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## To cite this version:

Pierre Gaillard, Mickaël Gastineau. Families of solutions of order nine to the NLS equation with sixteen parameters. 2015. hal-01145780

## HAL Id: hal-01145780 <br> https://hal.science/hal-01145780

Preprint submitted on 26 Apr 2015

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# Families of solutions of order nine to the NLS equation with sixteen parameters. 

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April 26, 2015


#### Abstract

We construct new deformations of the Peregrine breather $\left(P_{9}\right)$ of order 9 with 16 real parameters. With this method, we obtain explicitly new families of quasi-rational solutions to the NLS equation in terms of a product of an exponential depending on $t$ by a ratio of two polynomials of degree 90 in $x$ and $t$; when all the parameters are equal to 0 , we recover the classical $P_{9}$ breather. We construct new patterns of different types of rogue waves as triangular configurations of 45 peaks as well as rings and concentric rings.


PACS: 35Q55, 37K10.
Keywords : NLS equation, Peregrine breather, rogue waves.

## 1 Introduction

After the works of the precursors, Zakharov and Shabat in 1972 [1, 2], a considerable number of studies were carried out by a multiplicity of authors. The first expressions of the quasi-rational solutions were given by Peregrine in 1983 [3]. Akhmediev, Eleonski and Kulagin constructed the first higher order analogue of the Peregrine breather $[4,5]$ in 1986. Akhmediev et al. [6, 7], constructed other families of order 3 and 4, using the Darboux transformations.
It has been found in [10] solutions for the order $N$ (for determinants of order 2 N ) to the NLS equation depending on $2 N-2$ real parameters.
With this method, we construct explicitly here new quasi rational solutions to the NLS equation for the order 9 depending on sixteen parameters. When all the parameters are equal to zero, we recover the famous (analogue) Peregrine breather $P_{9}$. We obtain new deformations at order 9 with 16 real parameters of the $P_{9}$ breather as a ratio of two polynomials of degree 90 in $x$ and $t$ by an exponential depending on $t$. Because of the length of these expressions, we cannot present them here; we only present patterns depending on the choices of the parameters and make the analysis of the evolution of the solutions.
ing to construct deformations of the $P_{9}$ breather, solutions to the NLS equation.

Theorem 2.1 The function $v$ defined by
$v(x, t)=\frac{\operatorname{det}\left(\left(n_{\left.j k)_{j, k \in[1,2 N]}\right)}\right)\right.}{\operatorname{det}\left(\left(d_{j k)_{j, k \in[1,2 N]}}\right)\right.} e^{(2 i t-i \varphi)}$
is a quasi-rational solution of the NLS equation

$$
i v_{t}+v_{x x}+2|v|^{2} v=0
$$

where

$$
\begin{aligned}
& n_{j 1}=f_{j, 1}(x, t, 0), \\
& n_{j k}=\frac{\partial^{2 k-2} f_{j, 1}}{\partial \epsilon^{2 k-2}}(x, t, 0), \\
& n_{j N+1}=f_{j, N+1}(x, t, 0), \\
& n_{j N+k}=\frac{\partial^{2 k-2} f_{j, N+1}}{\partial \epsilon^{2 k-2}}(x, t, 0), \\
& d_{j 1}=g_{j, 1}(x, t, 0), \\
& d_{j k}=\frac{\partial^{2 k-2} g_{j, 1}}{\partial \epsilon^{2 k-2}}(x, t, 0), \\
& d_{j N+1}=g_{j, N+1}(x, t, 0), \\
& d_{j N+k}=\frac{\partial^{2 k-2} g_{j, N+1}}{\partial \epsilon^{2 k-2}}(x, t, 0), \\
& 2 \leq k \leq N, 1 \leq j \leq 2 N
\end{aligned}
$$

The functions $f$ and $g$ are defined for $1 \leq k \leq N$ by :

$$
\begin{aligned}
& f_{4 j+1, k}=\gamma_{k}^{4 j-1} \sin A_{k}, \\
& f_{4 j+2, k}=\gamma_{k}^{4 j} \cos A_{k}, \\
& f_{4 j+3, k}=-\gamma_{k}^{4 j+1} \sin A_{k}, \\
& f_{4 j+4, k}=-\gamma_{k}^{4 j+2} \cos A_{k}, \\
& f_{4 j+1, N+k}=\gamma_{k}^{2 N-4 j-2} \cos A_{N+k}, \\
& f_{4 j+2, N+k}=-\gamma_{k}^{2 N-4 j-3} \sin A_{N+k}, \\
& f_{4 j+3, N+k}=-\gamma_{k}^{2 N-4 j-4} \cos A_{N+k}, \\
& f_{4 j+4, k}=\gamma_{k}^{2 N-4 j-5} \sin A_{N+k}, \\
& g_{4 j+1, k}=\gamma_{k}^{4 j-1} \sin B_{k}, \\
& -g_{4 j+2, k}=\gamma_{k}^{4 j} \cos B_{k}, \\
& g_{4 j+3, k}=-\gamma_{k}^{4 j+1} \sin B_{k}, \\
& g_{4 j+4, k}=-\gamma_{k}^{4 j+2} \cos B_{k}, \\
& g_{4 j+1, N+k}=\gamma_{k}^{2 N-4 j-2} \cos B_{N+k}, \\
& g_{4 j+2, N+k}=-\gamma_{k}^{2 N-4 j-3} \sin B_{N+k}, \\
& g_{4 j+3, N+k}=-\gamma_{k}^{2 N-4 j-4} \cos B_{N+k}, \\
& g_{4 j+4, N+k}=\gamma_{k}^{2 N-4 j-5} \sin B_{N+k},
\end{aligned}
$$

We recall briefly the results obtained in $[11,12]$ that we use in the follow-

## 2 Determinant represen tation of solutions to NLS equation

The arguments $A_{\nu}$ and $B_{\nu}$ of these functions are given for $1 \leq \nu \leq 2 N$ by

$$
\begin{aligned}
& A_{\nu}=\kappa_{\nu} x / 2+i \delta_{\nu} t-i x_{3, \nu} / 2-i e_{\nu} / 2 \\
& B_{\nu}=\kappa_{\nu} x / 2+i \delta_{\nu} t-i x_{1, \nu} / 2-i e_{\nu} / 2
\end{aligned}
$$

The terms $\kappa_{\nu}, \delta_{\nu}, \gamma_{\nu}$ are defined by $1 \leq \nu \leq 2 N$

$$
\begin{align*}
& \kappa_{j}=2 \sqrt{1-\lambda_{j}^{2}}, \delta_{j}=\kappa_{j} \lambda_{j} \\
& \gamma_{j}=\sqrt{\frac{1-\lambda_{j}}{1+\lambda_{j}}}, \kappa_{N+j}=\kappa_{j}  \tag{2}\\
& \delta_{N+j}=-\delta_{j}, \gamma_{N+j}=1 / \gamma_{j} \\
& 1 \leq j \leq N
\end{align*}
$$

where $\lambda_{j}$ are given for $1 \leq j \leq N$ by :

$$
\begin{equation*}
\lambda_{j}=1-2 j^{2} \epsilon^{2}, \quad \lambda_{N+j}=-\lambda_{j} \tag{3}
\end{equation*}
$$

The terms $x_{r, \nu}(r=3,1)$ are defined for $1 \leq \nu \leq 2 N$ by :

$$
\begin{equation*}
x_{r, \nu}=(r-1) \ln \frac{\gamma_{\nu}-i}{\gamma_{\nu}+i} \tag{4}
\end{equation*}
$$

The parameters $e_{\nu}$ are given by

$$
\begin{align*}
& e_{j}=i \sum_{k=1}^{N-1} \tilde{a}_{j} \epsilon^{2 k+1} j^{2 k+1} \\
& -\sum_{k=1}^{N-1} \tilde{b}_{j} \epsilon^{2 k+1} j^{2 k+1} \\
& e_{N+j}=i \sum_{k=1}^{N-1} \tilde{a}_{j} \epsilon^{2 k+1} j^{2 k+1}  \tag{5}\\
& +\sum_{k=1}^{N-1} \tilde{b}_{j} \epsilon^{2 k+1} j^{2 k+1} \\
& 1 \leq j \leq N,
\end{align*}
$$

## 3 Deformations of the $P_{9}$ breather with sixteen parameters

We have constructed in $[17,16,19,18$, $20,22,14]$ and in a series of articles on the archive hal, solutions for the cases from $N=1$ until $N=8$ with $2 N-2$ parameters.
Here we make the study for the order nine. We don't give the analytic expression of the solution of NLS equation of order 9 with sixteen parameters
because of the length of the expression. For simplicity, we denote

$$
\begin{aligned}
d_{3} & :=\operatorname{det}\left(\left(n_{j k)_{j, k \in[1,2 N]}}\right)\right. \\
d_{1} & :=\operatorname{det}\left(\left(d_{j k}\right)_{j, k \in[1,2 N]}\right) .
\end{aligned}
$$

The number of terms of the polynomials of the numerator $d 3$ and denominator $d 1$ of the solutions are shown in the table below (Table 1) when other $a_{i}$ and $b_{i}$ are set to 0 .

| $\mathrm{N}=9$ | number of terms |
| :---: | :---: |
| $d_{3}\left(a_{1}, b_{1}, x, t\right)$ | 184554 |
| $d_{1}\left(a_{1}, b_{1}, x, t\right)$ | 94332 |
| $d_{3}\left(a_{2}, b_{2}, x, t\right)$ | 72174 |
| $d_{1}\left(a_{2}, b_{2}, x, t\right)$ | 36894 |
| $d_{3}\left(a_{3}, b_{3}, x, t\right)$ | 39813 |
| $d_{1}\left(a_{3}, b_{3}, x, t\right)$ | 20347 |

Table 1: Number of terms for the polynomials $d_{3}$ and $d_{1}$ of the solutions of the NLS equation.

We construct figures to show deformations of the ninth Peregrine breather. The computations were done using the computer algebra systems Maple and TRIP [26]. For example, the computations of the determinants of $Q(x, t)$ take only 3 hours for the figure 1 using the parallel kernel of the software TRIP on a workstation with 6 cores.
We get different types of symmetries in the plots in the $(x, t)$ plane. We give some examples of this fact in the following.

### 3.1 Peregrine breather of order 9

If we choose $\tilde{a}_{i}=\tilde{b}_{i}=0$ for $1 \leq i \leq 8$, we obtain the classical Peregrine breather


Figure 1: Solution to NLS, $\mathrm{N}=9$, all parameters equal to 0 , Peregrine breather $P_{9}$.

In the following, we give different deformations of this $P_{9}$ breather depending on the choices of the parameters. The triangular patterns have already been explained for the orders until $N=7$ in [27]. The ring patterns, and the classification of the solutions to the NLS equation were already given in [28] using numerical methods until order $N=6$. In that paper, the cases of order $7,8,9$ and 10 given in table page 9 have been extrapolated. It was also pointed out that the number of peaks in the different figures for $N=9$ is $N(N+1) / 2=45$.
The results here are obtained from explicitly exact solutions.

### 3.2 Variation of parameters

With other choices of parameters, we obtain all types of configurations : triangles and multiple concentric rings configurations with a maximum of 45 peaks.


Figure 2: Solution to NLS, $N=9, \tilde{a}_{1}=$ $10^{3}$ : triangle with 45 peaks; in bottom, sight of top.


Figure 3: Solution to NLS, $N=9, \tilde{b}_{1}=$ $10^{3}$ : triangle with 45 peaks; in bottom, sight of top.


Figure 4: Solution to NLS, $\mathrm{N}=9, \tilde{a}_{2}=$ $10^{5}: 6$ rings without a peak in the center; in bottom, sight of top.


Figure 5: Solution to NLS, $\mathrm{N}=9, \tilde{b}_{2}=$ $10^{5}$ : 6 rings without a peak in the center; in bottom, sight of top.

Figure 6: Solution to NLS, $N=9, \tilde{a}_{3}=$ $10^{7}: 4$ rings with in the center $P_{2}$; in bottom, sight of top.


Figure 7: Solution to NLS, $N=9, \tilde{b}_{3}=$ $10^{7}: 4$ rings with in the center $P_{2}$; in bottom, sight of top.

Similar figures were given in [28], and in all these figures, we have $N(N+$ 1) $/ 2=45$ peaks as it was already extrapolated in the table page 9 of that article. In the same paper, it was al-


Figure 8: Solution to NLS, $N=9, \tilde{a}_{4}=$ $10^{9}: 5$ rings without a central peak; in bottom, sight of top.


Figure 9: Solution to NLS, $N=9, \tilde{b}_{4}=$ $10^{9}: 5$ rings without a central peak; in bottom, sight of top.

Figure 11: Solution to NLS, $N=9, \tilde{b}_{5}=$ $10^{15}$ : 4 rings with in the center one peak; in bottom, sight of top.
peaks surrounding the $P_{N-2}=P_{7}$ breather. It can be noted that when a more convenient normalization of the NLS equation is used, these rings can be transformed in circles.


Figure 12: Solution to NLS, $N=9, \tilde{a}_{6}=$ $10^{15}$ : 3 rings with in the center the Peregrine breather of order 3; in bottom, sight of top.



Figure 13: Solution to NLS, $\mathrm{N}=9, \tilde{b}_{6}=$ $10^{15}$ : 3 rings with in the center the Peregrine breather of order 3 ; in bottom, sight of top.

## 4 Conclusion

In the present paper we have constructed explicitly solutions to the NLS equa-


Figure 14: Solution to NLS, $N=9, \tilde{a}_{7}=$ $10^{20}$ : two rings with in the center the Peregrine breather of order 5; in bottom, sight of top.


Figure 15: Solution to NLS, $\mathrm{N}=9, \tilde{b}_{7}=$ $10^{20}$ : two rings with in the center the Peregrine breather of order 5; in bottom, sight of top.
tion of order 9 with 16 real parameters. The explicit representation in terms of polynomials in $x$ and $t$ is too large to


Figure 16: Solution to NLS, $N=9, \tilde{a}_{8}=$ $10^{20}$ : one ring with in the center the Peregrine breather of order 7; in bottom, sight of top.


Figure 17: Solution to NLS, $N=9, \tilde{b}_{8}=$ $10^{20}$ : one ring with in the center the Peregrine breather of order 7 ; in bottom, sight of top.
be published here.
In the case of the variation of one parameter, we obtain different types of
configurations with a maximum of 45 peaks.
By different choices of these parameters, we obtained new patterns in the $(x ; t)$ plane; we recognized ring shape as already observed in the case of deformations depending on two parameters $[12,10]$. We get new triangular shapes and multiple concentric rings. Many applications in nonlinear optics and hydrodynamics have been given recently : we can mention in particular the works of Chabchoub et al. [29] or Kibler et al. [30].
These explicit solutions of order 9 and their deformations with eighteen parameters are presented for the first time to our knowledge. These deformations of the ninth Peregrine breather give a better understanding of the phenomenon of appearance of rogue waves and their asymptotic behavior.
The study of the solutions to the NLS equation has been done until order $N=$ 6 by Akhmediev et al. in [28] and extrapolated until order $N=10$. Present work gives explicitly exact solutions to the NLS equation at order 9; it verifies the conjectured classification given in [28].
It would be important to continue this study to try to classify solutions to the NLS equation in the general case of or$\operatorname{der} N(N>10)$.

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