mutant at each apparition, and σ is the mutation step.

We then search for the evolutionary singular strategy. We find that $da/d\theta = 0$ for $a^* = T$, which means that the optimal temperature trait tends to equal the environment temperature (Fig. 1 B). Note that the same results can be obtained with the Droop model (not presented here for sake of brevity).

We investigate the stability of the singular strategy examining the sign of the fitness invasion second order derivative.

Proposition 4. a^* is Convergent Stable Strategy (CSS), which means that the singular strategy is attractive, and Evolutionary Stable Strategy (ESS), which means that the resident a^* can't be invade by another mutant.

Proof: The following conditions are respected:

$$\mathbf{H1.} \quad \left. \frac{\partial^2 f(a_{mut}, a)}{\partial a_{mut}^2} \right|_{a_{mut}=a=a^*} < \left. \frac{\partial^2 f(a_{mut}, a)}{\partial a^2} \right|_{a_{mut}=a=a^*}$$

$$\mathbf{H2.} \quad \left. \frac{\partial^2 f(a_{mut}, a)}{\partial a_{mut}^2} \right|_{a_{mut}=a=a^*} < 0$$

Indeed, we have:

$$\begin{aligned} \left. \frac{\partial^2 f(a_{mut}, a)}{\partial a_{mut}^2} \right|_{a_{mut}=a=a^*} &= -D \frac{\beta''\beta - 2\beta'^2}{\beta^2} \\ \left. \frac{\partial^2 f(a_{mut}, a)}{\partial a^2} \right|_{a_{mut}=a=a^*} &= D \frac{\beta''}{\beta} \end{aligned} \quad (15)$$

$\beta'(a^*) = 0$, and so it is sufficient to prove:

$$-D \frac{\beta''(a^*)}{\beta(a^*)} < 0 \quad (16)$$

Yet, we have $\beta(a^*) = -(T - T_{min})(T - T_{max})(T_{min} - T) < 0$ because $T_{min} < T < T_{max}$.

Also, $\beta''(a^*) = 2T_{max} + 4T_{min} - 6T$. $\beta''(a^*) < 0$ is equivalent to $T > (2T_{min} + T_{max})/3$ which is true because of (11). Thus, (16) is true. H1 and H2 are verified.

3.3 Structural link between adaptive traits

In a real case of evolution, it is possible that several adaptive traits evolve concurrently, *e.g.* T_{max} and T_{opt} . In bacteria, Rosso et al. (1993) point out that there exists a structural link between T_{max} and T_{opt} . That is, T_{max} can be written as a linear function of T_{opt} . We compiled data from Miller and Castenholz (2000) for cyanobacteria and from Bernard and Rémond (2012) for microalgae to verify whether this relation is true for these organisms (Fig. 2). The data reveal a good correlation:

$$T_{max} = mT_{opt} + p \quad (17)$$

with $m = 0.93$ and $p = 9.83$, $r = 0.941$ (Fig. 2). This property must depend on the genetic interactions between traits, but the exact reason for such a link is still to understand.

In expression (9), we replace T_{max} by (17). By taking the adaptive trait $a = T_{opt}$, the selection gradient becomes:

$$\left. \frac{\partial f(a_{mut}, a)}{\partial a_{mut}} \right|_{a_{mut}=a} = D \frac{\lambda'(a)\beta(a) - \beta'(a)\lambda(a)}{\lambda(a)\beta(a)} \quad (18)$$

with

$$\lambda'(a) = -m(T - T_{min})^2 \quad (19)$$

We search for the evolutionary singular strategy by setting Eq.(18) equal to zero. As for (14), we find that $a^* = T$, and thus $T_{max}^* = mT + p$ (Fig. 1 C). This is an important result in the context of co-evolving species, because whether T_{max} evolve or not, the thermal reaction norm will differently affect the species fitness in its environment, and so its competition with other species or stains.

4. FLUCTUATING TEMPERATURE

4.1 Ecological timescale

We now study the system (M) in the context of fluctuating temperature:

$$\begin{cases} T = T_{inf} & \text{if } t \bmod \tau \in [\epsilon; \tau_1 - \epsilon[\\ T = T_{supp} & \text{if } t \bmod \tau \in [\tau_1 + \epsilon; \tau - \epsilon] \end{cases} \quad (20)$$

with $T_{inf} < T_{supp}$. By applying the conservation principle, we have the equality $X + S = S_{in}$ at equilibrium, that is to say $S = S_{in} - X$. Thus, the system (M) can be reduced to a one dimension differential equation:

$$\dot{X} = [\mu(T(t))\rho(S_{in} - X) - D]X \quad (21)$$

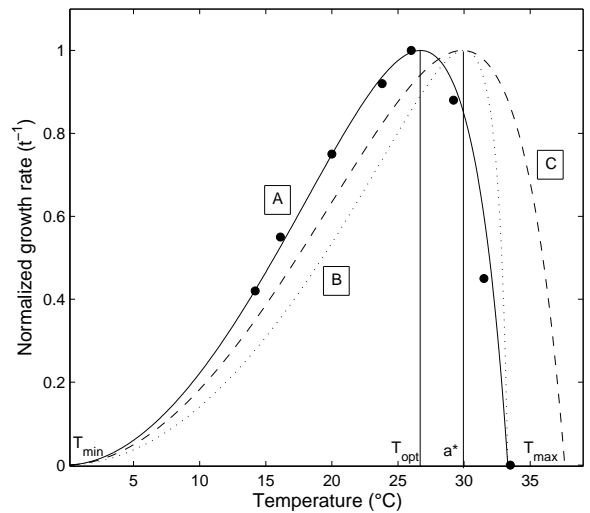


Fig. 1. Thermal reaction norm for *Nannochloropsis oceanica* (A). (B) and (C) are the evolutionary cases for $T = 30^\circ\text{C}$ with $a^* = T$, and with $T_{max} = ma^* + p$ for (C). The black circles are data from Sandnes et al. (2005) for *N. oceanica*.

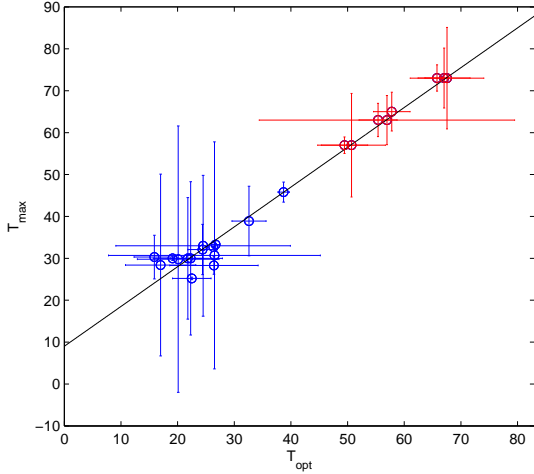


Fig. 2. Linear relationship between T_{max} and T_{opt} . Data from Bernard and Rémond (2012) are represented in blue, data from Miller and Castenholz (2000) are in red. From Miller and Castenholz (2000), we took data for the *Synechococcus* strains C9, SH94, OH2,20,30,29,28,4. Vertical and horizontal bars errors are represented, computed with the same method as Bernard and Rémond (2012).

where $T(t)$ is given by Eq. (20). We define $g(t, X) \stackrel{\text{def}}{=} \mu(T(t))\rho(S_{in} - X) - D$.

We follow the same reasoning as Butler et al. (1985) and Butler and Freedman (1981) who study a similar Monod-type model, but with a time varying dilution rate $D(t)$ instead, and a predator-prey system with periodic coefficients, respectively.

Theorem 1: For $D < \min(\mu(T)\rho(S_{in}))$, equation (21) has a unique nontrivial positive periodic solution $\psi(t)$ which is globally orbitally asymptotically stable. Moreover, we have $\min_{s \in [0, \tau]} A(s) \leq \psi(t) \leq \max_{s \in [0, \tau]} A(s)$ with

$$A(t) := [D(S_{in} + K) - \mu(T(t))S_{in}] / (D - \mu(T(t)))$$

(which corresponds to the steady state biomass concentration for constant T).

Proof: $\partial g / \partial X$ exists and is continuous for $(t, X) \in R^1 \times R_+^1$ with $R_+^1 = \{X \in R^1 : X \geq 0\}$.

Moreover, $\exists A(t) > 0 \ni [X - A(t)]g(t, X) < 0 \quad \forall X > 0, X \neq A(t)$. Indeed, given $D < \min(\mu(T)\rho(S_{in}))$, we have:

$$\begin{aligned} [X - A(t)]g(t, X) &< 0 \\ \forall X > 0, X \neq A(t) \end{aligned} \quad (22)$$

Thus, Massera's theorem (Massera, 1950) easily implies the existence of a periodic solution $\psi(t)$ of (21) satisfying $\min_{s \in [0, \tau]} A(s) \leq \psi(t) \leq \max_{s \in [0, \tau]} A(s)$.

$\psi(t)$ is the unique solution of (21), given that $X \partial g(t, X) / \partial X < 0$ for all $(t, X) \in R^1 \times R_+^1$. Indeed, we have:

$$X \frac{\partial g(t, X)}{\partial X} = \mu(T)X \frac{-K}{(S_{in} - X + K)^2} \quad (23)$$

Because $X > 0, K > 0, X \partial g(t, X) / \partial X < 0$ is always verified. Following the same reasoning as Butler and Freedman (1981), $X \partial g(t, X) / \partial X < 0$ implies that $g(t, X)$ is strictly

decreasing as a function of X , for $X > 0$, for all t . So, if two solutions $\psi(t)$ and $\psi_2(t)$ exist, with $\psi(t) < \psi_2(t)$, it implies that $\psi'(t)/\psi(t) = g(t, \psi(t)) > g(t, \psi_2(t)) = \psi_2'(t)/\psi_2(t)$ for all t . Integrating this inequality over $[0, \tau]$ leads to a contradiction which prove the uniqueness of $\psi(t)$.

Using theorem 1, we have the following inequality with the periodic temperature (20):

$$\begin{aligned} \frac{D(S_{in} + K) - \mu(T_{inf})S_{in}}{D - \mu(T_{inf})} &\leq X(t) \\ &\leq \frac{D(S_{in} + K) - \mu(T_{supp})S_{in}}{D - \mu(T_{supp})} \end{aligned} \quad (24)$$

which means that

$$\frac{DK}{\mu(T_{inf}) - D} \leq S(t) \leq \frac{DK}{\mu(T_{supp}) - D} \quad (25)$$

If the time spent at each temperature is sufficiently long, then, when $T = T_{inf}$ (resp. $T = T_{supp}$), the substrate concentration converges towards its equilibrium $S_{inf}^* = DK / (\mu(T_{inf}) - D)$ (resp. $S_{supp}^* = DK / (\mu(T_{supp}) - D)$). We assume that the state transition between S_{inf}^* and S_{supp}^* is negligible. From a biological point of view, this assumption implies that the growth rates at both temperatures are quite similar.

4.2 Evolutionary timescale

As explained previously, we assume that the resident population reaches rapidly its equilibrium for each temperature applied. In the case with the adaptive trait $a = T_{opt}$ and considering a constant T_{max} , we obtain:

$$\begin{cases} f(a_{mut}, a) = D \left(\frac{\beta(T_{inf})}{\beta_{mut}(T_{inf})} - 1 \right) & \text{if } T = T_{inf} \\ f(a_{mut}, a) = D \left(\frac{\beta(T_{supp})}{\beta_{mut}(T_{supp})} - 1 \right) & \text{if } T = T_{supp} \end{cases} \quad (26)$$

It is possible to use the average mutant growth rate in the resident population at steady-state (Ripa and Dieckmann, 2013):

$$\bar{f}(a_{mut}, a) = \frac{1}{\tau} \int_0^\tau f(a_{mut}, a; t) dt$$

which is equivalent to:

$$\begin{aligned} \bar{f}(a_{mut}, a) = D \left[\frac{\beta(a, T_{inf})}{\beta(a_{mut}, T_{inf})} \frac{\tau_1}{\tau} \right. \\ \left. + \frac{\beta(a, T_{supp})}{\beta(a_{mut}, T_{supp})} \frac{\tau - \tau_1}{\tau} - 1 \right] \end{aligned} \quad (27)$$

We thus deduce the selection gradient:

$$\left. \frac{\partial \bar{f}(a_{mut}, a)}{\partial a_{mut}} \right|_{a_{mut}=a} = D \left[- \frac{\tau_1}{\tau} \frac{\beta'(T_{inf})}{\beta(T_{inf})} - \frac{\tau - \tau_1}{\tau} \frac{\beta'(T_{supp})}{\beta(T_{supp})} \right] \quad (28)$$

Table 1. Model parameters for evolutionary branching.

Parameters	Unit
T_{min} , minimal temperature for growth	4°C
T_{max} , maximal temperature for growth	40°C
T_{inf} , inferior temperature applied	20°C
T_{supp} , superior temperature applied	36°C
τ_1 , time during which T_{inf} is applied	12 h
τ , period of temperature fluctuation	24 h

and so the canonical equation of the adaptive trait with fluctuating temperature is:

$$\frac{da}{d\theta} = M_p \sigma \bar{X}^* \left. \frac{\partial \bar{f}(a_{mut}, a)}{\partial a_{mut}} \right|_{a_{mut}=a} \quad (29)$$

Equation (29) is not analytically tractable. We perform simulations of (29) to search for the evolutionary singular strategy. Result (Fig. 3 C) shows that, at steady, T_{opt} does not converge to the average temperature T_{mean} . This result is still available if we replace T_{max} by a linear function of T_{opt} (Fig. 3 B), even if a^* is different. The asymmetric property of the thermal reaction norm is probably the reason for such outcome.

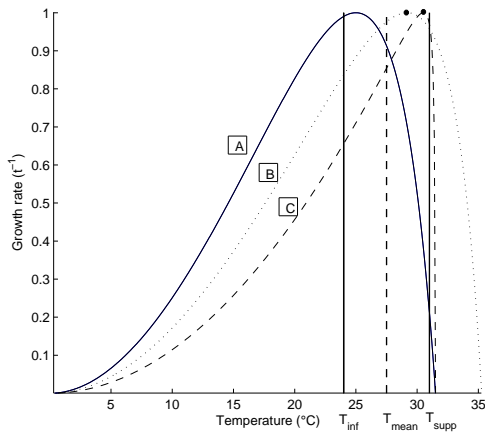


Fig. 3. Evolution of the thermal reaction norm for $T_{inf} = 24^\circ\text{C}$, $T_{supp} = 31^\circ\text{C}$. (A) is the initial reaction norm, (B) is the evolutionary reaction norm for $a = T_{opt}$ and $T_{max} = mT_{opt} + p$ at steady-state, and (C) is the evolutionary reaction norm for $a = T_{opt}$ and $T_{max} = 31.5^\circ\text{C}$ at steady state. The singular strategies are represented by black points.

4.3 Evolutionary Branching conditions

Under particular conditions, Adaptive Dynamics predict that evolution can converge to a specific singular strategy called ‘branching point’ where selection becomes disruptive, so that two strains move apart (Metz et al., 1992). The mutant can invade the resident and reciprocally, and both stably coexist.

We investigate the evolutionary dynamics of (3) by using a pairwise invasibility plot (PIP) (Fig. 4) (Geritz et al., 1998). This allow us to determine the stability of the singular strategy, studying graphically the sign of the mutant invasion fitness. For temperatures comprised between $T_{inf} = 20^\circ\text{C}$ and $T_{supp} = 36^\circ\text{C}$, and with $T_{min} = 4^\circ\text{C}$, $T_{max} = 40^\circ\text{C}$, we find that there exist a branching point (Fig. 4 A). Indeed, at this point, mutant and resident

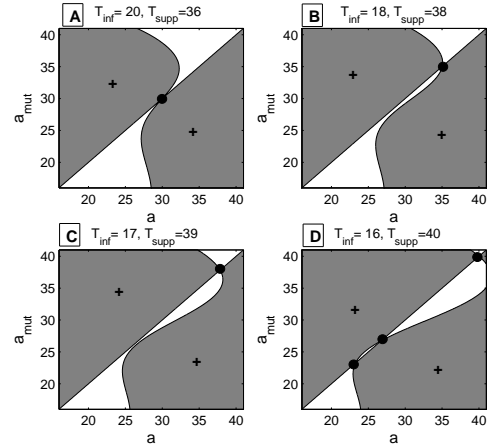


Fig. 4. Pairwise Invasibility Plots across a range of values of T_{inf} and T_{supp} (expressed in $^\circ\text{C}$). The other parameters are listed in Table 1. Grey and white represent positive and negative invasion fitness, respectively. Evolutionary singular strategy are represented by a black circle.

can mutually invade in such a way that two strains separate. For a range of values around the previous ones, we observe that the singular strategy bifurcates in three singular strategies (Fig. 4 B, C, D). In Fig. 4 D, there is still a branching point, which seems to allow a strong separation between strains because of the large area of positive invasion fitness.

To confirm the results found previously, we perform a simulation of an evolutionary branching. In line with Mirrahimi et al. (2011), we consider a model where the adaptive trait a becomes a continuous trait:

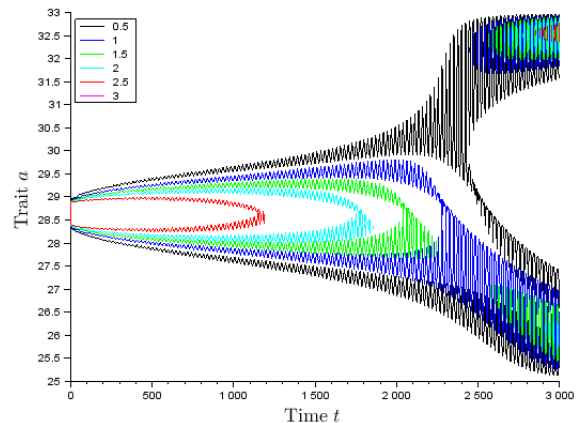


Fig. 5. Evolutionary branching of trait a for the parameters of Table 1. The trait a is expressed in $^\circ\text{C}$, t is arbitrarily expressed in hours. The colored lines correspond to a population size for each trait value.

$$\begin{cases} \partial_t X(a, t) = X(a, t)[\mu(a, T)S(t) - D] + \epsilon \Delta X(a, t) \\ S(t) = \frac{DS_{in}}{D + \int \mu(a, T)X(a, t) da} \end{cases} \quad (30)$$

where $X(a, t)$ is the species density with trait $a = T_{opt}$, $S(t)$ is the quasi-static approximation of resource dynamics, ϵ is the mutation rate. We use the parameters listed

in Table 1. Fig. 5 shows that an evolutionary branching really occurs. Two general morphs with two distinct T_{opt} appear and stabilize.

5. CONCLUSION

We have proposed a simple model of temperature effect on microalgae growth based on the model developed by Bernard and Rémond (2012). We used it in an evolutionary perspective thanks to the adaptive dynamics. We found that, under constant temperature, the optimal temperature tends to equal the environment temperature.

In a second part, we showed that there exists a structural link between T_{max} and T_{opt} . In the evolutionary sense, the optimal temperature also tends to the environment temperature, but the thermal reaction norm is different. Thus, the outcome of evolution if several microalgae species coexist is mostly influenced by this structural link, which may differ between species depending upon the genetics and physiological constraints.

Finally, we studied our model under a simple fluctuating temperature signal. We showed that a stable periodic solution exists. At evolutionary time scale, the fluctuating temperature allows strains to separate, and evolutionary branching occurs. We simulated the strain separation and found results consistent with our theoretical approach. This may be a first step to understand how species coexist under fluctuating temperature. It could serve to find a criterion for selecting species with the highest growth rate under particular temperature conditions, which is of key interest for microalgae outdoor production.

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