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# INFLUENCE OF PERIODIC DIFFUSIVE INCLUSIONS ON THE BIDOMAIN MODEL

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## Introduction

We present a new mathematical model of the electric activity of the heart. The main **drawback of the standard bidomain model** is that it assumes the existence of excitable cells (myocytes) everywhere in the heart, while it is known that there exist non-small regions where non-excitable cells (fibroblasts and collagen) take place. The problems that we are trying to address with our model are:

- The **laminar structure** of the myocardium that shows the presence of collagen, especially between muscle layers. See Fig 1. [2, 3]. In the diseased tissue the diffusion space plays an important role.

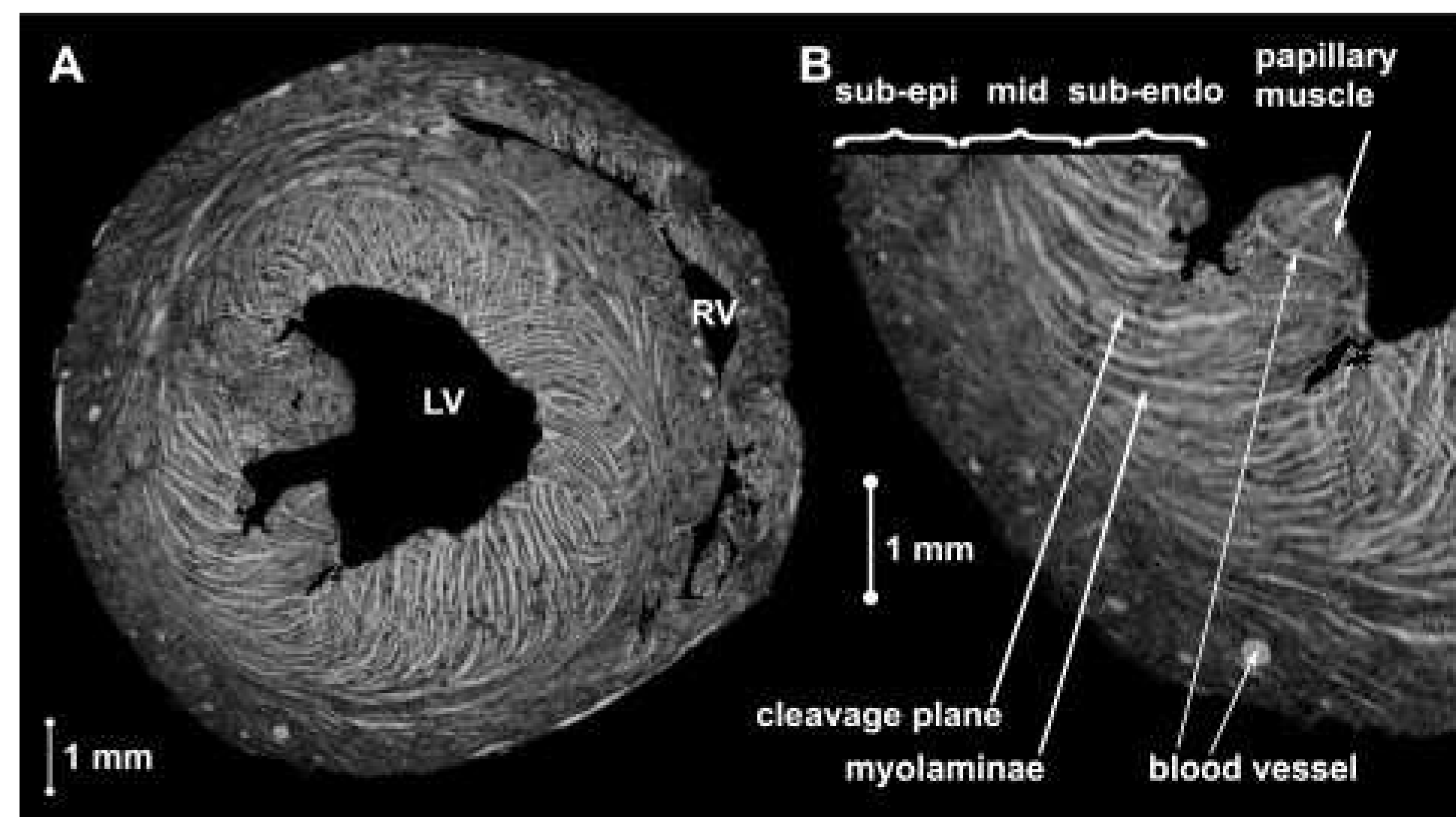


Figure 1: The laminar structure of myocardium. [3]

- The **infarct border zone**. After infarction in the heart some number of myocytes die and they are replaced by collagen and a few cells of fibroblasts. We consider that there are no cells of fibroblasts in the extracellular space.

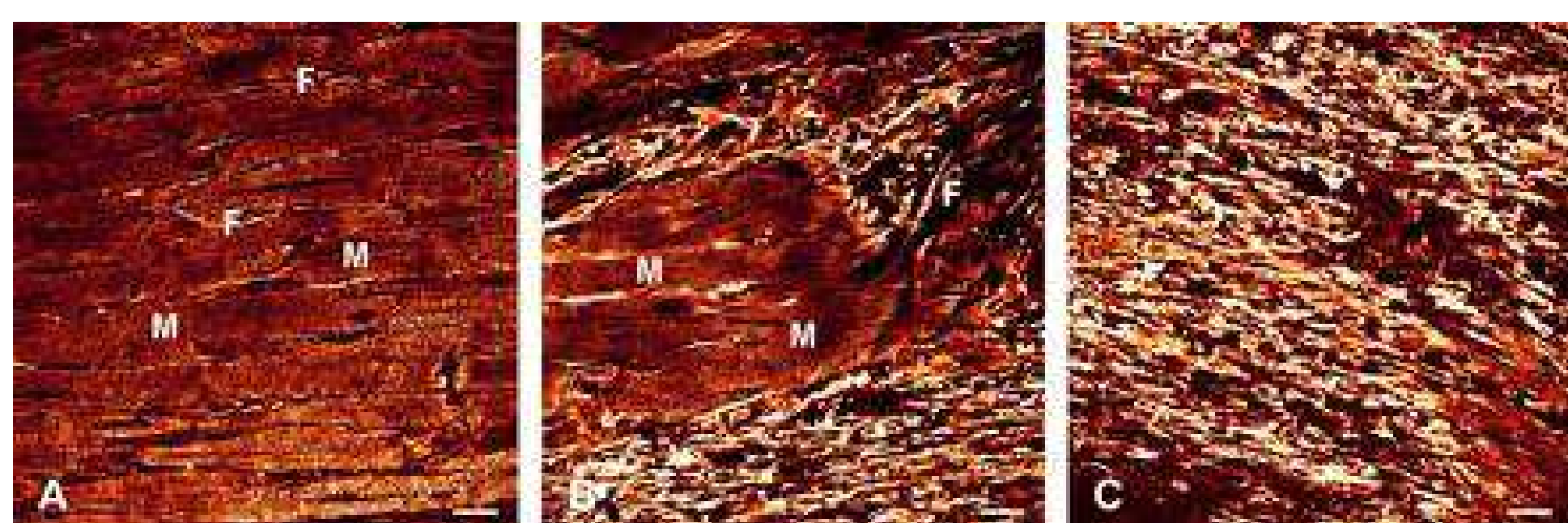


Figure 2: Fibroblast organization in sheep: normal ventricular myocardium (A), infarct border zone (B) and centre (C), 1 week after infarction. [4]

## I - Mesoscopic model

Modelling assumptions:

- Periodic distribution of added extracellular space
- Extracellular space is a passive conductor

The domain  $\Omega$  is split into two regions  $B_\varepsilon$  and  $D_\varepsilon$ :

- $B_\varepsilon$  represents the bidomain layer.
  - intra- and extra-cellular conductivities  $\sigma_\varepsilon^i$  and  $\sigma_\varepsilon^e$
  - unknown potentials  $u_\varepsilon^i$  and  $u_\varepsilon^e$
  - define  $v_\varepsilon := u_\varepsilon^i - u_\varepsilon^e$
- $D_\varepsilon$  represents the diffusive inclusions.
  - conductivity  $\sigma_\varepsilon^d$
  - unknown potential  $u_\varepsilon^d$
- $\Sigma_\varepsilon = \partial B_\varepsilon \cap \partial D_\varepsilon$  is the interface.

The bidomain model

$$\begin{aligned} \partial_t v_\varepsilon + cv_\varepsilon &= \nabla \cdot (\sigma_\varepsilon^i \nabla u_\varepsilon^i), & \text{in } B_\varepsilon, \\ \partial_t v_\varepsilon + cv_\varepsilon &= -\nabla \cdot (\sigma_\varepsilon^e \nabla u_\varepsilon^e), & \text{in } B_\varepsilon. \end{aligned}$$

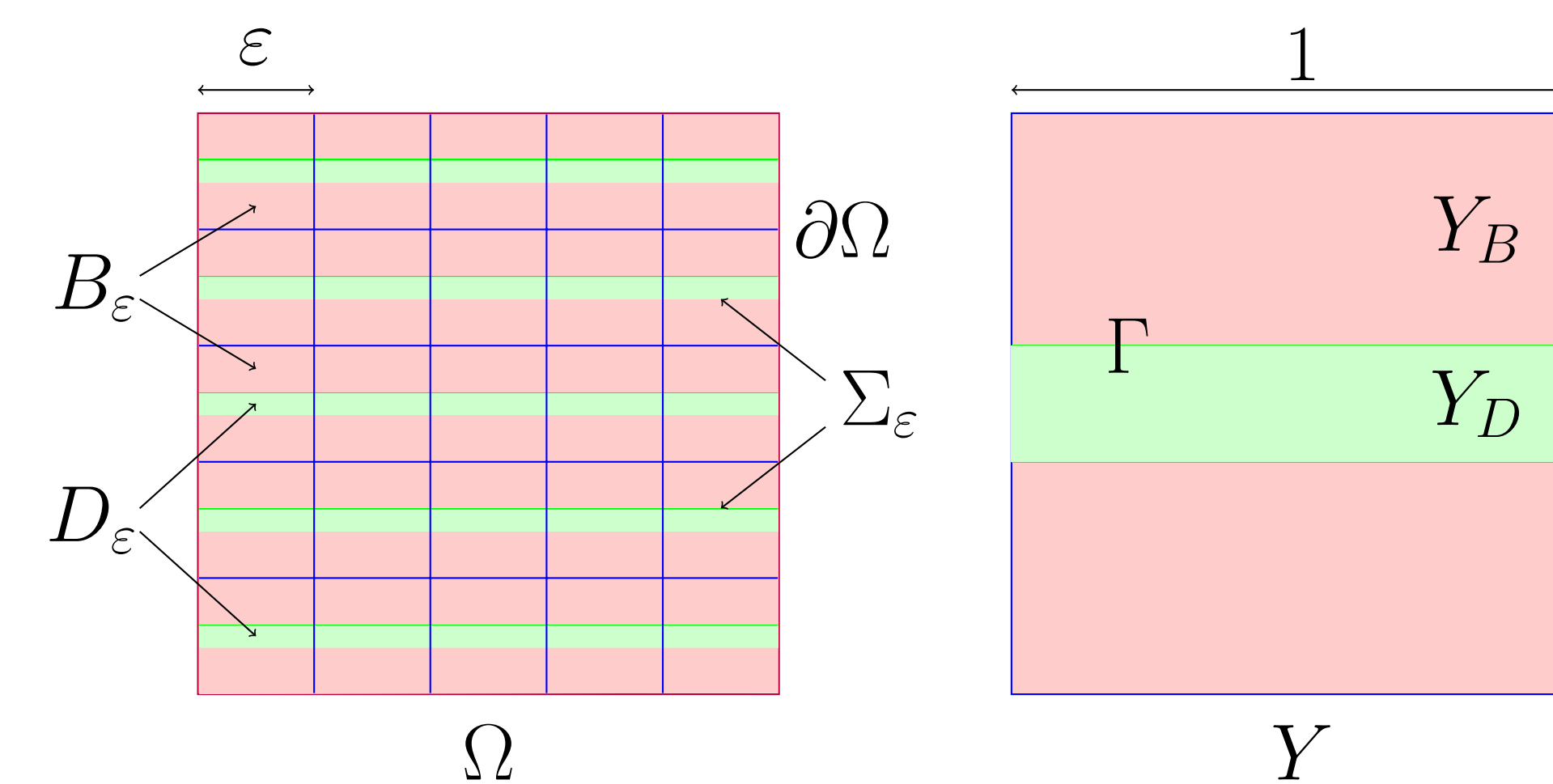
The diffusive inclusion

$$0 = -\nabla \cdot (\sigma_\varepsilon^d \nabla u_\varepsilon^d), \quad \text{in } D_\varepsilon.$$

The transmission conditions

$$\left. \begin{aligned} \sigma_\varepsilon^i \nabla u_\varepsilon^i \cdot \mathbf{n}_{\Sigma_\varepsilon} &= 0, \\ \sigma_\varepsilon^e \nabla u_\varepsilon^e \cdot \mathbf{n}_{\Sigma_\varepsilon} &= \sigma_\varepsilon^d \nabla u_\varepsilon^d \cdot \mathbf{n}_{\Sigma_\varepsilon}, \\ u_\varepsilon^e &= u_\varepsilon^d, \end{aligned} \right\} \quad \text{on } \Sigma_\varepsilon,$$

The initial and the boundary conditions are given.



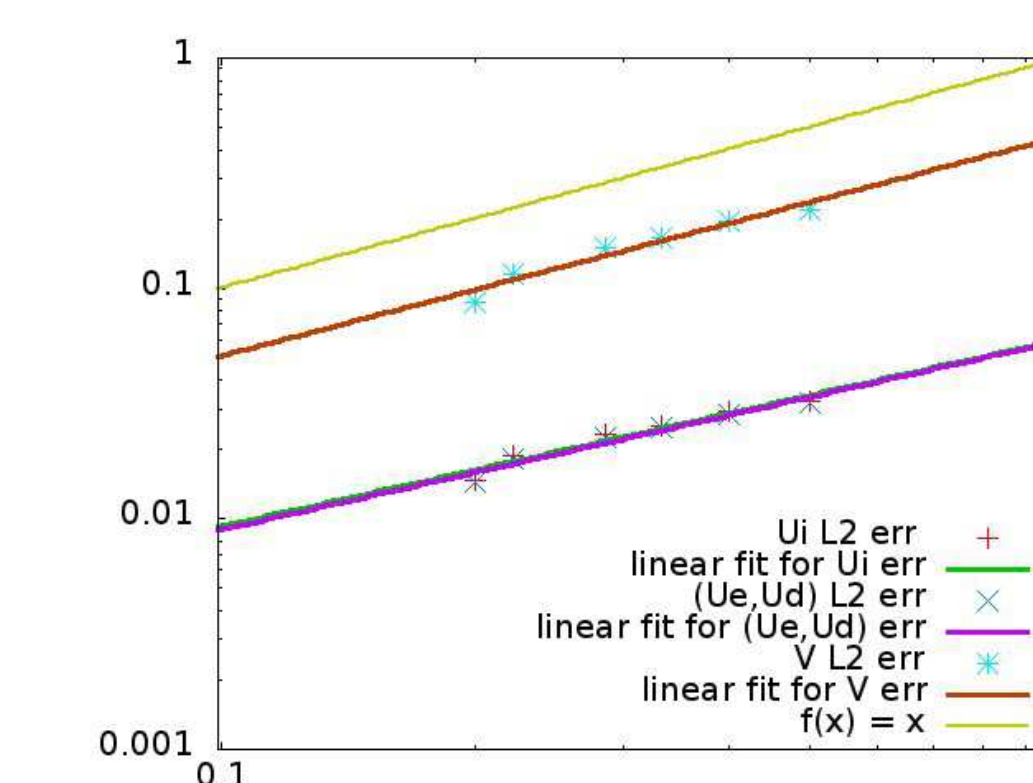
## II - Homogenisation

**Heuristics:** Redefine the problem in two scales and obtain the limit problem that depends only on the large scale variable. It gives a global, macroscopic behaviour of the unknown functions. Assume

$$\sigma_\varepsilon^i(x) = \sigma_i\left(\frac{x}{\varepsilon}\right), \quad \sigma_\varepsilon^e(x) = \sigma_e\left(\frac{x}{\varepsilon}\right), \quad \sigma_\varepsilon^d(x) = \sigma_d\left(\frac{x}{\varepsilon}\right).$$

According to [1], *a priori* estimates ensure the two-scale convergence:

$$u_\varepsilon^i \rightharpoonup u_0^i, \quad (u_\varepsilon^e, u_\varepsilon^d) \rightharpoonup u_0.$$



## III - Macroscopic equations

The limit problem is still a **bidomain model** with updated conductivities:

$$\begin{aligned} \nabla \cdot (\tilde{\sigma}_i \nabla u_0^i) &= \partial_t v_0 + cv_0, & \text{in } \Omega, \\ \nabla \cdot ((\tilde{\sigma}_e + \tilde{\sigma}_d) \nabla u_0) &= -(\partial_t v_0 + cv_0), & \text{in } \Omega, \\ \tilde{\sigma}_i &= \sigma_i + \frac{1}{|Y_B|} A_i, & \tilde{\sigma}_e &= \sigma_e + \frac{1}{|Y_B|} A_e, & \tilde{\sigma}_d &= \frac{|Y_D|}{|Y_B|} \sigma_d + \frac{1}{|Y_B|} A_d, \end{aligned}$$

where  $A_i, A_e$  and  $A_d$  are constant matrices that depend on the geometry of the unit cell.

**Remark:**

- The modified conductivities depend on the volume fraction of the diffusive part and on the geometry of the unit cell.
- If  $|Y_D| = 0$  we obtain the standard bidomain model. If  $|Y_B| = 0$  we get only the diffusion model.
- If  $\sigma_e = \sigma_d$  the modified extracellular conductivity will not depend on the geometry but only on the volume fraction.

## IV - Numerical verification on simple 2D geometries

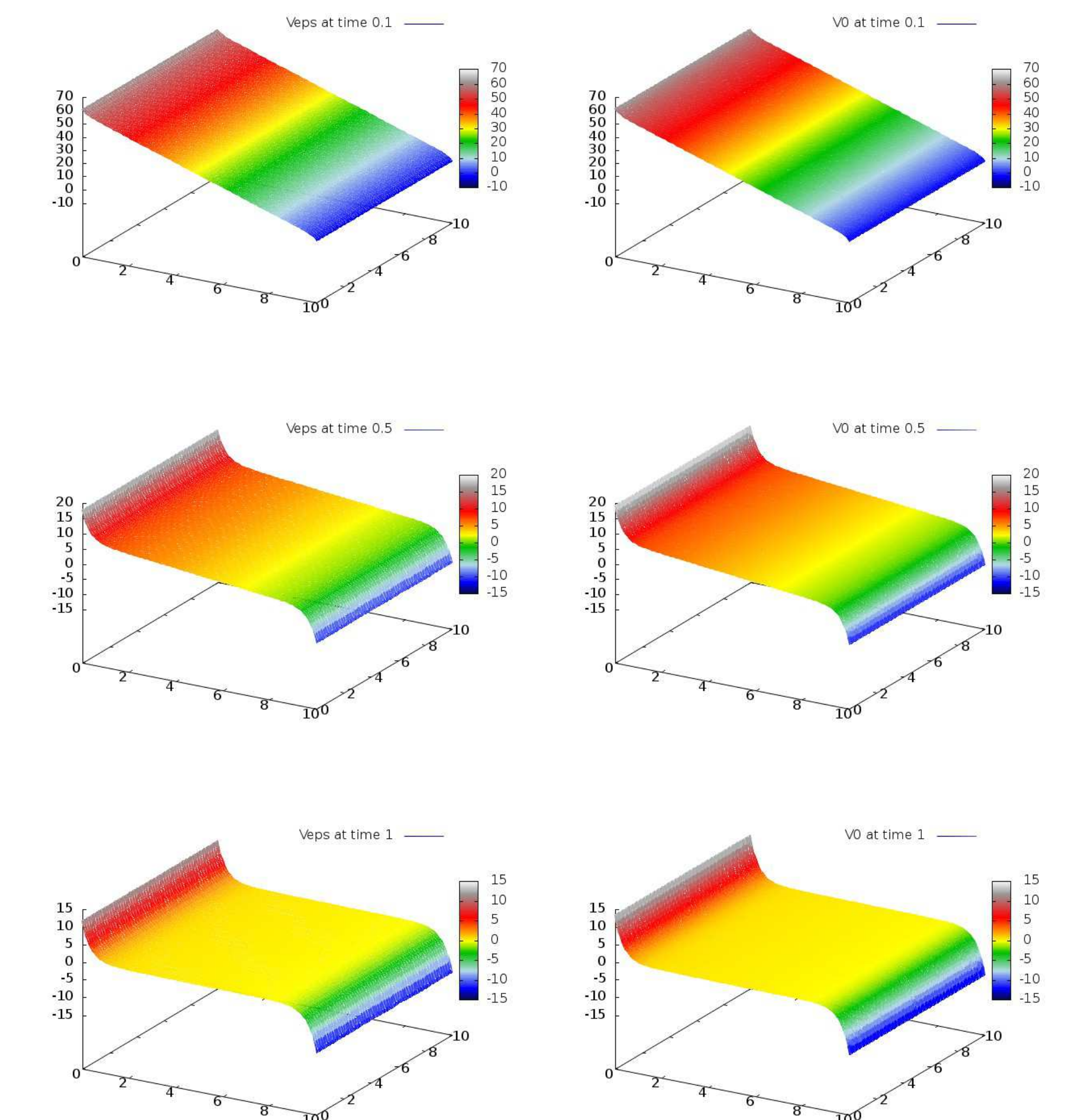


Figure 3: The time evolution for  $v_\varepsilon = u_\varepsilon^i - u_\varepsilon^e, \varepsilon = 0.16$  (left). The time evolution for  $v_0 = u_0^i - u_0$  (right).

## Perspectives

- Apply the ionic current model.
- Use real data, *e.g.* the late enhancement MRI provides the volume fraction of the extracellular space.

## References

- [1] Allaire, *Homogenisation and two-scale convergence*, 1992.
- [2] Hooks *et al.*, *Laminar Arrangement of Ventricular Myocytes Influences Electrical Behavior of the Heart*, 2007.
- [3] Gilbert *et al.*, *Visualization and quantification of whole rat heart laminar structure using high-spatial resolution contrast-enhanced MRI*, 2012.
- [4] Camelliti *et al.*, *Structural and functional characterisation of cardiac fibroblasts*, 2005.