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INVERSE POTENTIAL PROBLEMS FOR DIVERGENCE OF MEASURES WITH TOTAL VARIATION REGULARIZATION

L. BARATCHART, C. VILLALOBOS GUILLEN, D.P. HARDIN, M.C. NORTHINGTON,
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ABSTRACT. We study inverse problems for the Poisson equation with source term the divergence of an \mathbb{R}^3 -valued measure; that is, the potential Φ satisfies

$$\Delta\Phi = \operatorname{div} \boldsymbol{\mu},$$

and $\boldsymbol{\mu}$ is to be reconstructed knowing (a component of) the field $\operatorname{grad} \Phi$ on a set disjoint from the support of $\boldsymbol{\mu}$. Such problems arise in several electro-magnetic contexts in the quasi-static regime, for instance when recovering a remanent magnetization from measurements of its magnetic field. We develop methods for recovering $\boldsymbol{\mu}$ based on total variation regularization. We provide sufficient conditions for the unique recovery of $\boldsymbol{\mu}$, asymptotically when the regularization parameter and the noise tend to zero in a combined fashion, when it is uni-directional or when the magnetization has a support which is sparse in the sense that it is purely 1-unrectifiable.

Numerical examples are provided to illustrate the main theoretical results.

1. INTRODUCTION

intro

This work is concerned with inverse potential problems with source term in divergence form. That is, a \mathbb{R}^3 -valued measure on \mathbb{R}^3 has to be recovered knowing (one component of) the field of the Newton potential of its divergence on a piece of surface, away from the support. Such issues typically arise in source identification from field measurements for Maxwell's equations, in the quasi-static regime. They occur for instance in electroencephalography (EEG), magneto-encephalography (MEG), geomagnetism and paleomagnetism, as well as in several non-destructive testing problems, see *e.g.* [3, 4, 12, 36, 37] and their bibliographies. A model problem of our particular interest is inverse scanning magnetic microscopy, as considered for instance in [7, 34, 5] to recover magnetization distributions of thin rock samples, but the considerations below are of a more general and abstract nature. Our main objective is to introduce notions of sparsity that help recovery in this infinite-dimensional context, when regularization is performed by penalizing the total variation of the measure.

Key words and phrases. divergence free, distributions, solenoidal, total variation of measures, magnetization, inverse problems, purely 1-unrectifiable.

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1.1. **Two Basic Extremal Problems.** For a closed subset $S \subset \mathbb{R}^3$, let $\mathcal{M}(S)$ denote the space of finite signed Borel measures on \mathbb{R}^3 whose support lies in S . In this paper we consider inverse problems associated with the equation

equationPhi

$$(1) \quad \Delta \Phi = \operatorname{div} \boldsymbol{\mu}$$

where $\boldsymbol{\mu}$ is an unknown measure in $\mathcal{M}(S)^3$ to be recovered. Under suitable conditions on the decay of Φ at infinity, there is a unique solution $\Phi = \Phi(\boldsymbol{\mu})$ (see Section 2). Such problems arise in magnetostatics (our primary motivation) where $\boldsymbol{\mu}$ models a *magnetization* distribution. Then

bextension

$$(2) \quad \mathbf{b}(\boldsymbol{\mu}) = \mu_0 (\boldsymbol{\mu} - \operatorname{grad} \Phi(\boldsymbol{\mu}))$$

is the *magnetic field* $\mathbf{b}(\boldsymbol{\mu})$ generated by $\boldsymbol{\mu}$ and it follows from (1) that $\mathbf{b}(\boldsymbol{\mu})$ is divergence-free. The term $\mathbf{h}(\boldsymbol{\mu}) := \mu_0 \operatorname{grad} \Phi(\boldsymbol{\mu})$ is called the *magnetic intensity* generated by $\boldsymbol{\mu}$. We refer to (1) as a *Poisson-Hodge* equation since (2) provides a decomposition of $\boldsymbol{\mu}$ into curl-free and divergence-free terms.

The mapping $\boldsymbol{\mu} \rightarrow \mathbf{b}(\boldsymbol{\mu})$ is, in general, not injective. We say that magnetizations $\boldsymbol{\mu}, \boldsymbol{\nu} \in \mathcal{M}(S)^3$ are *S-equivalent* if $\mathbf{b}(\boldsymbol{\mu})$ and $\mathbf{b}(\boldsymbol{\nu})$ agree on $\mathbb{R}^3 \setminus S$, in which case we write

$$\boldsymbol{\mu} \stackrel{S}{\equiv} \boldsymbol{\nu}.$$

A magnetization $\boldsymbol{\mu}$ is said to be *S-silent* (or *silent in* $\mathbb{R}^3 \setminus S$) if it is *S-equivalent* to the zero magnetization; i.e., if $\mathbf{b}(\boldsymbol{\mu})$ vanishes on $\mathbb{R}^3 \setminus S$. It is suggestive from (1) (see Theorem 2.2 below) that a divergence-free magnetization $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is *S-silent*. For the converse direction we introduce the following terminology.

defslender

Definition 1. We will call a closed set $S \subset \mathbb{R}^3$ *slender* (with respect to \mathbb{R}^3) if its Lebesgue measure $\mathcal{L}_3(S) = 0$ and each connected component C of $\mathbb{R}^3 \setminus S$ has $\mathcal{L}_3(C) = \infty$.

For example, a subset of a plane is slender in \mathbb{R}^3 , while a sphere or a ball is not.

Slender sets form a family of S for which the above-mentioned inverse problem makes contact with geometric measure theory, since Theorem 2.2 shows that if S is a slender set, then any *S-silent* magnetization is divergence-free. Smirnov [40] characterizes divergence-free magnetizations (also known as *solenoids*) in \mathbb{R}^3 in terms of magnetizations that are absolutely continuous with respect to 1-dimensional Hausdorff measure, and we shall make use of this characterization.

We will assume that scalar data of the form $f = A(\boldsymbol{\mu}) := v \cdot \mathbf{b}(\boldsymbol{\mu})$, for some fixed nonzero vector $v \in \mathbb{R}^3$, is given on a closed subset $Q \subset \mathbb{R}^3 \setminus S$, where A is the *forward operator* mapping $\boldsymbol{\mu}$ to the restriction on Q of $v \cdot \mathbf{b}(\boldsymbol{\mu})$. We will consider the situation where

- (a) $A : \mathcal{M}(S)^3 \rightarrow L^2(Q)$ boundedly and
- (b) $A(\boldsymbol{\mu}) = 0$ if and only if $\boldsymbol{\mu}$ is *S-silent*.

Since $\mathbf{b}(\boldsymbol{\mu})$ is harmonic on $\mathbb{R}^3 \setminus S$, condition (a) will hold if Q and S are positively separated, and if Q is compact (which is the case in practice), then A is a compact operator. Theorem 2.3 provides sufficient conditions for (b) to hold. Condition (b) means that the observation is “faithful”, i.e. if the v -component of the field on Q is zero then the field is indeed zero everywhere off S . In this case, the null space of A (which is a crucial ingredient

of the inverse problem) coincides with S -silent magnetizations which depend solely on the geometry of S and can be studied using potential and measure-theoretic tools. As we shall see, A has, in general, a nontrivial null space and, if Q is “thin enough”, it has dense range (see Lemma 3.1). Note that a typical magnetic sensor is a coil measuring the component of the field parallel to its axis, which is why measurements are assumed to be of the form $v \cdot \mathbf{b}(\boldsymbol{\mu})$. The fact that v is a constant vector means that the orientation of the sensor is kept fixed. In some cases, *e.g.* in magneto-encephalography, v would depend on the point where a measurement is made. We do not consider this more complicated situation, as we are particularly motivated by applications to scanning magnetic microscopy (SMM) where measurements of the vertical component of the magnetic field, namely $b_3(\boldsymbol{\mu}) = e_3 \cdot \mathbf{b}(\boldsymbol{\mu})$ (e_i denotes the i -th unit vector of the canonical basis for \mathbb{R}^3 , for $i = 1, 2, 3$), are taken on a rectangle Q in a plane $x_3 = h$ for some $h > 0$, while S is contained in the half-space $x_3 \leq 0$.

The focus of this paper is on the use of the total variation norm for measures to regularize the ill-posed inverse problem of recovering $\boldsymbol{\mu}$ from $A(\boldsymbol{\mu})$. Let us briefly recall that the total variation measure of $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is the positive Borel measure $|\boldsymbol{\mu}| \in \mathcal{M}(S)$ such that $d\boldsymbol{\mu} = \mathbf{u}_\mu d|\boldsymbol{\mu}|$ for some \mathbb{R}^3 -valued Borel measurable function \mathbf{u}_μ satisfying $|\mathbf{u}_\mu| = 1$ at $|\boldsymbol{\mu}|$ -a.e. point, see (8). The *total variation norm* of $\boldsymbol{\mu}$ is then defined as

$$\boxed{\text{TVnormdef}} \quad (3) \quad \|\boldsymbol{\mu}\|_{TV} := |\boldsymbol{\mu}|(\mathbb{R}^3).$$

If \mathbf{u}_μ is constant $|\boldsymbol{\mu}|$ -a.e., then we call $\boldsymbol{\mu}$ *uni-directional*. Such magnetizations arise in geological samples whose remanent magnetizations were formed in a uniform external field.

We investigate two extremal problems for magnetization recovery. The first is that of minimizing the total variation over magnetizations S -equivalent to a given one. To fix notation, for $\boldsymbol{\mu} \in \mathcal{M}(S)^3$, let

$$\boxed{\text{defmineq}} \quad (4) \quad M(\boldsymbol{\mu}) := \inf\{\|\boldsymbol{\nu}\|_{TV} : \boldsymbol{\nu} \stackrel{S}{\equiv} \boldsymbol{\mu}\}.$$

Extremal Problem 1 (EP-1). For $\boldsymbol{\mu}_0 \in \mathcal{M}(S)^3$, find $\boldsymbol{\mu} \stackrel{S}{\equiv} \boldsymbol{\mu}_0$ such that $\|\boldsymbol{\mu}\|_{TV} = M(\boldsymbol{\mu}_0)$.

The second extremal problem involves minimizing the following functional defined for $\boldsymbol{\mu} \in \mathcal{M}(S)^3$, $f \in L^2(Q)$, and $\lambda > 0$, by

$$\boxed{\text{defcrit0}} \quad (5) \quad \mathcal{F}_{f,\lambda}(\boldsymbol{\mu}) := \|f - A\boldsymbol{\mu}\|_{L^2(Q)}^2 + \lambda\|\boldsymbol{\mu}\|_{TV}.$$

We consider the problem of finding some $\boldsymbol{\mu}_\lambda \in \mathcal{M}(S)^3$ minimizing $\mathcal{F}_{f,\lambda}$:

Extremal Problem 2 (EP-2). Given $f \in L^2(Q)$, find $\boldsymbol{\mu}_\lambda \in \mathcal{M}(S)^3$ such that

$$\boxed{\text{crit0}} \quad (6) \quad \mathcal{F}_{f,\lambda}(\boldsymbol{\mu}_\lambda) = \inf_{\boldsymbol{\mu} \in \mathcal{M}(S)^3} \mathcal{F}_{f,\lambda}(\boldsymbol{\mu}).$$

We remark that the total variation norm is convex on $\mathcal{M}(S)^3$ but not strictly convex and so there may be multiple $\boldsymbol{\mu}$ that solve EP-1 for a given $\boldsymbol{\mu}_0$. Still, we show in Section 2 that, under the assumption that S is slender, EP-1 uniquely recovers the magnetization

in two important cases: (a) the magnetization is supported on a purely 1-unrectifiable set (see Theorem 2.6) and (b) the magnetization is piecewise uni-directional (see Theorem 2.11). The notion of purely 1-unrectifiable set is classical from geometric measure theory, and means that intersection of the set with any rectifiable arc has 1-dimensional Hausdorff measure zero. In particular, any set with 1-dimensional Hausdorff measure zero is purely 1-unrectifiable and so EP-1 recovers a large class of magnetizations with ‘sparse’ support including, for example, those with countable or Cantor-like (with Hausdorff dimension less than 1) support. Note that uni-directionality involves some sort of sparsity as well, this time direction-wise.

In Section 2, we also consider the *net moment* of $\boldsymbol{\mu}$ given by

$$\boxed{\text{netmomdef}} \quad (7) \quad \langle \boldsymbol{\mu} \rangle := \boldsymbol{\mu}(\mathbb{R}^3),$$

which is a quantity of physical interest and is useful in our analysis of uni-directional magnetizations. Under the assumption that S is compact or a slender set, Theorem 2.2 and Lemma 2.8 show that S -equivalent magnetizations must have the same net moment. Hence, $\langle \boldsymbol{\mu} \rangle$ is uniquely defined by the field in this case. However, Remark 2.4 after that proposition shows this needs not hold if S is neither compact nor slender.

Solutions to EP-2 connect to those of EP-1 as follows. If $f = A(\boldsymbol{\mu}) + e$ with ‘noise’ $e \in L^2(Q)$, any weak-* accumulation point of the solutions to EP-2 when $\lambda \rightarrow 0$ and $e/\sqrt{\lambda} \rightarrow 0$ must be a solution of EP-1. This Tikhonov-like regularization theory is by now essentially understood in a more general context [14, 31, 13]. In Section 4, we improve on previous results by showing that this property holds not only for the \mathbb{R}^3 -valued measures involved but also for their total variation measures, and that weak-* convergence strengthens to narrow convergence (see Theorem 4.3). Hence, “no mass is lost” in the limit. This is important for if EP-1 has a unique solution, it implies that solutions to EP-2 asymptotically put mass exactly where the “true” magnetization does, a property which is usually not satisfied under weak-* convergence of measures. Another feature of solutions to EP-2, which is more specific to the present situation, is that they are supported on “small” sets, of codimension at least 1 in S (see Corollary 4.2).

Altogether, when the “true” magnetization is sparse in one of the senses mentioned above (that is, if it can be recovered by EP-1), we obtain in Theorem 4.4 and Corollary 4.5 asymptotic recovery results which, from the strict point of view of inverse problems, recap the main contributions of the paper. Finally, let us mention that EP-2 has a unique solution when S is contained in a plane, even though $\mathcal{F}_{f,\lambda}$ is not strictly convex if A has a nontrivial null space (the usual case considered in this paper). A proof would take us too far into the structure of divergence-free measures in the plane, and is beyond the scope of this paper, see [8]. Whether uniqueness holds more generally is an intriguing open question.

1.2. Background and Related Work. The class of inverse problems investigated below can be viewed as deconvolution from partial observations. Note that we work with vector-valued measures supported on sets of dimension greater than 1, and the forward operator A , whose convolution kernel is (weakly) singular, has a rather large null-space which depends

in a fairly complicated manner on the geometry of the observation set Q and of the set S *a priori* containing the support of the unknown measure.

The main contribution of the present work is perhaps to connect geometric measure theory with regularization theory for such inverse problems. This connection essentially rests on the structure of the null-space of the forward operator: in fact, the conditions on Q and S set up in this work for unique recovery of sparse measures to hold are designed so that this null-space consists exactly of divergence-free measures. These conditions generalize those given in [7] for planar distributions to magnetizations supported on more general sets in \mathbb{R}^3 . This characterization of the null-space is central to the present approach, and allows us to rely on classical tools from geometric measure theory and on material from [40] to proceed with the proof of our main result about EP-1, namely the minimality of the total variation of a sparse measure in its coset modulo the null-space. More general situations, notably the case where S is a closed surface or where it has positive Lebesgue measure in \mathbb{R}^3 (and thus is not slender), are left for future research. Such situations are relevant in particular to EEG and MEG inverse problems, as well as to SMM on volumic samples.

Whether the property of having purely 1-unrectifiable support qualifies a measure as being “sparse” is debatable: for instance the support could still disconnect the space (like the Koch curve does in 2-D). Nevertheless, it comprises standard notions of sparsity, such as being a finite sum of Dirac masses, which is why we consider purely 1-unrectifiability of the support as a generalized notion of sparsity in this context. Moreover if S is slender we shall see that another convenient notion of sparsity is to be a finite sum of measures with disjoint supports, each of which assumes constant direction (see Theorem 2.11). Thus, sparsity in the present context may also be envisaged “direction-wise”.

Let us now briefly discuss connections to other works dealing with sparse recovery. After early studies [22, 42] and the seminal work in [16, 17, 23, 24], approximately solving under-determined systems of linear equations in \mathbb{R}^n by minimizing the residuals while penalizing the l^1 -norm has proved to be quite successful in identification. In fact, under appropriate assumptions on the matrix of the system, this kind of approximation favors the recovery of sparse solutions, *i.e.* solutions having a large number of zero components, when they exist. This has resulted in the theory of compressed sensing, which shows by and large that a sparse signal can be recovered from much less linear observations than is *a priori* needed, see for example [28] and the bibliography therein.

In recent years, natural analogs in infinite-dimensional settings have been investigated by several authors, but then the situation is much less understood. A Tikhonov-like regularization theory was developed in [14, 31, 13] for linear equations whose unknown is a (possibly \mathbb{R}^n -valued) measure, by minimizing the residuals while penalizing the total variation. As expected from the non-reflexive character of spaces of measures, consistency estimates generally hold in a rather weak sense, such as weak-* convergence of subsequences to solutions of minimum total variation, or convergence in the Bregman distance when the so-called source condition holds. Algorithms and proofs typically rely on Fenchel duality, and reference [13] contains an extension of the soft thresholding algorithm to the case where the unknown gets parametrized as a finite linear combination of Dirac masses,

whose location no longer lies on a fixed grid in contrast with the discrete case. References [19, 18], which deal with inverse source problems for elliptic operators, dwell on the same circle of ideas but suggest a different thresholding method, connected with a Newton step, or else a finite element discretization of the equation having a linear combination of Dirac masses amongst its solutions. These methods yield constructive algorithms to approximate a solution of minimum total variation to the initial equation by a sequence of discrete measures, which is always possible in theory since these are weak-* dense in the space of measures supported on an open subset of \mathbb{R}^n . To obtain an asymptotic recovery result, in the weak-* sense as the regularizing parameter goes to zero, it remains to identify conditions on a measure ensuring that it is the unique element of least total variation in its coset modulo the null-space of the forward operator. The gist of compressed sensing is that, in finite-dimension, sparsity is such a condition for “most” operators. In infinite-dimension, when total variation-penalization replaces l^1 -penalization, it is unclear which assumptions imply the desired uniqueness and why they should be connected with sparsity. Moreover, granted the diversity of operators involved in applications, it is likely that such assumptions will much depend on the situation under examination.

Still, a striking connection with sparsity, in the sense of being a sum of Dirac masses, was recently established for 1-D deconvolution issues, where a train of spikes is to be recovered from filtered observation thereof [11, 19, 21, 15, 25]. More precisely, following the work in [15] which identifies a sum of Dirac masses as the solution of minimum total variation in the case of an ideal low-pass filter, it was shown in [25] that a finite train of spikes can be recovered arbitrarily well by total variation regularization, when the noise and the regularization parameter are small enough, provided that the filter satisfies nondegeneracy conditions and the spikes are sufficiently separated. Also, the algorithm in [13] can be used for that purpose. The present work may be viewed as investigating in higher dimension such deconvolution issues, the filter being now defined by the Biot-Savart law, and Corollary 4.5 may be compared to [25, Thm. 2].

1.3. Notation. We conclude this section with some details of our notation, as used above, for vectors and vector-valued functions, and with some preliminaries concerning measures and distributions. For a vector x in the Euclidean space \mathbb{R}^3 , we denote the j -th component of x by x_j and the partial derivative with respect to x_j by ∂_{x_j} . For $x, y \in \mathbb{R}^3$, $x \cdot y$ and $|x| = \sqrt{x \cdot x}$ denote the usual Euclidean scalar product and norm, respectively. We denote the gradient with respect to y by grad_y and drop the subscript if the variable with respect to which we differentiate is unambiguous. By default, we consider vectors x as column vectors; e.g., for $x \in \mathbb{R}^3$ we write $x = (x_1, x_2, x_3)^T$ where “ T ” denotes “transpose”. We use bold symbols to represent vector-valued functions and measures, and the corresponding nonbold symbols with subscripts to denote the respective components; e.g., $\boldsymbol{\mu} = (\mu_1, \mu_2, \mu_3)$ or $\mathbf{b}(\boldsymbol{\mu}) = (b_1(\boldsymbol{\mu}), b_2(\boldsymbol{\mu}), b_3(\boldsymbol{\mu}))$. We let δ_x stand for the Dirac delta measure at a point $x \in \mathbb{R}^3$ and refer to a magnetization of the form $\boldsymbol{\mu} = v\delta_x$ for some $v \in \mathbb{R}^3$ as the *point dipole at x with moment v* .

For $x \in \mathbb{R}^3$ and $R > 0$, we let $\mathbb{B}(x, R)$ denote the open ball centered at x with radius R and $\mathbb{S}(x, R)$ the boundary sphere. Given a finite measure $\mu \in \mathcal{M}(\mathbb{R}^3)$ and a Borel set

$E \subset \mathbb{R}^3$, we denote by $\mu|_E$ the finite measure obtained by restricting μ to E (i.e. for every Borel set $B \subset \mathbb{R}^3$, $\mu|_E(B) := \mu(E \cap B)$).

For $\mu \in \mathcal{M}(S)^3$, the *total variation measure* $|\mu|$ is defined on Borel sets $B \subset \mathbb{R}^3$ by

totVarMeasDef

$$(8) \quad |\mu|(B) := \sup_{\mathcal{P}} \sum_{P \in \mathcal{P}} |\mu(P)|,$$

where the supremum is taken over all finite Borel partitions \mathcal{P} of B . Since $|\mu|$ is a Radon measure, the Radon-Nikodym derivative $\mathbf{u}_\mu := d\mu/d|\mu|$ exists and satisfies $|\mathbf{u}_\mu| = 1$ a.e. with respect to $|\mu|$.

We shall identify $\mu \in \mathcal{M}(\mathbb{R}^k)^k$ with the linear form on $(C_c(\mathbb{R}^k))^k$ (the space of \mathbb{R}^k -valued continuous functions on \mathbb{R}^k with compact support equipped with the sup-norm) given by

defmesf

$$(9) \quad \langle \mathbf{f}, \mu \rangle := \int \mathbf{f} \cdot d\mu, \quad \mathbf{f} \in (C_c(\mathbb{R}^k))^k.$$

The norm of the functional (9), is $\|\mu\|_{TV}$. It extends naturally with the same norm to the space $(C_0(\mathbb{R}^k))^k$ of \mathbb{R}^k -valued continuous functions on \mathbb{R}^k vanishing at infinity.

At places, we also identify μ with the restriction of (9) to $(C_c^\infty(\mathbb{R}^3))^3$, where $C_c^\infty(\mathbb{R}^3)$ is the space of C^∞ -smooth functions with compact support, equipped with the usual topology making it an LF-space [39]. We refer to a continuous linear functional on $(C_c^\infty(\mathbb{R}^m))^n$ as being a distribution.

We denote Lebesgue measure on \mathbb{R}^3 by \mathcal{L}_3 and d -dimensional Hausdorff measure on \mathbb{R}^3 by \mathcal{H}_d . We normalize \mathcal{H}_d for $d = 1$ and 2 so that it coincides with arclength and surface area for smooth curves and surfaces, respectively. We denote the Hausdorff dimension of a set $E \subset \mathbb{R}^3$ by $\dim_{\mathcal{H}}(E)$.

2. EQUIVALENT MAGNETIZATIONS, NET MOMENTS, AND TOTAL VARIATION

Sect3D

In this section, we discuss some regularity issues for magnetic fields and potentials, and we study the connection between S -silent sources and \mathbb{R}^3 -valued measures which are distributionally divergence-free. This rests on the notion of a *slender set* introduced in Definition 1. Subsequently, we solve Extremal Problem 1 for certain classes of magnetizations when S is slender. Such magnetizations are “sparse”, in the sense that either their support is purely 1-unrectifiable (see Theorem 2.6) or they assume a single direction on each piece of some finite partition of S (see Theorem 2.11). We also give conditions on S and Q ensuring that the forward operator has kernel the space of S -silent magnetizations (see Lemma 2.3).

We first define the *magnetic field* $\mathbf{b}(\mu)$ and the *scalar magnetic potential* $\Phi(\mu)$ (see [32]) generated by a magnetization distribution μ at points x not in the support of μ as

bDef1

$$(10) \quad \begin{aligned} \mathbf{b}(\mu)(x) &:= -c \left(\int \frac{1}{|x-y|^3} d\mu(y) - 3 \int (x-y) \frac{(x-y) \cdot d\mu(y)}{|x-y|^5} \right) \\ &= -c \operatorname{grad} \int \operatorname{grad}_y \frac{1}{|x-y|} \cdot d\mu(y) \\ &= -\mu_0 \operatorname{grad} \Phi(\mu)(x), \quad x \notin \operatorname{supp} \mu, \end{aligned}$$

where $c = 10^{-7} H m^{-1}$, $\mu_0 = 4\pi \times c$, and $\Phi(\boldsymbol{\mu})$ is given by

$$\boxed{\text{PhiDef}} \quad (11) \quad \Phi(\boldsymbol{\mu})(x) := \frac{1}{4\pi} \int \text{grad}_y \frac{1}{|x-y|} \cdot d\boldsymbol{\mu}(y) = \frac{1}{4\pi} \int \frac{x-y}{|x-y|^3} \cdot d\boldsymbol{\mu}(y), \quad x \notin \text{supp } \boldsymbol{\mu}.$$

Note that the second equality in (10) follows from the smoothness of the Newton kernel $1/|x-y|$ for $x \neq y$ and the formula

$$\text{grad}_x \left(\text{grad}_y \frac{1}{|x-y|} \cdot a \right) = \frac{a}{|x-y|^3} - 3(x-y) \frac{(x-y) \cdot a}{|x-y|^5},$$

for a fixed $a \in \mathbb{R}^3$. Note that $\Phi(\boldsymbol{\mu})$ and the components of $\mathbf{b}(\boldsymbol{\mu})$ are harmonic functions on $\mathbb{R}^3 \setminus S$. We show in Proposition 2.1 that Φ can be extended to a locally integrable function on \mathbb{R}^3 that satisfies, in the sense of distributions, the Poisson-Hodge equation (1). Furthermore, this proposition shows that the magnetic field $\mathbf{b}(\boldsymbol{\mu})$ extends to a divergence-free \mathbb{R}^3 -valued distribution.

$\boxed{\text{LocIntPhi}}$

Proposition 2.1. *Let S be a closed proper subset of \mathbb{R}^3 and $\boldsymbol{\mu} \in \mathcal{M}(S)^3$. Then, the integral in the right-hand side of (11) converges absolutely for a.e. $x \in \mathbb{R}^3$, thereby defining an extension of $\Phi(\boldsymbol{\mu})$ to all of \mathbb{R}^3 . Denoting this extension by $\Phi(\boldsymbol{\mu})$ again, it holds for each p, q with $1 \leq p < 3/2 < q \leq \infty$ that $\Phi(\boldsymbol{\mu}) \in L^p(\mathbb{R}^3) + L^q(\mathbb{R}^3)$, and that, in the distributional sense,*

$$\Delta \Phi(\boldsymbol{\mu}) = \text{div } \boldsymbol{\mu}.$$

Furthermore, the distribution $\mu_0(\boldsymbol{\mu} - \text{grad } \Phi(\boldsymbol{\mu}))$ is divergence-free and extends $\mathbf{b}(\boldsymbol{\mu})$ from a distribution on $\mathbb{R}^3 \setminus S$ to a distribution on \mathbb{R}^3 . Denoting this extension by $\mathbf{b}(\boldsymbol{\mu})$ again, we have that $\langle T_\alpha \mathbf{f}, \mathbf{b}(\boldsymbol{\mu}) \rangle \rightarrow 0$ as $|\alpha| \rightarrow \infty$ for every $\mathbf{f} \in (C_c^\infty(\mathbb{R}^3))^3$ (here $T_\alpha \mathbf{f}$ denotes the translation of the argument of \mathbf{f} by α).

Remark. The decomposition $\boldsymbol{\mu} = \mathbf{b}(\boldsymbol{\mu})/\mu_0 + \text{grad } \Phi(\boldsymbol{\mu})$ is the Helmholtz-Hodge decomposition of the \mathbb{R}^3 -valued measure $\boldsymbol{\mu}$ into the sum of a gradient and a divergence-free term. Although $\boldsymbol{\mu}$ is a distribution of order 0, note that the summands will generally have order -1. Hereafter, we let $\Phi(\boldsymbol{\mu})$ and $\mathbf{b}(\boldsymbol{\mu})$ denote the extensions given in Proposition 2.1.

Proof. Set $\mathbf{k}(y) = y/(4\pi|y|^3)$ for $y \in \mathbb{R}^3$ and let $\phi \in C_c^\infty(\mathbb{R}^3)$ be valued in $[0, 1]$, identically 1 on $\mathbb{B}(0, 1)$ and 0 outside of $\mathbb{B}(0, 2)$. Writing $\mathbf{f}_1 = \phi \mathbf{k}$ and $\mathbf{f}_2 = (1 - \phi) \mathbf{k}$, we have that $|\mathbf{f}_1| \in L^p(\mathbb{R}^3)$ for $1 \leq p < 3/2$ and $|\mathbf{f}_2| \in L^q(\mathbb{R}^3)$. From (11), we get that $\Phi(\boldsymbol{\mu}) = \mathbf{k} * \boldsymbol{\mu} = \mathbf{f}_1 * \boldsymbol{\mu} + \mathbf{f}_2 * \boldsymbol{\mu}$ off S (the product under the convolution integral is here the scalar product). For any $r \in [1, \infty]$, Jensen's inequality implies that the convolution of a finite signed measure with an L^r function is an L^r function with norm not exceeding the mass of the measure times the initial norm, and Fubini's theorem entails that the integrals converge absolutely a.e. Therefore $\mathbf{f}_1 * \boldsymbol{\mu} \in L^p(\mathbb{R}^3)$ and $\mathbf{f}_2 * \boldsymbol{\mu} \in L^q(\mathbb{R}^3)$, showing that $\mathbf{k} * \boldsymbol{\mu} \in L^p(\mathbb{R}^3) + L^q(\mathbb{R}^3)$.

We next show that

$\boxed{\text{DelPhi}}$

$$(12) \quad \Delta \Phi(\boldsymbol{\mu}) = \Delta(\mathbf{k} * \boldsymbol{\mu}) = \text{div } \boldsymbol{\mu}.$$

Let $N(y) = -1/(4\pi|y|)$ for $y \in \mathbb{R}^3$, pick $\psi \in C_c^\infty(\mathbb{R}^3)$ and recall that $G := N * \psi$ is a smooth function vanishing at infinity such that $\Delta G = \psi$ [1, Cor. 4.3.2&4.5.4]. Now, differentiating

under the integral sign, we have that $\mathbf{k} * \psi = \text{grad } G$, therefore

$$\begin{aligned} \langle \Delta \Phi(\boldsymbol{\mu}), \psi \rangle &= \langle \mathbf{k} * \boldsymbol{\mu}, \Delta \psi \rangle = -\langle \boldsymbol{\mu}, \mathbf{k} * \Delta \psi \rangle = -\langle \boldsymbol{\mu}, \Delta(\mathbf{k} * \psi) \rangle = -\langle \boldsymbol{\mu}, \Delta(\text{grad } G) \rangle \\ &= -\langle \boldsymbol{\mu}, \text{grad } \Delta G \rangle = \langle \text{div } \boldsymbol{\mu}, \Delta G \rangle = \langle \text{div } \boldsymbol{\mu}, \psi \rangle \end{aligned}$$

which proves (12).

Finally, the distribution $\mathbf{d} := \mu_0(\boldsymbol{\mu} - \text{grad } \Phi(\boldsymbol{\mu}))$ coincides with $\mathbf{b}(\boldsymbol{\mu})$ on $\mathbb{R}^3 \setminus S$, by (10), and it follows from (12) that it satisfies $\text{div } \mathbf{d} = \mu_0(\text{div } \boldsymbol{\mu} - \Delta \Phi(\boldsymbol{\mu})) = 0$; i.e., \mathbf{d} is divergence-free. Renaming \mathbf{d} as $\mathbf{b}(\boldsymbol{\mu})$ since it extends the latter, the finiteness of $\boldsymbol{\mu}$ and the fact that $\Phi(\boldsymbol{\mu}) \in L^1(\mathbb{R}^3)^3 + L^2(\mathbb{R}^3)^3$ show:

$$\langle T_\alpha \mathbf{f}, \mathbf{b}(\boldsymbol{\mu}) \rangle = \mu_0 (\langle T_\alpha \mathbf{f}, \boldsymbol{\mu} \rangle - \langle \text{div } T_\alpha \mathbf{f}, \Phi(\boldsymbol{\mu}) \rangle) \rightarrow 0$$

as $|\alpha| \rightarrow \infty$, for every $\mathbf{f} \in (C_c^\infty(\mathbb{R}^3))^3$. □

secDivSilent

2.1. Divergence-free and silent magnetizations. The equation $\Delta \Phi(\boldsymbol{\mu}) = \text{div } \boldsymbol{\mu}$ is suggestive of the existence of a relationship between S -silent and divergence-free magnetizations. As will be seen from the lemma below, for any closed set S all divergence-free magnetizations supported on S are S -silent, but the converse is not always true as follows from the next construction.

balldipole

Example 2.1. Let $S = \mathbb{B}(0, 1)$ be the closed unit Euclidean ball centered at the origin, $\mu = \mathcal{L}_3 \llcorner S$, and $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ the \mathbb{R}^3 -valued measure equal to $(4\pi/3)^{-1} \mu e_1$. Then, by the mean value theorem, we get that

$$\Phi(\boldsymbol{\mu})(x) = \frac{1}{4\pi} \int \text{grad}_y \frac{1}{|x-y|} \cdot d\boldsymbol{\mu}(y) = \frac{1}{4\pi} \frac{x_1}{|x|^3}, \quad x \notin S,$$

since $\frac{1}{4\pi} \frac{x_1}{|x|^3}$ is harmonic on $\mathbb{R}^3 \setminus \{0\}$. Note that $\frac{1}{4\pi} \frac{x_1}{|x|^3}$ is also the magnetic potential generated by the dipole $\boldsymbol{\nu} := \delta_0 e_1$; therefore $\boldsymbol{\mu}$ and $\boldsymbol{\nu}$ are S -equivalent, that is $\boldsymbol{\mu} - \boldsymbol{\nu}$ is S -silent. However, this magnetization is not divergence-free since, for every $f \in C_c^\infty(\mathbb{R}^3)$ supported in $\mathbb{B}(0, 1)$, it holds that $\langle f, \text{div } (\boldsymbol{\mu} - \boldsymbol{\nu}) \rangle = -\langle f, \text{div } \boldsymbol{\nu} \rangle = -\partial_{x_1} f(0)$.

An analogous argument shows that, for a_3 the area of $\mathbb{S}(0, 1)$ and $\tilde{\mu} := \mathcal{H}_2 \llcorner \mathbb{S}(0, 1)$, the \mathbb{R}^3 -valued measure $a_3^{-1} \tilde{\mu} e_1$ is likewise S -equivalent to $\boldsymbol{\nu}$.

The following variant of this example is also instructive: there is a sequence $x_n \in \mathbb{B}(0, 1)$ and a sequence c_n of real numbers with $\sum_n |c_n| < \infty$ such that the measure $\boldsymbol{\alpha}$ defined by $\boldsymbol{\alpha} = e_1 \sum_n c_n \delta_{x_n}$ is S -equivalent to $a_3^{-1} \tilde{\mu} e_1$ which is S -equivalent to $\boldsymbol{\nu}$. Therefore $\boldsymbol{\alpha} - \boldsymbol{\nu}$ is an S -silent magnetization consisting of countably many point dipoles. To see that x_n and c_n exist, recall Bonsall's theorem (whose proof in the ball is the same as in the disk, see [38, Thms. 5.21 & 5.22]) that whenever x_n is a sequence in $\mathbb{B}(0, 1)$ which is nontangentially dense in $\mathbb{S}(0, 1)$, each function $h \in L^1(\tilde{\mu})$ can be written as $h(\xi) = \sum_n c_n P_{x_n}(\xi)$ where $P_{x_n}(\xi) = (1/4\pi)(1 - |x_n|^2)/|\xi - x_n|^3$ is the familiar Poisson kernel of the unit ball at x_n , and c_n is a sequence of real numbers with absolutely convergent sum. Choosing $h \equiv 1$ and observing that, for $y \notin S$,

$$\frac{y - x_n}{|y - x_n|^3} = \int \frac{y - \xi}{|y - \xi|^3} P_{x_n}(\xi) d\tilde{\mu}(\xi)$$

by the Poisson representation of harmonic functions, we easily check that $e_1 \sum_n c_n \delta_{x_n}$ is S equivalent to $a_3^{-1} \tilde{\mu} e_1$, as desired.

Example 2.1 shows us that a \mathbb{R}^3 -valued measure on S which is S -silent needs not be divergence-free in general. It does, however, when S is slender:

divM **Theorem 2.2.** *Let $S \in \mathbb{R}^3$ be closed and $\boldsymbol{\mu} \in \mathcal{M}(S)^3$. If $\operatorname{div} \boldsymbol{\mu} = 0$, then $\boldsymbol{\mu}$ is S -silent. Furthermore, if S is a slender set and $\boldsymbol{\mu}$ is S -silent, then $\operatorname{div} \boldsymbol{\mu} = 0$.*

Proof. Since $\Phi(\boldsymbol{\mu}) \in L^1(\mathbb{R}^3)^3 + L^2(\mathbb{R}^3)^3$ by Proposition 2.1, we get from the Schwarz inequality:

inegPhig (13)
$$\int_E |\Phi(\boldsymbol{\mu})| d\mathcal{L}_3 \leq C_1 + C_2 \mathcal{L}_3^{1/2}(E)$$

for some constants C_1, C_2 and each Borel set E of finite measure. If $\operatorname{div} \boldsymbol{\mu} = 0$, then $\Phi(\boldsymbol{\mu})$ is harmonic on \mathbb{R}^3 by the same lemma, therefore it is constant by (13) and the mean value theorem (see proof of [2, Thm. 2.1]). Consequently $\boldsymbol{\mu}$ is S -silent.

For the second statement, assume that S is a slender set and that $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is S -silent. Since $\mathbf{b}(\boldsymbol{\mu}) = \operatorname{grad} \Phi(\boldsymbol{\mu})$ and $\boldsymbol{\mu}$ is S -silent, then $\Phi(\boldsymbol{\mu})$ is constant on each connected component of $\mathbb{R}^3 \setminus S$. If \mathfrak{C} is such a component, we can apply (13) with $E = \mathfrak{C} \cap B(0, n)$ and let $n \rightarrow \infty$ to conclude that the corresponding constant is zero, because $\mathcal{L}_3(\mathfrak{C}) = +\infty$. Hence, $\Phi(\boldsymbol{\mu})$ must be zero on $\mathbb{R}^3 \setminus S$, and since $\mathcal{L}_3(S) = 0$ it follows that $\Phi(\boldsymbol{\mu})$ is zero as a distribution, so that $\operatorname{div} \boldsymbol{\mu} = \Delta \Phi(\boldsymbol{\mu}) = 0$. \square

In typical Scanning Magnetic Microscopy experiments, data consists of point-wise values of one component of the magnetic field taken on a plane not intersecting S . Of course, finitely many values do not characterize the field, but it is natural to ask how one can choose the measurement points to ensure that infinitely many of them would, in the limit, determine $\mathbf{b}(\boldsymbol{\mu})$ uniquely. We next provide a sufficient condition that such data (more generally, data measured on an analytic surface which needs not be a plane) determines the field in the complement of S . The condition dwells on the remark that a nonzero real analytic function on a connected open subset of \mathbb{R}^k has a zero set of Hausdorff dimension at most $k - 1$. It is so because the zero set is locally a countable union of smooth (even real-analytic) embedded submanifolds of strictly positive codimension, see [33, thm 5.2.3]. This fact sharpens the property that the zero set of a nonzero real analytic function in \mathbb{R}^k has Lebesgue measure zero, and will be used at places in the paper. Using local coordinates, it is immediately checked that the previous bound on the Hausdorff dimension remains valid when \mathbb{R}^k is replaced by a smooth real-analytic manifold embedded in \mathbb{R}^m for some $m > k$.

We also need at this point a version of the Jordan-Brouwer separation theorem for a connected, properly embedded (*i.e.* complete but not necessarily compact) surface in \mathbb{R}^3 , to the effect that the complement of such a surface has two connected components. In the smooth case which is our concern here, we give in Appendix A a short, differential topological argument for this result which we assume is known but for which we could not find a published reference.

Hereafter, we will say that $E, F \subset \mathbb{R}^3$ are positively separated from each other if $\text{dist}(E, F) > 0$, where $\text{dist}(E, F) := \inf_{x \in E, y \in F} |x - y|$.

Lem4

Lemma 2.3. *Let $S \subset \mathbb{R}^3$ be closed and suppose $\mathbb{R}^3 \setminus S$ is connected and contains a nonempty open half-cylinder of direction $v \in \mathbb{R}^3 \setminus \{0\}$. Furthermore, let \mathcal{A} be a smooth complete and connected real analytic surface in $\mathbb{R}^3 \setminus S$ that is positively separated from S and such that S lies entirely within one of the two connected components of $\mathbb{R}^3 \setminus \mathcal{A}$. Let also $Q \subset \mathbb{R}^3 \setminus S$ be such that the closure of $Q \cap \mathcal{A}$ has Hausdorff dimension strictly greater than 1. If $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is such that $v \cdot \mathbf{b}(\boldsymbol{\mu})$ vanishes on $Q \cap \mathcal{A}$, then $\boldsymbol{\mu}$ is S -silent.*

Proof. Suppose $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is such that $v \cdot \mathbf{b}(\boldsymbol{\mu})$ vanishes on $Q \cap \mathcal{A}$. As $v \cdot \mathbf{b}(\boldsymbol{\mu})$ is harmonic in $\mathbb{R}^3 \setminus S$ it is real-analytic there, and since $v \cdot \mathbf{b}(\boldsymbol{\mu})$ vanishes on the closure of $Q \cap \mathcal{A}$ which has Hausdorff dimension > 1 it must vanish identically on \mathcal{A} .

Observe now that $\mathbb{R}^3 \setminus \mathcal{A}$ has two connected components (see Theorem A.1), and let \mathcal{U} be the one not containing S . Note, using (10), that if $\boldsymbol{\nu} \in \mathcal{M}(S)^3$ and $x \notin \text{supp}(\boldsymbol{\nu})$, then,

eqn:b_bound

$$(14) \quad |\mathbf{b}(\boldsymbol{\nu})(x)| \leq 4c (\text{dist}(x, \text{supp}(\boldsymbol{\nu})))^{-3} \|\boldsymbol{\nu}\|_{TV}.$$

For $R > 0$ let $\boldsymbol{\mu}_R := \boldsymbol{\mu}|_{B(0, R)}$ and $\tilde{\boldsymbol{\mu}}_R := \boldsymbol{\mu} - \boldsymbol{\mu}_R$. Then

$$\mathbf{b}(\boldsymbol{\mu})(x) = \mathbf{b}(\boldsymbol{\mu}_R)(x) + \mathbf{b}(\tilde{\boldsymbol{\mu}}_R)(x),$$

and applying (14) to $\boldsymbol{\mu}_R$ and $\tilde{\boldsymbol{\mu}}_R$ for R large enough, using that \mathcal{A} is positively separated from S , we get that $\limsup_{x \in \mathcal{U}, |x| \rightarrow \infty} |\mathbf{b}(\boldsymbol{\mu})(x)| < \varepsilon$ for any $\varepsilon > 0$, hence $\mathbf{b}(\boldsymbol{\mu})(x) \rightarrow 0$ as $x \rightarrow \infty$ in \mathcal{U} . Since $v \cdot \mathbf{b}(\boldsymbol{\mu})$ vanishes on the boundary of \mathcal{U} , we may use the maximum principle to conclude that $v \cdot \mathbf{b}(\boldsymbol{\mu})$ vanishes on \mathcal{U} and therefore on $\mathbb{R}^3 \setminus S$ as the complement of S is connected.

This implies that the magnetic potential $\Phi(\boldsymbol{\mu})$ is constant on every line segment parallel to v not intersecting S . Now, $\mathbb{R}^3 \setminus S$ contains a half-cylinder \mathcal{C} of direction v , and shrinking the latter if necessary we may assume it is positively separated from S . From (11) we get

eqn:Phi_bound

$$(15) \quad |\Phi(\boldsymbol{\nu})(x)| \leq \frac{1}{4\pi} (\text{dist}(x, \text{supp}(\boldsymbol{\nu})))^{-2} \|\boldsymbol{\nu}\|_{TV}$$

for $\boldsymbol{\nu} \in \mathcal{M}(S)^3$ and $x \notin \text{supp}(\boldsymbol{\nu})$, and we conclude that $\Phi(\boldsymbol{\mu})(x)$ goes to zero as $x \rightarrow \infty$ in \mathcal{C} . Hence its value on each half line contained in \mathcal{C} is zero, so $\Phi(\boldsymbol{\mu}) \equiv 0$ in \mathcal{C} . Consequently it vanishes identically (thus also $\mathbf{b}(\boldsymbol{\mu})$) in the connected open set $\mathbb{R}^3 \setminus S$, by real analyticity. \square

The following example shows that Lemma 2.3 needs not hold if S is not contained in a single component of $\mathbb{R}^3 \setminus \mathcal{A}$ or if \mathcal{A} fails to be analytic.

Example 2.2. Let S be equal to $\{e_3, -e_3\}$, $\boldsymbol{\mu} = (\delta_{e_3} + \delta_{-e_3})e_2$, $v = e_3$, and $\mathcal{A} = \{x_3 = 0\}$. Then $e_3 \cdot \mathbf{b}(\boldsymbol{\mu})$ is zero on \mathcal{A} but $\boldsymbol{\mu}$ is not S -silent. Also, whenever Q is a bounded subset of $\{x_3 = 0\}$ with $\dim_{\mathcal{H}} Q > 1$, there is a closed C^∞ -smooth surface Z containing Q such that e_3 and $-e_3$ lie in the same component of $\mathbb{R}^3 \setminus Z$; however, Z cannot be analytic.

Lemma 2.3 does not hold either if $\mathbb{R}^3 \setminus S$ is disconnected: for instance, if $S = \{x_3 = 0\}$, there is a $\boldsymbol{\mu}$ of the form $\mathbf{hd}\mathcal{L}_2$ with $\mathbf{h} \in (\mathfrak{h}^1(\mathbb{R}^2))^3$, where $\mathfrak{h}^1(\mathbb{R}^2)$ is the real Hardy space on \mathbb{R}^2 , such that $\mathbf{b}(\boldsymbol{\mu})$ vanishes identically for $\{x_3 < 0\}$ but not for $\{x_3 > 0\}$ [7, Thm. 3.2].

In particular, if we let $Q = \mathcal{A} = \{x_3 = -1\}$, we have that $\mathbf{b}(\boldsymbol{\mu})$ vanishes on Q but $\boldsymbol{\mu}$ is not S -silent. However, if $\mathbb{R}^3 \setminus S$ is not connected but has a connected component \mathcal{V} containing a half-cylinder, then by replacing S with $\tilde{S} := \mathbb{R}^3 \setminus \mathcal{V}$ and selecting an appropriate \mathcal{A} , Q and v , Lemma 2.3 may be applied to the effect that $\boldsymbol{\mu}$ is \tilde{S} -silent whenever $v \cdot \mathbf{b}(\boldsymbol{\mu})$ vanishes on Q . Thus, if each component \mathcal{V}_i of $\mathbb{R}^3 \setminus S$ contains a half-cylinder and can be associated with suitable \mathcal{A}_i , Q_i and v_i , and if $v_i \cdot \mathbf{b}(\boldsymbol{\mu})$ vanishes on Q_i for all i , then $\boldsymbol{\mu}$ is S -silent.

Remark. If S is as Lemma 2.3 and $Q \subset \mathbb{R}^3$ is positively separated from S and has closure \bar{Q} of Hausdorff dimension > 2 , then the conclusion of the lemma still holds and the proof is easier. In this case indeed, it follows directly from the hypothesis on Q that $v \cdot \mathbf{b}(\boldsymbol{\mu})$ is identically zero in $\mathbb{R}^3 \setminus S$ as soon as it vanishes on Q , and the rest of the proof is as before. We shall not investigate this situation which, from the point of view of inverse problems, corresponds to the case where measurements of the field are taken in a volume rather than on a surface. Though more information can be gained this way, the experimental and computational burden often becomes deterring.

components

Lemma 2.4. *Let $S = S_0 \cup S_1 \subset \mathbb{R}^3$ for some disjoint closed sets S_0 and S_1 . If $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is S -silent, then for $i = 0, 1$ the restriction $\boldsymbol{\mu}|_{S_i}$ is S_i -silent.*

Proof. Let $\mathbf{b}_0 = \mathbf{b}(\boldsymbol{\mu}|_{S_0})$ and $\mathbf{b}_1 = \mathbf{b}(\boldsymbol{\mu}|_{S_1})$. Note that \mathbf{b}_0 and \mathbf{b}_1 are harmonic in $\mathbb{R}^3 \setminus S_0$ and $\mathbb{R}^3 \setminus S_1$ respectively. Also, as $\boldsymbol{\mu}$ is S -silent, it holds that $\mathbf{b}_0(x) = -\mathbf{b}_1(x)$ for $x \notin S$. Hence the function

$$\tilde{\mathbf{b}}(x) = \begin{cases} \mathbf{b}_0(x) & x \in \mathbb{R}^3 \setminus S_0, \\ -\mathbf{b}_1(x) & x \in \mathbb{R}^3 \setminus S_1, \end{cases}$$

is harmonic on \mathbb{R}^3 (note that the two definitions agree on $\mathbb{R}^3 \setminus S$). Moreover, Proposition 2.1 implies that for every $\mathbf{f} \in (C_c^\infty(\mathbb{R}^3))^3$:

$$(16) \quad |\langle T_\alpha \mathbf{f}, \tilde{\mathbf{b}} \rangle| \leq |\langle T_\alpha \mathbf{f}, \mathbf{b}_0 \rangle| + |\langle T_\alpha \mathbf{f}, \mathbf{b}_1 \rangle| \rightarrow 0 \quad \text{as} \quad |\alpha| \rightarrow \infty.$$

Since $\tilde{\mathbf{b}}$ is harmonic, the mean value property applied to (16) with $\mathbf{f} = \varphi e_j$, where φ is a non-negative radially symmetric smooth function with support $\bar{\mathbb{B}}(0, 1)$, implies that $\tilde{\mathbf{b}}(x)$ vanishes as $x \rightarrow \infty$ and therefore is identically 0 by Liouville's theorem. Thus both \mathbf{b}_0 and \mathbf{b}_1 are zero on $\mathbb{R}^3 \setminus S_0$ and $\mathbb{R}^3 \setminus S_1$ respectively, and hence $\boldsymbol{\mu}|_{S_0}$ is S_0 -silent and $\boldsymbol{\mu}|_{S_1}$ is S_1 -silent. \square

sparse3D:Sec

2.2. Decomposition of divergence-free magnetizations and recovery of magnetizations with sparse support. A set $E \subset \mathbb{R}^2$ is said to be *1-rectifiable* (e.g., see [35, Def. 15.3]) if there exist Lipschitz maps $f_i : \mathbb{R} \rightarrow \mathbb{R}^n$, $i = 1, 2, \dots$, such that

$$\mathcal{H}_1 \left(E \setminus \bigcup_{i=1}^{\infty} f_i(\mathbb{R}) \right) = 0.$$

A set $B \subset \mathbb{R}^n$ is *purely 1-unrectifiable* if $\mathcal{H}_1(E \cap B) = 0$ for every 1-rectifiable set E . Clearly a set of \mathcal{H}_1 -measure zero is purely 1-unrectifiable. A purely 2-unrectifiable set is defined in the same way, only with \mathcal{H}_2 instead of \mathcal{H}_1 .

We call a Lipschitz mapping $\gamma : [0, \ell] \rightarrow \mathbb{R}^3$ a *rectifiable curve* and let $\Gamma := \gamma([0, \ell])$ denote its image. If γ is an arclength parametrization of Γ ; i.e., if γ satisfies

$$\boxed{\text{arc}} \quad (17) \quad \mathcal{H}_1(\gamma([\alpha, \beta])) = \beta - \alpha, \quad \forall [\alpha, \beta] \subset [0, \ell],$$

then we call γ an *oriented rectifiable curve*. By Rademacher's Theorem (see [26]), γ is differentiable a.e. on $[0, \ell]$. Furthermore, it follows from (17) that $|\gamma'(t)| = 1$ a.e. on $[0, \ell]$. For a given oriented rectifiable curve γ we define $\mathbf{R}_\gamma \in \mathcal{M}(S)^3$ through the relation

$$(18) \quad \langle \mathbf{R}_\gamma, \mathbf{f} \rangle = \int_0^\ell \mathbf{f}(\gamma(t)) \cdot \gamma'(t) dt,$$

for $\mathbf{f} \in C_0(\mathbb{R}^3)^3$. Alternatively, since \mathbf{R}_γ is absolutely continuous with respect to \mathcal{H}_1 we may consider the Radon-Nikodym derivative τ of \mathbf{R}_γ with respect to \mathcal{H}_1 and we remark that $\tau(\gamma(t)) = \gamma'(t)$ for a.e. $t \in [0, \ell]$. Then, for a Borel set $B \subset \mathbb{R}^3$ we have

$$\boxed{\text{Rgamma}} \quad (19) \quad \mathbf{R}_\gamma(B) = \int_B \tau d(\mathcal{H}_1 \llcorner \Gamma).$$

We remark that if B is purely 1-unrectifiable, then $|\mathbf{R}_\gamma|(B) = \mathcal{H}_1(B \cap \Gamma) = 0$ and, furthermore, Fubini's Theorem implies $\mathcal{L}_3(B) = 0$.

Let $\mathcal{C} \subset \mathcal{M}(S)^3$ denote the collection of oriented rectifiable curves with topology inherited from $\mathcal{M}(S)^3$. Suppose $\text{div } \boldsymbol{\mu} = 0$ (as a distribution). Smirnov [40, Theorem A] shows that $\boldsymbol{\mu}$ can be decomposed into elements from \mathcal{C} . In particular, it can be proven that there is a positive measure ρ on \mathcal{C} such that

$$\boxed{\text{Smi1}} \quad (20) \quad \boldsymbol{\mu}(B) = \int \mathbf{R}(B) d\rho(\mathbf{R}),$$

and

$$\boxed{\text{Smi2}} \quad (21) \quad |\boldsymbol{\mu}|(B) = \int |\mathbf{R}|(B) d\rho(\mathbf{R}),$$

for any Borel set $B \subset \mathbb{R}^3$. From the representation (20) of a divergence-free magnetization, we immediately obtain the following lemma.

$\boxed{\text{1unrecdiv}}$ **Lemma 2.5.** *Suppose $S \subset \mathbb{R}^3$ is closed and purely 1-unrectifiable. If $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is divergence-free, then $\boldsymbol{\mu} = 0$.*

$\boxed{\text{unrec3D}}$ **Theorem 2.6.** *Suppose $S \subset \mathbb{R}^3$ is a closed, slender set. If $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ has support that is purely 1-unrectifiable and $\boldsymbol{\nu} \in \mathcal{M}(S)^3$ is S -equivalent to $\boldsymbol{\mu}$, then $\|\boldsymbol{\nu}\|_{TV} > \|\boldsymbol{\mu}\|_{TV}$ unless $\boldsymbol{\nu} = \boldsymbol{\mu}$.*

Proof. Since $\boldsymbol{\theta} := \boldsymbol{\nu} - \boldsymbol{\mu}$ is S -silent, Theorem 2.2 implies $\text{div } \boldsymbol{\theta} = 0$ and so $\boldsymbol{\theta}$ can be represented in the form (20), where (21) holds. Since the support of $\boldsymbol{\mu}$ is purely 1-unrectifiable, it follows from (21) and the remark after (19) that the measures $\boldsymbol{\mu}$ and $\boldsymbol{\theta}$ are mutually singular. Thus, $\|\boldsymbol{\nu}\|_{TV} = \|\boldsymbol{\mu}\|_{TV} + \|\boldsymbol{\theta}\|_{TV} > \|\boldsymbol{\mu}\|_{TV}$ unless $\boldsymbol{\nu} = \boldsymbol{\mu}$. \square

Snotslender

Example 2.3. Recall from Example 2.1 that if $S = \overline{\mathbb{B}}(0, 1)$, the magnetizations $\boldsymbol{\mu}$ modeling a uniformly magnetized ball and $\boldsymbol{\nu}$ which is a point dipole at 0, with the same net moment as $\boldsymbol{\mu}$, are S -equivalent. Moreover, it is easy to verify that $\|\boldsymbol{\mu}\|_{TV} = \|\boldsymbol{\nu}\|_{TV} = |\langle \boldsymbol{\nu} \rangle|$. Since the support of $\boldsymbol{\nu}$ is a single point, it is purely 1-unrectifiable, hence the assumption that S is a slender set cannot be eliminated from Theorem 2.6.

The previous example entails that total variation minimization is not sufficient alone to distinguish magnetizations with purely 1-unrectifiable support among all equivalent magnetizations supported on S , when S is not slender. However, as the following result shows, the recovery problem for general S has at most one solution when restricted to magnetizations whose support is purely 1-unrectifiable *and* has finite \mathcal{H}_2 measure.

To see this, we shall need a consequence of the Besicovitch-Federer Projection Theorem [35, Thm 18.1]; namely that the complement of a closed, purely 2-unrectifiable set with finite \mathcal{H}_2 measure is connected. We will also need the fact that a purely 1-unrectifiable set is purely 2-unrectifiable. We are confident these facts are known (e.g., see the introduction in [20]), but since we have not explicitly found proofs in the literature, we provide outlines of the arguments.

With regard to the first fact, let $F \subset \mathbb{R}^3$ be a closed, purely 2-unrectifiable set with finite \mathcal{H}_2 measure and suppose $\mathbb{B}(x, r), \mathbb{B}(y, r)$ are disjoint balls in $\mathbb{R}^3 \setminus F$. By the Besicovitch-Federer Projection Theorem there exists a plane P such that the intersection of the orthogonal projections of these balls onto P minus the orthogonal projection of F onto P is nonempty and therefore the balls can be joined by a line segment not intersecting F .

As to the second fact, it follows from [27, Lemma 3.2.18] that it is enough for a set F to be purely 2-unrectifiable that $\mathcal{H}_2(F \cap \psi(K)) = 0$ for any compact set $K \subset \mathbb{R}^2$ and any bi-Lipschitz mapping $\psi : K \rightarrow \mathbb{R}^3$. Since bi-Lipschitz maps preserve unrectifiability we may restrict our considerations to \mathbb{R}^2 where the result follows easily from Fubini's theorem.

As a consequence of these facts, the complement of a closed, purely 1-unrectifiable set with finite \mathcal{H}_2 measure must be connected.

Cor-unrec3D

Corollary 2.7. *Suppose S is a closed, proper subset of \mathbb{R}^3 and that $\boldsymbol{\nu} \in \mathcal{M}(S)^3$ has purely 1-unrectifiable support of finite \mathcal{H}_2 measure. If $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is S -equivalent to $\boldsymbol{\nu}$ but not equal to $\boldsymbol{\nu}$, then the support of $\boldsymbol{\mu}$ is not a purely 1-unrectifiable set with finite \mathcal{H}_2 measure.*

Proof. Suppose $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ is S -equivalent to $\boldsymbol{\nu}$ and has support that is purely 1-unrectifiable with finite \mathcal{H}_2 measure. Then the support \tilde{S} of $\boldsymbol{\mu} - \boldsymbol{\nu}$ is also purely 1-unrectifiable with finite \mathcal{H}_2 measure. Therefore, its complement is connected and thus \tilde{S} is slender.

Moreover, $\boldsymbol{\mu} - \boldsymbol{\nu}$ is S -silent, hence its field vanishes on the nonempty open set $\mathbb{R}^3 \setminus S$, and since \tilde{S} is closed with $\mathcal{L}_3(\tilde{S}) = 0$ (because it is slender), the field must vanish on a nonempty open subset of $\mathbb{R}^3 \setminus \tilde{S}$. Since \tilde{S} has connected complement, we conclude that $\boldsymbol{\mu} - \boldsymbol{\nu}$ is \tilde{S} -silent, by real analyticity. Consequently, it is divergence-free by Theorem 2.2 and hence, by Lemma 2.5, $\boldsymbol{\mu} - \boldsymbol{\nu}$ is the zero measure. \square

Corollary 2.7 applies in particular if $\boldsymbol{\nu}$ is a finite sum of point dipoles. However, in view of Example 2.1, it does not apply in general to a convergent series of point dipoles.

2.3. The net moment of silent magnetizations. Our next result shows that, under certain assumptions on their support, silent measures have vanishing moment:

MomZero

Lemma 2.8. *Let $S \subset \mathbb{R}^3$ be a closed set and $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ be S -silent. Assume that one of the following conditions is satisfied:*

- (a) S is compact,
- (b) $\operatorname{div} \boldsymbol{\mu} = 0$.

Then the net moment $\langle \boldsymbol{\mu} \rangle = 0$.

Proof. Fix $i \in \{1, 2, 3\}$. Let $\phi \in C_0^\infty(\mathbb{R}^3)$ be supported in $\mathbb{B}(0, 2)$, $\phi(x) = x_i$ on $\mathbb{B}(0, 1)$ and for any $n > 0$ let $\phi_n(x) := n\phi(x/n)$. Note that $\|\operatorname{grad} \phi\|_\infty = \|\operatorname{grad} \phi_n\|_\infty$, that ϕ_n is supported in $\mathbb{B}(0, 2n)$, and for $x \in \overline{\mathbb{B}(0, n)}$ that $\phi_n(x) = x_i$, $\operatorname{grad} \phi_n(x) = e_i$ and $\Delta \phi_n(x) = 0$.

If S is compact take $n > 0$ such that $S \subset \mathbb{B}(0, n)$. Then $\langle \operatorname{grad} \phi_n, \boldsymbol{\mu} \rangle = \langle \boldsymbol{\mu} \rangle_i$, the i -th component of the moment of $\boldsymbol{\mu}$. Since $\boldsymbol{\mu}$ is S -silent $\Phi(\boldsymbol{\mu})$ is constant in each connected component of $\mathbb{R}^3 \setminus S$, and applying (13) with $E = \mathbb{B}(0, R) \setminus \mathbb{B}(0, n)$ we conclude on letting $R \rightarrow \infty$ that $\Phi(\boldsymbol{\mu}) \equiv 0$ in the unbounded component. Thus, $\Phi(\boldsymbol{\mu})$ is supported on $\mathbb{B}(0, n)$ and since $\Delta \Phi(\boldsymbol{\mu}) = \operatorname{div} \boldsymbol{\mu}$ by Proposition 2.1, we obtain:

$$\langle \boldsymbol{\mu} \rangle_i = \langle \operatorname{grad} \phi_n, \boldsymbol{\mu} \rangle = \langle \phi_n, \operatorname{div} \boldsymbol{\mu} \rangle = \langle \phi_n, \Delta \Phi(\boldsymbol{\mu}) \rangle = \langle \Delta \phi_n, \Phi(\boldsymbol{\mu}) \rangle = 0.$$

Therefore, taking $i = 1, 2, 3$, we get that $\langle \boldsymbol{\mu} \rangle = 0$, as announced.

Assume next that $\operatorname{div} \boldsymbol{\mu} = 0$. For any integer $m > 0$, let $D_m := \mathbb{B}(0, 2^{m+1}) \setminus \mathbb{B}(0, 2^m)$ and $M_m := |\boldsymbol{\mu}|(D_m)$. Because $\sum_m M_m \leq \|\boldsymbol{\mu}\|_{TV} < \infty$, we have that

liman

$$(22) \quad |\langle \operatorname{grad} \phi_{2^m}|_{D_m}, \boldsymbol{\mu} \rangle| \leq M_m \|\operatorname{grad} \phi\|_\infty \rightarrow 0 \quad \text{as } m \rightarrow \infty.$$

Now, let \mathbf{E}_i be the constant function equal to e_i on \mathbb{R}^3 . By (22), we see that

$$\lim_{m \rightarrow \infty} \langle \operatorname{grad} \phi_{2^m}, \boldsymbol{\mu} \rangle = \lim_{m \rightarrow \infty} \langle \operatorname{grad} \phi_{2^m}|_{\mathbb{B}(0, 2^m)}, \boldsymbol{\mu} \rangle = \langle \mathbf{E}_i, \boldsymbol{\mu} \rangle = \langle \boldsymbol{\mu} \rangle_i,$$

and since $\langle \operatorname{grad} \phi_n, \boldsymbol{\mu} \rangle = \langle \phi_n, \operatorname{div} \boldsymbol{\mu} \rangle = 0$ for each $n > 0$, by our assumption, we conclude that $\langle \boldsymbol{\mu} \rangle = 0$, as desired. \square

Assumptions (a) or (b) cannot be dropped in Lemma 2.8, for it is not sufficient that a magnetization be S -silent for its net moment to vanish, as shown by the following example.

momsil

Example 2.4. Consider the case where $S = \mathbb{R}^3 \setminus \mathbb{B}(0, R)$ and let $\boldsymbol{\mu} = v\mathcal{H}_2[\mathbb{S}(0, R)]$ where $v \in \mathbb{R}^3 \setminus \{0\}$. The density of $\boldsymbol{\mu}$ with respect to $\mathcal{H}_2[\mathbb{S}(0, R)]$ is the constant map $\mathbf{v} : \mathbb{S}(0, R) \rightarrow \mathbb{R}^3$ given by $\mathbf{v}(x) = v$, which is the trace on $\mathbb{S}(0, R)$ of the gradient of the function $x \mapsto v \cdot x$ which is harmonic on a neighborhood of $\overline{\mathbb{B}(0, R)}$, hence \mathbf{v} *a fortiori* belongs to the Hardy space $\mathcal{H}_{+, R}^2$ of harmonic gradients in $\mathbb{B}(0, R)$. Therefore $\boldsymbol{\mu}$ is silent in that ball [6, Lemma 4.2], and still $\langle \boldsymbol{\mu} \rangle = 4\pi v$. Integrating this example over $R \in [1, \infty)$ against the weight $1/R^4$ further shows that the \mathbb{R}^3 -valued measure $d\nu(x) = \mathbf{v}|x|^{-4}\chi_{\{|x| \geq 1\}}(x)d\mathcal{L}_3(x)$, is silent in the ball $\mathbb{B}(0, 1)$ but has $\langle \boldsymbol{\nu} \rangle = 4\pi v$. This provides us with an example of a (non-compactly supported) measure with non-zero total moment which is silent in the complement of its support.

uni3D:Sec

2.4. Total variation and unidirectional magnetizations. For $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ we can write $d\boldsymbol{\mu} = \mathbf{u}_\mu d|\boldsymbol{\mu}|$, therefore the Cauchy Schwarz inequality yields that

$$(23) \quad |\langle \boldsymbol{\mu} \rangle|^2 = \int \langle \boldsymbol{\mu} \rangle \cdot \mathbf{u}_\mu d|\boldsymbol{\mu}| \leq |\langle \boldsymbol{\mu} \rangle| \|\boldsymbol{\mu}\|_{TV},$$

where equality holds if and only if either $\langle \boldsymbol{\mu} \rangle = 0$ or $\mathbf{u}_\mu = \langle \boldsymbol{\mu} \rangle / |\langle \boldsymbol{\mu} \rangle|$ a. e. with respect to $|\boldsymbol{\mu}|$. We say that $\boldsymbol{\mu}$ is *uni-directional* if \mathbf{u}_μ is constant a.e. with respect to $|\boldsymbol{\mu}|$ (note that the zero magnetization is uni-directional). Thus, (23) implies the following:

lemuni1

Lemma 2.9. *If $\boldsymbol{\mu} \in \mathcal{M}(S)^3$, then $|\langle \boldsymbol{\mu} \rangle| \leq \|\boldsymbol{\mu}\|_{TV}$ with equality if and only if $\boldsymbol{\mu}$ is uni-directional.*

We call a magnetization *uni-dimensional* if it is the difference of two uni-directional magnetizations. The next lemma states that a uni-dimensional magnetization which is divergence-free must be the zero magnetization.

unidiv0

Lemma 2.10. *If $\boldsymbol{\mu} \in \mathcal{M}(\mathbb{R}^3)^3$ is uni-dimensional and $\operatorname{div} \boldsymbol{\mu} = 0$, then $\boldsymbol{\mu} = 0$.*

Proof. Suppose $\boldsymbol{\mu} \in \mathcal{M}(\mathbb{R}^3)^3$ is uni-dimensional and divergence-free. Then $\boldsymbol{\mu} = \mu v$ for some $v \in \mathbb{R}^3$ and $\mu \in \mathcal{M}(\mathbb{R}^3)$, with $0 = \operatorname{div}(\mu v) = v \cdot \operatorname{grad} \mu$. A standard argument (see below) shows that μ is translation invariant with respect to any vector parallel to v and therefore μ is finite only if it is zero.

Without loss of generality we may assume that $v = (1, 0, 0)$. To see the translation invariance of μ , take any $f \in C_c^\infty(\mathbb{R}^3)$ and let \tilde{f} be a translation of f in the x_1 direction. Then $f - \tilde{f} = \partial_{x_1} g$ with $g \in C_c^\infty(\mathbb{R}^3)$ defined by:

$$g(x_1, x_2, x_3) := \int_{-\infty}^{x_1} (f - \tilde{f})(y, x_2, x_3) d\mathcal{L}_1(y),$$

and so $\mu(f - \tilde{f}) = \langle \partial_{x_1} g, \mu \rangle = -\langle g, \partial_{x_1} \mu \rangle = 0$. □

UnidirThm

Theorem 2.11. *Let $S = \bigcup_{i=1}^n S_i$ for some disjoint closed sets S_1, S_2, \dots, S_n in \mathbb{R}^3 and suppose that S is either compact or slender. Let $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ be such that $\boldsymbol{\mu}_i := \boldsymbol{\mu}|_{S_i}$ is uni-directional for $i = 1, 2, \dots, n$.*

If $\boldsymbol{\nu} \in \mathcal{M}(S)^3$ is S -equivalent to $\boldsymbol{\mu}$, then $\boldsymbol{\nu}_i := \boldsymbol{\nu}|_{S_i}$ and $\boldsymbol{\mu}_i$ are S_i -equivalent for $i = 1, 2, \dots, n$, moreover

uniminTV

$$(24) \quad \|\boldsymbol{\mu}\|_{TV} \leq \|\boldsymbol{\nu}\|_{TV},$$

with equality in (24) if and only if $\boldsymbol{\nu}_i$ is uni-directional in the same direction as $\boldsymbol{\mu}_i$ for $i = 1, 2, \dots, n$. Furthermore, if S is slender and equality holds in (24), then $\boldsymbol{\mu} = \boldsymbol{\nu}$.

Proof. Since $\boldsymbol{\mu}$ and $\boldsymbol{\nu}$ are S -equivalent, their difference $\boldsymbol{\tau} := \boldsymbol{\nu} - \boldsymbol{\mu}$ is S -silent. In addition, it follows from Theorem 2.2 that if S is slender then $\operatorname{div} \boldsymbol{\tau} = 0$. By Lemma 2.4, the restriction $\boldsymbol{\tau}_i := \boldsymbol{\tau}|_{S_i}$ is S_i -silent and thus $\boldsymbol{\mu}_i$ and $\boldsymbol{\nu}_i$ are S_i -equivalent. Since either S is compact

or $\operatorname{div} \boldsymbol{\tau} = 0$, the same is true of each S_i , $\boldsymbol{\tau}_i$ and we can use Lemma 2.8 to obtain that $\langle \boldsymbol{\mu}_i \rangle = \langle \boldsymbol{\nu}_i \rangle$ for $i = 1, 2, \dots, n$. Then

$$\boxed{\text{numuTV}} \quad (25) \quad \|\boldsymbol{\nu}\|_{TV} = \sum_{i=1}^n \|\boldsymbol{\nu}_i\|_{TV} \geq \sum_{i=1}^n |\langle \boldsymbol{\nu}_i \rangle| = \sum_{i=1}^n |\langle \boldsymbol{\mu}_i \rangle| = \sum_{i=1}^n \|\boldsymbol{\mu}_i\|_{TV} = \|\boldsymbol{\mu}\|_{TV},$$

where the next to last equality follows from the uni-directionality of $\boldsymbol{\mu}_i$. By Lemma 2.9, equality holds in (25) if and only if each $\boldsymbol{\nu}_i$ is uni-directional, and it must have the direction of $\boldsymbol{\mu}_i$ since their moments agree. In particular $\boldsymbol{\tau}$ is then unidimensional, hence if in addition S is slender so that $\operatorname{div} \boldsymbol{\tau} = 0$, we get from Lemma 2.10 that equality holds in (24) only when $\boldsymbol{\mu} = \boldsymbol{\nu}$. \square

Note that if S is not slender, $\boldsymbol{\mu}$ and $\boldsymbol{\nu}$ in the previous theorem can be different, even when equality holds in (24). Indeed, for $n = 1$ already, Example 2.1 yields such a situation.

3. MAGNETIZATION-TO-FIELD OPERATORS

$\boxed{\text{MtoFop}}$

Let $S \subset \mathbb{R}^3$ be closed and $Q \subset \mathbb{R}^3 \setminus S$ be compact. For $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ and v a unit vector in \mathbb{R}^3 , the component of the magnetic field $\mathbf{b}(\boldsymbol{\mu})$ in the direction v at $x \notin S$ is given, in view of (10), by

$$\boxed{\text{b3K}} \quad (26) \quad b_v(\boldsymbol{\mu})(x) := v \cdot \mathbf{b}(\boldsymbol{\mu})(x) = -\frac{\mu_0}{4\pi} \int \mathbf{K}_v(x-y) \cdot d\boldsymbol{\mu}(y),$$

where

$$\boxed{\text{noyauK}} \quad (27) \quad \mathbf{K}_v(x) = \frac{v}{|x|^3} - 3x \frac{v \cdot x}{|x|^5} = \operatorname{grad} \left(\frac{v \cdot x}{|x|^3} \right).$$

Consider a finite, positive Borel measure ρ with support contained in Q and let $A : \mathcal{M}(S)^3 \rightarrow L^2(Q, \rho)$ be the operator defined by

$$\boxed{\text{Adef}} \quad (28) \quad A(\boldsymbol{\mu})(x) := b_v(\boldsymbol{\mu})(x), \quad x \in Q.$$

Since \mathbf{K}_v is continuous on $\mathbb{R}^3 \setminus \{0\}$ and Q and S are positively separated, it follows that b_v is continuous on Q and consequently A does indeed map $\mathcal{M}(S)^3$ into $L^2(Q, \rho)$.

If $\Psi \in L^2(Q, \rho)$, then using Fubini's Theorem and (26) we have that

$$(29) \quad \langle \Psi, A(\boldsymbol{\mu}) \rangle_{L^2(Q, \rho)} = -\frac{\mu_0}{4\pi} \iint \Psi(x) \mathbf{K}_v(x-y) \cdot d\boldsymbol{\mu}(y) d\rho(x) = \langle A^*(\Psi), \boldsymbol{\mu} \rangle,$$

where for $x \in S$ the adjoint operator A^* is given by

$$\boxed{\text{b3*W}} \quad (30) \quad A^*(\Psi)(x) := -\frac{\mu_0}{4\pi} \int \Psi(y) \mathbf{K}_v(x-y) d\rho(y), \quad x \in S.$$

In view of (27), a compact way of re-writing (30) is

$$\boxed{\text{b3*Wbis}} \quad (31) \quad A^*(\Psi)(x) := -\mu_0 \operatorname{grad} (\operatorname{grad} U^{\rho, \psi} \cdot v)(x), \quad U^{\rho, \psi}(x) = -\frac{1}{4\pi} \int \frac{\Psi(y)}{|x-y|} d\rho(y).$$

Since Q and S are positively separated it follows as in the proof of Lemma 2.3 that $A^*(\Psi) \in C_0(S)^3$ and thus $A^* : (L^2(Q, \rho))^* \sim L^2(Q, \rho) \rightarrow C_0(S)^3 \subset (\mathcal{M}(S)^3)^*$.

From the point of view of the inverse problem described in Section 1, consisting of recovering $\boldsymbol{\mu}$ from measurements of the field of the potential of its divergence in the direction v , one should think of S as a set *a priori* containing the support of the sources to be recovered, and of Q as the set on which the component of the field in the direction v is measured. In the proposition below, we single out two additional assumptions on the pair (S, Q) , namely:

- (I) $\mathbb{R}^3 \setminus Q$ is connected, $\mathcal{L}_3(Q) = 0$ and $\mathcal{H}_d(S) > 2$.
- (II) $\mathbb{R}^3 \setminus Q$ is connected, $\mathcal{L}_3(Q) = 0$ and there is a smooth complete real analytic surface \mathcal{B} such that Q lies in a single connected component of $\mathbb{R}^3 \setminus \mathcal{B}$, while $\mathcal{H}_d(S \cap \mathcal{B}) > 1$.

Acomp

Proposition 3.1. *Let $S \subset \mathbb{R}^3$ be closed, $Q \subset \mathbb{R}^3 \setminus S$ be compact, ρ be a finite, positive Borel measure with support contained in Q , and v a unit vector in \mathbb{R}^3 .*

- (a) *The operator $A : \mathcal{M}(S)^3 \rightarrow L^2(Q, \rho)$ defined in (28) is compact.*
- (b) *Each function in the range of A^* is the restriction to S of a real-analytic \mathbb{R}^3 -valued function on $\mathbb{R}^3 \setminus Q$.*
- (c) *If either assumption (I) or (II) holds, then A^* is injective, hence A has dense range.*
- (d) *If Q, S, v satisfy the assumptions of Lemma 2.3 and the support of ρ contains $Q \cap \mathcal{A}$, then every element in the kernel of A is S -silent.*

Proof. Let $h := \text{dist}(S, Q) > 0$. Outside an open ball of radius h , the kernel \mathbf{K}_v and its first order derivatives are bounded, say by some constant C . Thus, if $\boldsymbol{\mu}_n$ is a sequence in the unit ball of $\mathcal{M}(S)^3$, then $|b_u(\boldsymbol{\mu}_n)|$ and its partial derivatives are bounded by C on Q . Therefore $b_v(\boldsymbol{\mu}_n)$ is a uniformly bounded family of equicontinuous functions on the compact set Q , hence it is relatively compact in the uniform topology by Ascoli's Theorem. *A fortiori*, this family is relatively compact in $L^2(Q, \rho)$. Besides, since \mathbf{K}_v is a harmonic vector field in $\mathbb{R}^3 \setminus \{0\}$, differentiating (30) under the integral sign shows that the components of $A^*(\Psi)$ are harmonic in $\mathbb{R}^3 \setminus Q$, thus, *a fortiori* real analytic.

To see that A has dense range if either (I) or (II) is satisfied, we prove that A^* is injective in this case. For this, assume that $A^*\Psi = 0$ for some $\Psi \in L^2(Q, \rho)$ and let us show that Ψ is zero ρ -a.e. Assume first that (II) holds, and consider the \mathbb{R}^3 -valued function

$$\mathbf{D}(x) = \frac{1}{4\pi} \int \Psi(y) \frac{x-y}{|x-y|^3} d\rho(y), \quad x \in \mathbb{R}^3 \setminus Q.$$

Arguing as we did to get (12) and observing that $L^2(Q, \rho) \subset L^1(Q, \rho)$ since ρ is finite, we find that \mathbf{D} extends to a locally integrable function on \mathbb{R}^3 with $\text{div } \mathbf{D} = \Psi d\rho$ as distributions. Note that $\text{grad}(\mathbf{D} \cdot v)$ is a harmonic vector field on $\mathbb{R}^3 \setminus Q$ which is equal to $-A^*\Psi/\mu_0$ on S , hence it vanishes there. Since $S \cap \mathcal{B}$ has Hausdorff dimension strictly greater than 1 and \mathcal{B} is real analytic, it holds that $\text{grad}(\mathbf{D} \cdot v)$ vanishes on \mathcal{B} . Moreover, as \mathbf{K}_v vanishes at infinity and Ψ has compact support, $\text{grad}(\mathbf{D} \cdot v)$ vanishes at infinity as well. Thus, by the maximum principle, it must vanish in the component of $\mathbb{R}^3 \setminus \mathcal{B}$ which does not contain Q , therefore also in $\mathbb{R}^3 \setminus Q$ by real analyticity and since it is connected. This means that $\mathbf{D} \cdot v$ is constant in $\mathbb{R}^3 \setminus Q$, and it is in fact identically zero because it is clear from the compactness of Q that \mathbf{D} vanishes at infinity. Now, it holds that $\mathbf{D} = \text{grad}(N * (\psi d\rho))$, where $N(y) = -1/(4\pi|y|)$ is the Newton kernel already used in the proof of Proposition

2.1. Hence, $N * (\psi d\rho)$ must be constant on half lines parallel to v contained in $\mathbb{R}^3 \setminus Q$, and since it vanishes at infinity while Q is compact we find that $N * (\psi d\rho)$ is identically zero in $\mathbb{R}^3 \setminus Q$. Now, being the Newton potential of a finite measure, $N * (\psi d\rho)$ is a locally integrable function and, since $\mathcal{L}_3(Q) = 0$, we just showed that it is zero almost everywhere. Hence it is the zero distribution, and so is its weak Laplacian $\psi d\rho$. Consequently ψ is zero ρ -a.e., as desired. If (I) holds instead of (II) the proof of (c) is similar but easier, because we conclude directly that $\text{grad}(\mathbf{D} \cdot v) = 0$ on $\mathbb{R}^3 \setminus Q$, since it is harmonic there and vanishes on S which has Hausdorff dimension strictly greater than 2. Finally, to prove (d), observe that if $A(\boldsymbol{\mu}) = 0$ a.e. with respect to ρ , then by continuity $A(\boldsymbol{\mu}) = 0$ on the support of ρ and so on $Q \cap \mathcal{A}$. Thus, by Lemma 2.3 $\boldsymbol{\mu}$ is S -silent whenever $\boldsymbol{\mu}$ is in the kernel of A . \square

rmkdens

Remark. Note that if Q contains a nonempty open ball B (then of course $\mathcal{L}_3(Q) \neq 0$ and neither (I) nor (II) is satisfied), and if moreover the support of ρ contains B , then the image of A is not dense in $L^2(Q, \rho)$ since it consists of functions harmonic in B .

The density of $\text{ran}A$ plays no role in the forthcoming results, so from the sheer mathematical point of view we may forget about Proposition 3.1 (c). However, it is more often satisfied than not. For instance, in paleomagnetism, Q would be a rectangular region in a plane and S a rock sample which is either volumic (then (I) is met) or sanded down to a thin slab (then (II) is met, with \mathcal{B} a plane). In practice, if A has dense range then the data can be explained arbitrary well in terms of similarity between the measured and modeled fields, at the cost of proposing a model for $\boldsymbol{\mu}$ which is unreasonably large and therefore non-physical, see Theorem 4.3 point (b). This phenomenon is typical of ill-posed problems and a compelling reason why regularization is needed. That is, one must trade-off between the precision of the model against available data and its physical relevance, *e.g.* the regularization parameter λ should not be made too small in Theorem 4.3.

4. REGULARIZATION BY PENALIZING THE TOTAL VARIATION

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In this section, we consider the inverse magnetization problem of recovering $\boldsymbol{\mu} \in \mathcal{M}(S)^3$ from the knowledge of $A(\boldsymbol{\mu})$, where A is the operator defined in (28). We will study the regularization scheme EP-2, based on (6), that penalizes the total variation of the candidate approximant, and prove that solutions to EP-2 exist and are necessarily “localized”, in the sense that their support has dimension at most 2 if S has nonempty interior in \mathbb{R}^3 and dimension at most 1 if S is contained in some unbounded analytic surface where it has nonempty interior. The existence of a solution to EP-2, as well as the optimality condition given in Theorem 4.1, fall under the scope of [13, prop. 3.6] and could just have been referenced. We nevertheless provide a proof, partly because it may be interesting in its own right as it is independent from the Fenchel duality used in [13], but mainly because we want to discuss non-uniqueness in a specific manner. We conclude this section with a ‘consistency’ result showing that solutions to EP-2 approach those of EP-1, in the limit that the regularization parameter λ and the (additive) perturbation on the data vanish in a controlled manner. Our account of this regularization theory is new inasmuch as it includes the asymptotic behavior of total variation measures of the solutions, and deals with narrow convergence (not just weak-*).

Hereafter, as in Section 3, we let $S \subset \mathbb{R}^3$ be closed, $Q \subset \mathbb{R}^3 \setminus S$ be compact, ρ be a finite, positive Borel measure supported in Q , and v a unit vector in \mathbb{R}^3 . The operator A is then defined by (28). For $\boldsymbol{\mu} \in \mathcal{M}(S)^3$, $f \in L^2(Q, \rho)$, and $\lambda > 0$, we recall from (5) the definition of $\mathcal{F}_{f, \lambda}$:

$$\boxed{\text{defcrit}} \quad (32) \quad \mathcal{F}_{f, \lambda}(\boldsymbol{\mu}) := \|f - A\boldsymbol{\mu}\|_{L^2(Q, \rho)}^2 + \lambda \|\boldsymbol{\mu}\|_{TV},$$

and from (6) that $\boldsymbol{\mu}_\lambda \in \mathcal{M}(S)^3$ denotes a minimizer of $\mathcal{F}_{f, \lambda}$ whose existence is proved in Theorem 4.1 below; *i.e.*,

$$\boxed{\text{crit}} \quad (33) \quad \mathcal{F}_{f, \lambda}(\boldsymbol{\mu}_\lambda) = \inf_{\boldsymbol{\mu} \in \mathcal{M}(S)^3} \mathcal{F}_{f, \lambda}(\boldsymbol{\mu}).$$

$\boxed{\text{VPTm}}$ **Theorem 4.1.** *Notation and assumptions being as above, given $f \in L^2(Q, \rho)$ there is a solution to (33). A \mathbb{R}^3 -valued measure $\boldsymbol{\mu}_\lambda \in \mathcal{M}(S)^3$ is such a solution if and only if:*

$$\boxed{\text{CP}} \quad (34) \quad \begin{aligned} A^*(f - A\boldsymbol{\mu}_\lambda) &= \frac{\lambda}{2} \mathbf{u}_{\boldsymbol{\mu}_\lambda} \quad |\boldsymbol{\mu}_\lambda|\text{-a.e. and} \\ |A^*(f - A\boldsymbol{\mu}_\lambda)| &\leq \frac{\lambda}{2} \quad \text{everywhere on } S. \end{aligned}$$

Moreover, $\boldsymbol{\mu}'_\lambda \in \mathcal{M}(S)^3$ is another solution if and only if:

- (a) $A(\boldsymbol{\mu}'_\lambda - \boldsymbol{\mu}_\lambda) = 0$,
- (b) *there is a $|\boldsymbol{\mu}_\lambda|$ -measurable non-negative function g and a positive measure $\nu_s \in \mathcal{M}(S)$, singular to $|\boldsymbol{\mu}_\lambda|$ and supported on the set $\{x \in S : |A^*(f - A\boldsymbol{\mu}_\lambda)(x)| = \lambda/2\}$, such that*

$$\boxed{\text{equivsoldec}} \quad (35) \quad d\boldsymbol{\mu}'_\lambda = g d\boldsymbol{\mu}_\lambda + 2 \frac{A^*(f - A\boldsymbol{\mu}_\lambda)}{\lambda} d\nu_s.$$

Proof. Fix $\lambda > 0$ and let $\boldsymbol{\mu}_n$ a minimizing sequence for the right hand side of (33). By construction $\|\boldsymbol{\mu}_n\|_{TV}$ is bounded, hence we can find a subsequence that converges weak-* to some $\boldsymbol{\mu}_\lambda$, by the Banach-Alaoglu Theorem. Renumbering if necessary, let us denote this subsequence by $\boldsymbol{\mu}_n$ again. The Banach-Alaoglu Theorem also entails that

$$\boxed{\text{var1}} \quad (36) \quad \|\boldsymbol{\mu}_\lambda\|_{TV} \leq \liminf_n \|\boldsymbol{\mu}_n\|_{TV}.$$

Moreover, since A is compact, $f - A(\boldsymbol{\mu}_n)$ converges to $f - A(\boldsymbol{\mu}_\lambda)$ in $L^2(Q, \rho)$, hence

$$\boxed{\text{convL2}} \quad (37) \quad \|f - A(\boldsymbol{\mu}_\lambda)\|_{L^2(Q, \rho)} = \lim_n \|f - A(\boldsymbol{\mu}_n)\|_{L^2(Q, \rho)}.$$

Because $\boldsymbol{\mu}_n$ is minimizing, it now follows from (36) and (37) that $\boldsymbol{\mu}_\lambda$ meets (33) and that (36) is both an equality and a true limit.

Let now $\boldsymbol{\nu} \in \mathcal{M}(S)^3$ be absolutely continuous with respect to $|\boldsymbol{\mu}_\lambda|$ with Radon-Nykodim derivative $\boldsymbol{\nu}_a \in (L^1(\boldsymbol{\mu}_\lambda))^3$; that is to say: $d\boldsymbol{\nu} = \boldsymbol{\nu}_a d|\boldsymbol{\mu}_\lambda|$.

We evaluate $\mathcal{F}_{f, \lambda}(\boldsymbol{\mu}_\lambda + t\boldsymbol{\nu})$ for small t . On the one hand,

$$\boxed{\text{expq}} \quad (38) \quad \|f - A(\boldsymbol{\mu}_\lambda + t\boldsymbol{\nu})\|_{L^2(Q, \rho)}^2 = \|f - A(\boldsymbol{\mu}_\lambda)\|_{L^2(Q, \rho)}^2 - 2t \langle f - A(\boldsymbol{\mu}_\lambda), A(\boldsymbol{\nu}) \rangle + t^2 \|A(\boldsymbol{\nu})\|_{L^2(Q, \rho)}^2.$$

On the other hand, since it has unit norm $|\boldsymbol{\mu}_\lambda|$ -a.e., the Radon-Nykodim derivative $\mathbf{u}_{\boldsymbol{\mu}_\lambda}$ has a unique norming functional when viewed as an element of $(L^1(\boldsymbol{\mu}_\lambda))^3$, given by

$$\boldsymbol{\Psi} \mapsto \int \boldsymbol{\Psi} \cdot \mathbf{u}_{\boldsymbol{\mu}_\lambda} d|\boldsymbol{\mu}_\lambda|, \quad \boldsymbol{\Psi} \in (L^1(\boldsymbol{\mu}_\lambda))^3.$$

Hence, the $(L^1(\mu_\lambda))^3$ -norm is Gâteaux differentiable at \mathbf{u}_{μ_λ} [9, Part 3, Ch. 1, Prop. 2, Remark 1] and we get that

$$\boxed{\text{Gateaux}} \quad (39) \quad \|\mu_\lambda + t\nu\|_{TV} = \int |\mathbf{u}_{\mu_\lambda} + t\nu_a| d|\mu_\lambda| = \|\mu_\lambda\|_{TV} + t \int \nu_a \cdot \mathbf{u}_{\mu_\lambda} d|\mu_\lambda| + t\varepsilon(t),$$

where $\varepsilon(t) \rightarrow 0$ when $t \rightarrow 0$. From (38) and (39), we gather that

$$\mathcal{F}_{f,\lambda}(\mu_\lambda + t\nu) - \mathcal{F}_{f,\lambda}(\mu_\lambda) = -2t\langle A^*(f - A(\mu_\lambda)), \nu \rangle + t\lambda \int \nu_a \cdot \mathbf{u}_{\mu_\lambda} d|\mu_\lambda| + o(t).$$

The left hand side is nonnegative by definition of μ_λ , so the coefficient of t in the right hand side is zero otherwise we could adjust the sign for small $|t|$. Consequently

$$\int (-2A^*(f - A(\mu_\lambda)) + \lambda \mathbf{u}_{\mu_\lambda}) \cdot \nu_a d|\mu_\lambda| = 0, \quad \nu_a \in L^1(|\mu_\lambda|),$$

which implies the first equation in (34).

Assume next that the second inequality in (34) is violated:

$$\boxed{\text{contsi}} \quad (40) \quad |A^*(f - A\mu_\lambda)|(x) > \lambda/2$$

for some $x \in S$. Then $|\mu_\lambda|(\{x\}) = 0$ by the first part of (34) just proven, and the measure

$$\boxed{\text{defdels}} \quad (41) \quad \nu = \frac{A^*(f - A\mu_\lambda)(x)}{|A^*(f - A\mu_\lambda)|(x)} \delta_x,$$

is singular with respect to μ_λ . Hence, for $t > 0$,

$$\boxed{\text{tvsp}} \quad (42) \quad \|\mu_\lambda + t\nu\|_{TV} = \|\mu_\lambda\|_{TV} + t\|\nu\|_{TV} = \|\mu_\lambda\|_{TV} + t,$$

and it follows from (38), (41) and (42) that

$$\mathcal{F}_{f,\lambda}(\mu_\lambda + t\nu) - \mathcal{F}_{f,\lambda}(\mu_\lambda) = -2t|A^*(f - A\mu_\lambda)|(x) + t\lambda + O(t^2)$$

which is strictly negative for $t > 0$ small enough, in view of (40). But this cannot hold since μ_λ is a minimizer of (32), thereby proving the second inequality in (34) by contradiction.

Conversely, assume that (34) holds. Let $\nu \in \mathcal{M}(S)^3$ and write the Radon-Nykodim decomposition of ν with respect to μ_λ as $d\nu = \nu_a d|\mu_\lambda| + d\nu_s$, where $\nu_a \in L^1(|\mu_\lambda|)$ and ν_s is singular with respect to $|\mu_\lambda|$. Setting $t = 1$ in (38), we get that

$$\boxed{\text{sufopt}} \quad (43) \quad \begin{aligned} \|f - A(\mu_\lambda + \nu)\|_{L^2(Q,\rho)}^2 &- \|f - A(\mu_\lambda)\|_{L^2(Q,\rho)}^2 \geq -2\langle f - A(\mu_\lambda), A(\nu) \rangle \\ &= -2 \int A^*(f - A(\mu_\lambda)) \cdot \nu_a d|\mu_\lambda| - 2\langle A^*(f - A(\mu_\lambda)), \nu_s \rangle \\ &= -\lambda \int (\nu_a \cdot \mathbf{u}_{\mu_\lambda}) d|\mu_\lambda| - 2\langle A^*(f - A(\mu_\lambda)), \nu_s \rangle \\ &\geq -\lambda \int (\nu_a \cdot \mathbf{u}_{\mu_\lambda}) d|\mu_\lambda| - \lambda \|\nu_s\|_{TV}. \end{aligned}$$

In another connection, we have that

$$\|\mu_\lambda + \nu\|_{TV} = \int |\mathbf{u}_{\mu_\lambda} + \nu_a| d|\mu_\lambda| + \|\nu_s\|_{TV} = \int (1 + 2\nu_a \cdot \mathbf{u}_{\mu_\lambda} + |\nu_a|^2)^{1/2} d|\mu_\lambda| + \|\nu_s\|_{TV},$$

and since $|\mathbf{u}_{\mu_\lambda}| = 1$ a.e. with respect to $|\mu_\lambda|$, we obtain:

$$\boxed{\text{locm}} \quad (44) \quad (1 + 2\nu_a \cdot \mathbf{u}_{\mu_\lambda} + |\nu_a|^2)^{1/2} \geq |1 + \nu_a \cdot \mathbf{u}_{\mu_\lambda}|, \quad |\mu_\lambda|\text{-a.e.}$$

Thus, if we let E_+ (resp. E_-) be the subset of $\text{supp } |\boldsymbol{\mu}_\lambda|$ where $\mathbf{u}_{\boldsymbol{\mu}_\lambda} \cdot \boldsymbol{\nu}_a > -1$ (resp. ≤ -1), we obtain:

$$\boxed{\text{1dec}} \quad (45) \quad \lambda \|\boldsymbol{\mu}_\lambda + \boldsymbol{\nu}\|_{TV} \geq \lambda \int_{E_+} (1 + \boldsymbol{\nu}_a \cdot \mathbf{u}_{\boldsymbol{\mu}_\lambda}) d|\boldsymbol{\mu}_\lambda| + \lambda \|\boldsymbol{\nu}_s\|_{TV},$$

Besides, it follows from (43) that

$$\boxed{\text{2dec}} \quad (46) \quad \|f - A(\boldsymbol{\mu}_\lambda + \boldsymbol{\nu})\|_{L^2(Q, \rho)}^2 - \|f - A(\boldsymbol{\mu}_\lambda)\|_{L^2(Q, \rho)}^2 \geq -\lambda \int_{E_+} (\boldsymbol{\nu}_a \cdot \mathbf{u}_{\boldsymbol{\mu}_\lambda}) d|\boldsymbol{\mu}_\lambda| + \lambda \int_{E_-} d|\boldsymbol{\mu}_\lambda| - \lambda \|\boldsymbol{\nu}_s\|_{TV}.$$

Adding up (45) and (46), using that $\|\boldsymbol{\mu}_\lambda\|_{TV} = \int_{E_+} d|\boldsymbol{\mu}_\lambda| + \int_{E_-} d|\boldsymbol{\mu}_\lambda|$, we obtain:

$$\boxed{\text{optmeas}} \quad (47) \quad \mathcal{F}_{f, \lambda}(\boldsymbol{\mu}_\lambda + \boldsymbol{\nu}) - \mathcal{F}_{f, \lambda}(\boldsymbol{\mu}_\lambda) \geq 0,$$

thereby showing that $\boldsymbol{\mu}_\lambda$ indeed meets (33).

Finally, observe that in the previous estimates we neglected the term $t^2 \|A\boldsymbol{\nu}\|_{L^2(Q, \rho)}^2$ in (38) and the term $|\boldsymbol{\nu}_a|^2 - (\boldsymbol{\nu}_a \cdot \mathbf{u}_{\boldsymbol{\mu}_\lambda})^2$ in (44), as well as the term $\lambda \int_{E_-} (|\boldsymbol{\nu}_a \cdot \mathbf{u}_{\boldsymbol{\mu}_\lambda}| - 1) d|\boldsymbol{\mu}_\lambda|$ in (45) and (46), along with the term $\lambda \|\boldsymbol{\nu}_s\|_{TV} - 2\langle A^*(f - A(\boldsymbol{\mu}_\lambda)), \boldsymbol{\nu}_s \rangle$ in (43). Hence, equality holds in (47) if and only if they are all zero. Thus, for $\boldsymbol{\mu}'_\lambda = \boldsymbol{\mu}_\lambda + \boldsymbol{\nu}$ to be another solution to (33), it is necessary and sufficient that $A\boldsymbol{\nu} = 0$ and $\boldsymbol{\nu}_a = h\mathbf{u}_{\boldsymbol{\mu}_\lambda}$ with h a real-valued function such that $h \geq -1$, a.e. with respect to $|\boldsymbol{\mu}_\lambda|$, while $\boldsymbol{\nu}_s$ is supported on the subset of S where $|A^*(f - A\boldsymbol{\mu}_\lambda)| = \lambda/2$ and $A^*(f - A\boldsymbol{\mu}_\lambda) = (\lambda/2)\mathbf{u}_{\boldsymbol{\nu}_s}$ at $|\boldsymbol{\nu}_s|$ -a.e. point. Thus, $\boldsymbol{\mu}'_\lambda = g\boldsymbol{\mu}_\lambda + \boldsymbol{\nu}_s$ with $g := 1 + h \geq 0$, which gives us (a) and (b). \square

That any two minimizers of (32) must differ by a member of the kernel of A is but a simple consequence of the strict convexity of the $L^2(Q, \rho)$ -norm. In particular, if the assumptions on Q, S, \mathcal{A} and v of Lemma 2.3 hold and the support of ρ contains $Q \cap \mathcal{A}$, then any two minimizers are S -equivalent, by (d) of Proposition 3.1. The second assertion of Theorem 4.1 means that when (a) holds, then $\|\boldsymbol{\mu}'_\lambda\|_{TV} = \|\boldsymbol{\mu}_\lambda\|_{TV}$ if and only if (b) holds.

$\boxed{\text{corspar}}$ **Corollary 4.2.** *Assumptions and notation being as in Theorem 4.1, assume in addition that S is contained in the unbounded connected component of $\mathbb{R}^3 \setminus S$. Then, the union of the supports of all minimizers of (32), for fixed f and $\lambda > 0$, is contained in a finite union of points, embedded curves and surfaces, each of which is real-analytic and bounded. If \mathcal{B} (resp. \mathcal{C}) is an unbounded connected real analytic surface (resp. curve) such that $\mathcal{B} \cap Q$ does not disconnect \mathcal{B} (resp. $\mathcal{C} \cap Q = \emptyset$), then the aforementioned union of supports has an intersection with \mathcal{B} (resp. \mathcal{C}) which is contained in a finite union of points and embedded real analytic curves (resp. points).*

Proof. Recall from (b) in Proposition 3.1 that $A^*(f - A\boldsymbol{\mu}_\lambda)$ is the restriction to S of a \mathbb{R}^3 -valued real analytic vector field on $\mathbb{R}^3 \setminus Q$ that vanishes at infinity. Set $g = |A^*(f - A\boldsymbol{\mu}_\lambda)|^2$ which vanishes at infinity and is a real analytic function $\mathbb{R}^3 \setminus Q \rightarrow \mathbb{R}$. Theorem 4.1 implies that the support of $|\boldsymbol{\mu}_\lambda|$, and also of any other minimizer of (32), is included in the zero set of $h := g - \lambda^2/4$. Note that h is independent of the minimizer $\boldsymbol{\mu}_\lambda$ under consideration, since any two have the same image under A by Theorem 4.1. Now, since h is not the zero

function on the unbounded component of $\mathbb{R}^3 \setminus Q$ (because g vanishes at infinity), its zero set is a locally finite union of points and real analytic embedded curves and surfaces, see discussion after Theorem 2.2. Moreover, since h tends to $-\lambda^2/4$ at infinity, the zero set of h intersected with S is contained in a relatively compact open subset of $\mathbb{R}^3 \setminus Q$, therefore the support of μ_λ or any other minimizer is contained in finitely many of these points, curves and surfaces.

The proof of the second assertion is similar, reasoning in $\mathcal{B} \setminus Q$ (resp. \mathcal{C} rather than $\mathbb{R}^3 \setminus Q$). \square

Remark. Corollary 4.2 applies for instance in paleomagnetism, when trying to recover magnetizations on thin slabs of rock *via* the regularization scheme (33), a case in which \mathcal{B} is a plane. Note that if we omit the assumption that \mathcal{B} (resp. \mathcal{C}) is unbounded in Corollary 4.2, then we can only conclude that the support of μ_λ either is contained in finitely many points and arcs (resp. points) or else $h = 0$ on \mathcal{B} (resp. \mathcal{C}). This remark applies, *e.g.* in MEG inverse problems where \mathcal{B} is typically a closed surface.

Even if $f \in \text{range } A$, say $f = A(\mu_0)$ for some $\mu_0 \in (\mathcal{M}(S))^3$, it is clear from (34) that $\mu_\lambda \neq \mu_0$ when $\lambda > 0$, unless $\mu_0 = 0$. The purpose of the regularizing term $\lambda \|\mu\|_{TV}$ in (32) is rather to get a μ_λ which is not too far from μ_0 when f gets replaced by $f_e = f + e$ in (33). Here, e is some error (*e.g.* due to measurements) and f_e represents the actual data. To clarify the matter, whenever $f, e \in L^2(Q, \rho)$ we set $f_e := f + e$ and, for $\lambda > 0$, we let $\mu_{\lambda, e}$ be a minimizer of (32) when f gets replaced by f_e . Thus, with the notation of (33), we have that $\mu_\lambda = \mu_{\lambda, 0}$. Typical results to warrant a regularization approach based on approximating μ_0 by $\mu_{\lambda, e}$ are of “consistency” type, namely they assert that $\mu_{\lambda, e}$ yields information on μ_0 as $\|e\|_{L^2(Q, \rho)}$ and λ go to 0 in a combined fashion, see for example [14, Thms. 2&5] or [31, Thm. 3.5&4.4]. We give below a theorem of this type, which goes beyond [31, Thm. 3.5] in that we deal not just with weak- $*$ convergence of subsequences μ_{λ_n, e_n} , but more generally with narrow convergence of both μ_{λ_n, e_n} and $|\mu_{\lambda_n, e_n}|$. We will not consider quantitative convergence properties involving the Bregman distance, that require an additional source condition which needs not be satisfied here in general.

As an extra piece of notation, we define for $\mu_0 \in \mathcal{M}(S)^3$:

defineqA

$$(48) \quad \mathfrak{M}(\mu_0) := \min\{\|\mu\|_{TV} : A(\mu) = A(\mu_0)\}.$$

The infimum in the right-hand side of (48) is indeed attained, by the Banach-Alaoglu theorem and since the kernel of A is weak- $*$ closed. When S and Q satisfy the conditions of Lemma 2.3, then this kernel consists of S -silent magnetizations and $\mathfrak{M}(\mu_0)$ is just $M(\mu_0)$ defined in (4). But when these conditions are not satisfied (for instance if S is smooth compact surface), then the two quantities may not coincide.

Recall that a sequence $\mu_n \in \mathcal{M}(S)^3$ converges in the narrow sense to $\mu \in \mathcal{M}(S)^3$ if $\int \varphi \cdot d\mu_n \rightarrow \int \varphi \cdot d\mu$ as $n \rightarrow \infty$, whenever $\varphi : S \rightarrow \mathbb{R}^3$ is continuous and bounded. When S is compact this is equivalent to weak- $*$ convergence, but if S is unbounded it means that μ_n does not “lose mass at infinity”.

lamto0Thm

Theorem 4.3. *Assumptions and notation being as in Theorem 4.1, given $f \in L^2(Q, \rho)$, the following hold.*

lamto0a (a) If $f = A(\boldsymbol{\mu}_0)$ with $\boldsymbol{\mu}_0 \in (\mathcal{M}(S))^3$, while $e \in L^2(Q, \rho)$ and $\lambda > 0$, then

$$\mathbf{regbn} \quad (49) \quad \|\boldsymbol{\mu}_{\lambda,e}\|_{TV} \leq \frac{\|e\|_{L^2(Q,\rho)}^2}{\lambda} + \mathfrak{M}(\boldsymbol{\mu}_0)$$

and

$$\mathbf{limTVa} \quad (50) \quad \lim_{\lambda \rightarrow 0^+, \|e\|_{L^2(Q,\rho)}/\sqrt{\lambda} \rightarrow 0} \|\boldsymbol{\mu}_{\lambda,e}\|_{TV} = \mathfrak{M}(\boldsymbol{\mu}_0).$$

As $\lambda \rightarrow 0$ and $\|e\|_{L^2(Q,\rho)}/\sqrt{\lambda} \rightarrow 0$, any weak-* cluster point $\boldsymbol{\mu}^*$ of $\boldsymbol{\mu}_{\lambda,e}$ (there must be at least one since $\|\boldsymbol{\mu}_{\lambda,e}\|_{TV}$ is bounded) meets $A(\boldsymbol{\mu}^*) = A(\boldsymbol{\mu}_0) = f$ and satisfies:

$$\mathbf{minwc} \quad (51) \quad \|\boldsymbol{\mu}^*\|_{TV} = \mathfrak{M}(\boldsymbol{\mu}_0).$$

Moreover, if $\lambda_n \rightarrow 0^+$ and $\|e_n\|_{L^2(Q,\rho)}/\sqrt{\lambda_n} \rightarrow 0$, with λ_n, e_n such that $\boldsymbol{\mu}_{\lambda_n,e_n}$ converges weak-* to $\boldsymbol{\mu}^*$, we have that

$$\mathbf{limfinps} \quad (52) \quad \lim_{n \rightarrow \infty} \int \left| \frac{2A^*(f_{e_n} - A\boldsymbol{\mu}_{\lambda_n,e_n})}{\lambda_n} - \mathbf{u}_{\boldsymbol{\mu}^*} \right| d|\boldsymbol{\mu}^*| = 0,$$

also $\boldsymbol{\mu}_{\lambda_n,e_n}$ and $|\boldsymbol{\mu}_{\lambda_n,e_n}|$ converge respectively to $\boldsymbol{\mu}^*$ and $|\boldsymbol{\mu}^*|$ in the narrow sense.

- lamto0b** (b) If $f \notin \text{range } A$ and either assumption (I) or (II) in Section 3 holds, then $\|\boldsymbol{\mu}_{\lambda,e}\|_{TV} \rightarrow \infty$ as $\lambda \rightarrow 0$ and $e \rightarrow 0$.
- (c) If $\lambda \geq 2 \sup_{x \in S} |(A^*f)(x)|$, then the unique minimizer of the right-hand side of (4) is the zero magnetization.

Proof. If $f \in \text{range } A$, or if $\text{range } A$ is dense in $L^2(Q, \rho)$, it is clear that $\mathcal{F}_{f_e,\lambda}(\boldsymbol{\mu}_{\lambda,e}) \rightarrow 0$ as $\lambda \rightarrow 0$ and $e \rightarrow 0$, hence $\|f - A(\boldsymbol{\mu}_{\lambda,e})\|_{L^2(Q,\rho)} \rightarrow 0$ in this case. In particular, if $f \notin \text{range } A$ but either assumption (I) or (II) in Section 3 holds, then $\text{range } A$ is dense by Proposition 3.1 (c) and so $\|\boldsymbol{\mu}_{\lambda,e}\|_{TV} \rightarrow \infty$ otherwise a subsequence would converge weak-* to some $\boldsymbol{\mu}_0 \in (\mathcal{M}(S))^3$ implying in the limit that $f = A(\boldsymbol{\mu}_0)$, a contradiction which proves (b).

Next, let $\tilde{\boldsymbol{\mu}}_0$ be a minimizer of the right hand side of (4), so that $A(\tilde{\boldsymbol{\mu}}_0) = f$ and $\|\tilde{\boldsymbol{\mu}}_0\|_{TV} = \mathfrak{M}(\boldsymbol{\mu}_0)$. By the optimality of $\boldsymbol{\mu}_{\lambda,e}$, we have that

$$\mathbf{ptdep} \quad (53) \quad \begin{aligned} & \|f_e - A(\boldsymbol{\mu}_{\lambda,e})\|_{L^2(Q)}^2 + \lambda \|\boldsymbol{\mu}_{\lambda,e}\|_{TV} = \mathcal{F}_{f_e,\lambda}(\boldsymbol{\mu}_{\lambda,e}) \leq \mathcal{F}_{f_e,\lambda}(\tilde{\boldsymbol{\mu}}_0) \\ & = \|e\|_{L^2(Q,\rho)}^2 + \lambda \|\tilde{\boldsymbol{\mu}}_0\|_{TV} = \|e\|_{L^2(Q,\rho)}^2 + \lambda \mathfrak{M}(\boldsymbol{\mu}_0), \end{aligned}$$

implying that (49) holds. Thus, if $\boldsymbol{\mu}^*$ is a weak-* cluster point of $\{\boldsymbol{\mu}_{\lambda,e}\}$ as $\lambda \rightarrow 0^+$ with $\|e\|_{L^2(Q,\rho)} = o(\sqrt{\lambda})$, and if λ_n, e_n are sequences with these limiting properties such that $\boldsymbol{\mu}_{\lambda_n,e_n}$ converges weak-* to $\boldsymbol{\mu}^*$, we deduce from the Banach-Alaoglu Theorem that

$$\mathbf{prelim} \quad (54) \quad \|\boldsymbol{\mu}^*\|_{TV} \leq \liminf_n \|\boldsymbol{\mu}_{\lambda_n,e_n}\|_{TV} \leq \limsup_n \|\boldsymbol{\mu}_{\lambda_n,e_n}\|_{TV} \leq \mathfrak{M}(\boldsymbol{\mu}_0).$$

Also, since A is weak-* to weak continuous (it is even compact), we get that

$$\mathbf{seclim} \quad (55) \quad \|f - A(\boldsymbol{\mu}^*)\|_{L^2(Q,\rho)} \leq \lim_n \|f - A(\boldsymbol{\mu}_{\lambda_n,e_n})\|_{L^2(Q,\rho)} = 0,$$

where the last equality was obtained in the proof of (b). From (55) it follows that $A(\boldsymbol{\mu}^*) = f$, and from (54) we now see that (51) holds, by definition of $\mathfrak{M}(\boldsymbol{\mu}_0)$. Moreover, since a

weak-* convergent subsequence can be extracted from any subsequence of $\boldsymbol{\mu}_{\lambda_n, e_n}$, we deduce from what precedes that (50) takes place. Next, if λ_n, e_n are as before, we get in view of (50) and (53) that $\|f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n})\|_{L^2(Q, \rho)}^2 = o(\lambda_n)$, which is equivalent to

$$0 = \lim_n \frac{2}{\lambda_n} \langle f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}), f - A(\boldsymbol{\mu}_{\lambda_n, e_n}) \rangle + \frac{2}{\lambda_n} \langle f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}), e_n \rangle,$$

and since $\|e_n\|_{L^2(Q, \rho)} = o(\sqrt{\lambda_n})$ while $f = A(\boldsymbol{\mu}^*)$, we obtain:

$$\begin{aligned} (56) \quad & 0 = \lim_n \left\langle \frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n}, \boldsymbol{\mu}_{\lambda_n, e_n} - \boldsymbol{\mu}^* \right\rangle = \lim_n \left\langle \frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n}, \boldsymbol{\mu}_{\lambda_n, e_n} \right\rangle \\ & - \lim_n \left\langle \frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n}, \boldsymbol{\mu}^* \right\rangle = \lim_n \left(\|\boldsymbol{\mu}_{\lambda_n, e_n}\|_{TV} - \int \frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n} \cdot d\boldsymbol{\mu}^* \right) \\ & = \|\boldsymbol{\mu}^*\|_{TV} - \lim_n \int \frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n} \cdot \mathbf{u}_{\boldsymbol{\mu}^*} d|\boldsymbol{\mu}^*|, \end{aligned}$$

where we used the first relation in (34) to get the third equality and (50), (51) to get the last one. By the second relation in (34), we know that $|2A^*(f_{e_n} - A\boldsymbol{\mu}_{\lambda_n})/\lambda_n| \leq 1$ everywhere on S , hence (56) implies that for any $\varepsilon > 0$

$$\limsup_n |\boldsymbol{\mu}^*| \left\{ x \in S : \frac{2A^*(f_{e_n} - A\boldsymbol{\mu}_{\lambda_n, e_n})}{\lambda_n} \cdot \mathbf{u}_{\boldsymbol{\mu}^*} < 1 - \varepsilon \right\} = 0.$$

Therefore, using the Borel-Cantelli Lemma and a diagonal argument, we may extract a subsequence λ_{k_n} for which $2A^*(f_{e_{k_n}} - A\boldsymbol{\mu}_{\lambda_{k_n}, e_{k_n}})/\lambda_{k_n}$ converges pointwise $|\boldsymbol{\mu}^*|$ -a.e. to $\mathbf{u}_{\boldsymbol{\mu}^*}$. So, by dominated convergence, it follows that

$$\lim_{n \rightarrow \infty} \int \left| \frac{2A^*(f_{e_{k_n}} - A(\boldsymbol{\mu}_{\lambda_{k_n}, e_{k_n}}))}{\lambda_{k_n}} - \mathbf{u}_{\boldsymbol{\mu}^*} \right| d|\boldsymbol{\mu}^*| = 0,$$

and since the reasoning can be applied to any subsequence of λ_n , we obtain (52).

We now prove that $|\boldsymbol{\mu}_{\lambda_n, e_n}|$ converges weak-* to $|\boldsymbol{\mu}^*|$. For this, it is enough to show that if $|\boldsymbol{\mu}_{\lambda_n, e_n}|$ converges weak-* to $\boldsymbol{\nu} \geq 0 \in \mathcal{M}(S)$, then $\boldsymbol{\nu} = |\boldsymbol{\mu}^*|$. For this, let $\psi : S \rightarrow [0, 1]$ be a continuous function with compact support. Pick $\varepsilon > 0$, and then n_ε such that the integral in the left-hand side of (52) is less than ε for $n \geq n_\varepsilon$. As $|2A^*(f_{e_{n_\varepsilon}} - A(\boldsymbol{\mu}_{\lambda_{n_\varepsilon}, e_{n_\varepsilon}}))/\lambda_{n_\varepsilon}| \leq 1$ everywhere on S by (34), we obtain from the definition of n_ε and (51) that

$$\begin{aligned} \int \psi d\boldsymbol{\nu} &= \lim_n \int \psi d|\boldsymbol{\mu}_{\lambda_n, e_n}| \geq \lim_n \left| \int \psi \frac{2A^*(f_{e_{n_\varepsilon}} - A(\boldsymbol{\mu}_{\lambda_{n_\varepsilon}, e_{n_\varepsilon}}))}{\lambda_{n_\varepsilon}} \cdot d\boldsymbol{\mu}_{\lambda_n, e_n} \right| \\ &= \left| \int \psi \frac{2A^*(f_{e_{n_\varepsilon}} - A(\boldsymbol{\mu}_{\lambda_{n_\varepsilon}, e_{n_\varepsilon}}))}{\lambda_{n_\varepsilon}} \cdot d\boldsymbol{\mu}^* \right| \geq \int \psi d|\boldsymbol{\mu}^*| - \int \left| \frac{2A^*(f_{e_{n_\varepsilon}} - A(\boldsymbol{\mu}_{\lambda_{n_\varepsilon}, e_{n_\varepsilon}}))}{\lambda_{n_\varepsilon}} - \mathbf{u}_{\boldsymbol{\mu}^*} \right| \cdot d|\boldsymbol{\mu}^*| \\ &\geq \int \psi d|\boldsymbol{\mu}^*| - \varepsilon, \end{aligned}$$

where we used in the equality above that $A^*(f_{e_{n_\varepsilon}} - A(\boldsymbol{\mu}_{\lambda_{n_\varepsilon}, e_{n_\varepsilon}}))$ is continuous on S , by Proposition 3.1 (b). Since $\varepsilon > 0$ was arbitrary, we conclude that $\boldsymbol{\nu} - |\boldsymbol{\mu}^*| \geq 0$. However, since $\|\boldsymbol{\mu}^*\|_{TV} = \mathfrak{M}(\boldsymbol{\mu}_0)$ by (51), whereas $\|\boldsymbol{\nu}\|_{TV} \leq \mathfrak{M}(\boldsymbol{\mu}_0)$ by the Banach Alaoglu theorem, we conclude that $\boldsymbol{\nu} - |\boldsymbol{\mu}^*|$ is the zero measure, as desired.

To establish that $\boldsymbol{\mu}_{\lambda_n, e_n}$ converges to $\boldsymbol{\mu}^*$ in the narrow sense, pick $\varepsilon > 0$ and n_ε as before. Fix R_ε so large that $|\boldsymbol{\mu}^*|(S \cap \overline{B}(0, R_\varepsilon)) > \|\boldsymbol{\mu}^*\|_{TV} - \varepsilon$ and for each R let $\psi_R : S \rightarrow [0, 1]$ be

continuous, identically 1 on $S \cap B(0, R)$ and 0 on $S \setminus B(0, 2R)$. Reasoning as before, we get that

$$\begin{aligned}
 \|\boldsymbol{\mu}^*\|_{TV} &\geq \limsup_n |\boldsymbol{\mu}_{\lambda_n, e_n}|(S \cap \overline{B}(0, 2R_\varepsilon)) \geq \limsup_n \int \psi_{R_\varepsilon} d|\boldsymbol{\mu}_{\lambda_n, e_n}| \\
 &\geq \lim_n \left| \int \psi_{R_\varepsilon} \frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n} \cdot d\boldsymbol{\mu}_{\lambda_n, e_n} \right| = \left| \int \psi_{R_\varepsilon} \frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n} \cdot d\boldsymbol{\mu}^* \right| \\
 &\geq \left| \int \psi_{R_\varepsilon} \mathbf{u}_{\boldsymbol{\mu}^*} \cdot d\boldsymbol{\mu}^* \right| - \left| \int \psi_{R_\varepsilon} \left(\frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n} - \mathbf{u}_{\boldsymbol{\mu}^*} \right) \cdot d\boldsymbol{\mu}^* \right| \\
 &\geq \int \psi_{R_\varepsilon} d|\boldsymbol{\mu}^*| - \int \left| \frac{2A^*(f_{e_n} - A(\boldsymbol{\mu}_{\lambda_n, e_n}))}{\lambda_n} - \mathbf{u}_{\boldsymbol{\mu}^*} \right| \cdot d|\boldsymbol{\mu}^*| \geq |\boldsymbol{\mu}^*|(S \cap \overline{B}(0, R)) - \varepsilon \\
 &\geq \|\boldsymbol{\mu}^*\|_{TV} - 2\varepsilon.
 \end{aligned}$$

Hence, in view of (50) and (51), we see from what precedes that for n large enough $|\boldsymbol{\mu}_{\lambda_n, e_n}|(S \setminus \overline{B}(0, 2R_\varepsilon)) \leq 3\varepsilon$, say. Therefore, if we fix a bounded and continuous $\varphi : S \rightarrow \mathbb{R}^3$ with $|\varphi| \leq M$, we have since $\boldsymbol{\mu}_{\lambda_n, e_n}$ converges weak-* to $\boldsymbol{\mu}^*$ that

$$\begin{aligned}
 \limsup_n \left| \int \varphi \cdot d(\boldsymbol{\mu}_{\lambda_n, e_n} - \boldsymbol{\mu}^*) \right| &\leq \limsup_n \left| \int \psi_{2R_\varepsilon} \varphi \cdot d(\boldsymbol{\mu}_{\lambda_n, e_n} - \boldsymbol{\mu}^*) \right| \\
 + \limsup_n \left| \int (1 - \psi_{2R_\varepsilon}) \varphi \cdot d(\boldsymbol{\mu}_{\lambda_n, e_n} - \boldsymbol{\mu}^*) \right| &\leq 0 + 6M\varepsilon.
 \end{aligned}$$

Because $\varepsilon > 0$ was arbitrary, we deduce that $\boldsymbol{\mu}_{\lambda_n, e_n}$ converges to $\boldsymbol{\mu}^*$ in the narrow sense, and the fact that $|\boldsymbol{\mu}_{\lambda_n, e_n}|$ converges to $|\boldsymbol{\mu}^*|$ in the narrow sense as well can be shown in a similar way. This proves (a).

Finally, suppose $\lambda \geq 2 \sup_{x \in S} |(A^* f)(x)|$. Theorem 4.1 shows that the zero magnetization is a minimizer of $\mathcal{F}_{f, \lambda}$ and that any other minimizer $\boldsymbol{\mu}$ must be silent, but then $\mathcal{F}_{f, \lambda}(\boldsymbol{\mu}) = \|f\| + \lambda \|\boldsymbol{\mu}\|_{TV}$ showing that in fact the zero magnetization is the unique minimizer of $\mathcal{F}_{f, \lambda}$. \square

Assertion (a) of Theorem 4.3 entails that any sequence $\boldsymbol{\mu}_{\lambda_n, e_n}$ with $\lambda_n = o(1)$ and $\|e_n\|_{L^2(Q)} = o(\sqrt{\lambda_n})$ has a subsequence converging in the narrow sense to some $\boldsymbol{\mu}^*$ such that $A(\boldsymbol{\mu}^*) = A(\boldsymbol{\mu}_0) = f$ and $\|\boldsymbol{\mu}^*\|_{TV} = \mathfrak{M}(\boldsymbol{\mu}_0)$. If such a $\boldsymbol{\mu}^*$ is unique, we get narrow convergence of $\boldsymbol{\mu}_{\lambda, e}$ to $\boldsymbol{\mu}^*$ as soon as $\lambda \rightarrow 0$ and $\|e\|_{L^2(Q)}/\sqrt{\lambda} \rightarrow 0$. Using Theorems 2.6 and 2.11, we thus obtain a recovery result for ‘‘sparse magnetizations’’ as follows.

weak*to0

Theorem 4.4. *Let $S, Q \subset \mathbb{R}^3$ satisfy the assumptions of Lemma 2.3 with Q compact, and assume in addition that S is a slender set with $S = \bigcup_{i=1}^n S_i$ for some finite collection of disjoint closed sets S_1, S_2, \dots, S_n . Suppose $\boldsymbol{\mu}_0 \in \mathcal{M}(S)^3$ and set $f = A\boldsymbol{\mu}_0$. If either*

- (a) $\boldsymbol{\mu}_0 \llcorner S_i$ is uni-directional for $i = 1, 2, \dots, n$,
- (b) or $\text{supp } \boldsymbol{\mu}_0$ is purely 1-unrectifiable

then $\boldsymbol{\mu}_{\lambda, e}$ converges narrowly to $\boldsymbol{\mu}_0$ and $|\boldsymbol{\mu}_{\lambda, e}|$ converges narrowly to $|\boldsymbol{\mu}_0|$ as $\lambda \rightarrow 0$ and $\|e\|_{L^2(Q)}/\sqrt{\lambda} \rightarrow 0$.

Remark. In the setting of Theorem 4.3 (a), it is generally not true that $\|\boldsymbol{\mu}^* - \boldsymbol{\mu}_{\lambda_n, e_n}\|_{TV} \rightarrow 0$. For instance, let $S \subset \mathbb{R}^2 \times \{0\}$ be compact, assume that $Q \subset \mathbb{R}^2 \times \{h\}$ for some $h > 0$, let ρ be 2-dimensional Hausdorff measure and $\boldsymbol{\mu}_0 = \chi_S v$, where $v \in \mathbb{R}^3$. Then $\boldsymbol{\mu}_0$ is unidirectional, and we know from Theorem 4.4 that $\boldsymbol{\mu}_{\lambda, e}$ converges narrowly to $\boldsymbol{\mu}_0$ as $\lambda \rightarrow 0$ and $\|e\|_{L^2(Q)}/\sqrt{\lambda} \rightarrow 0$. Still, the support of $\boldsymbol{\mu}_{\lambda, e}$ has Hausdorff dimension at most

1, by Corollary 4.2, therefore $\boldsymbol{\mu}_{\lambda,e}$ and $\boldsymbol{\mu}_0$ are mutually singular. Hence $\|\boldsymbol{\mu}_{\lambda,e} - \boldsymbol{\mu}_0\|_{TV} = \|\boldsymbol{\mu}_{\lambda,e}\|_{TV} + \|\boldsymbol{\mu}_0\|_{TV}$ cannot go to zero when λ goes to zero.

As a particular case of Theorem 4.4, taking into account Corollary 4.2, we obtain:

points

Corollary 4.5. *Let $S, Q \subset \mathbb{R}^3$ satisfy the assumptions of Lemma 2.3 with Q compact and S slender. Suppose that $\boldsymbol{\mu}_0 \in \mathcal{M}(S)^3$ is a finite sum of point dipoles: $\boldsymbol{\mu}_0 = \sigma_{j=1}^N v_j \delta_{x^{(j)}}$ for some $v_j \in \mathbb{R}^3 \setminus \{0\}$ and $x^{(j)} \in S$, and set $f = A\boldsymbol{\mu}_0$. Then, to each $\varepsilon, r > 0$ there is $\eta > 0$ such that, whenever $0 < \lambda < \eta$ and $\|e\|_{L^2(Q)}/\sqrt{\lambda} < \eta$:*

- (a) $|\boldsymbol{\mu}_{\lambda,e}|(B(x^{(j)}, r)) - |v_j| < \varepsilon$ for $1 \leq j \leq N$,
- (b) $|\boldsymbol{\mu}_{\lambda,e}|(S \setminus \cup_{j=1}^N B(x^{(j)}, r)) < \varepsilon$,
- (c) $|\boldsymbol{\mu}_{\lambda,e}(B(x^{(j)}, r)) - v_j| < \varepsilon$ for $1 \leq j \leq N$,
- (d) $\text{supp } \boldsymbol{\mu}_{\lambda,e}$ is contained in a finite collection of analytic surfaces, curves and points.

Remark. If S is contained in some unbounded real analytic surface or curve (e.g. a plane or a line) which is disjoint from Q , point (d) of Corollary 4.5 involves only curves and points in case of a surface and only points in case of a curve. This follows from Corollary 4.2.

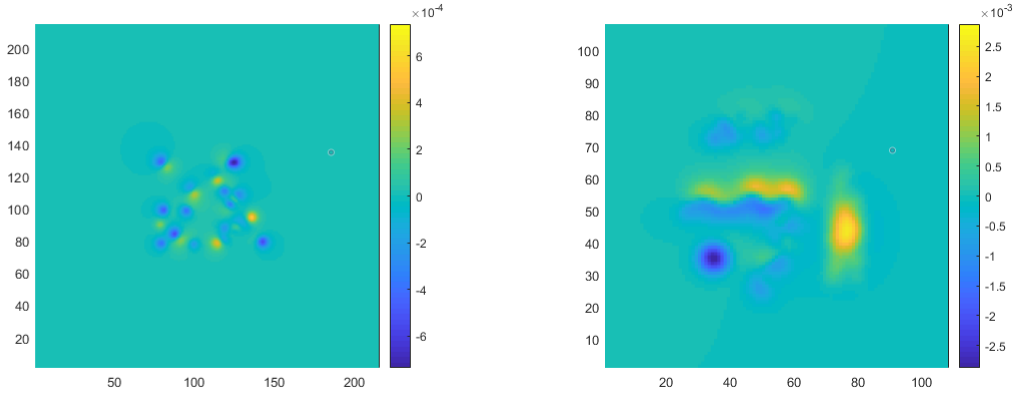


fig:Fields

FIGURE 1. The fields $b_3(\boldsymbol{\mu})$ for (left) a sparse magnetization consisting of 20 dipoles (see Figure 2 for $\boldsymbol{\mu}$ and reconstructions $\boldsymbol{\mu}_\lambda$) and (right) a piecewise unidirectional magnetization with S consisting of four connected components (see Figure 3).

Numerical

5. NUMERICAL EXAMPLES

In this section we present two examples of numerical reconstructions illustrating Theorem 4.4. In both cases, we are considering continuous problems with $v = e_3$ (the measured field is $b_3(\boldsymbol{\mu}_0)$), $e = 0$ (no noise), and S and Q compact subsets of the $z = 0$ and $z = h$ planes, respectively. We discretize the continuous problems by restricting to magnetizations $\boldsymbol{\mu}$'s consisting of a finite number of dipoles located on a rectangular grid in the

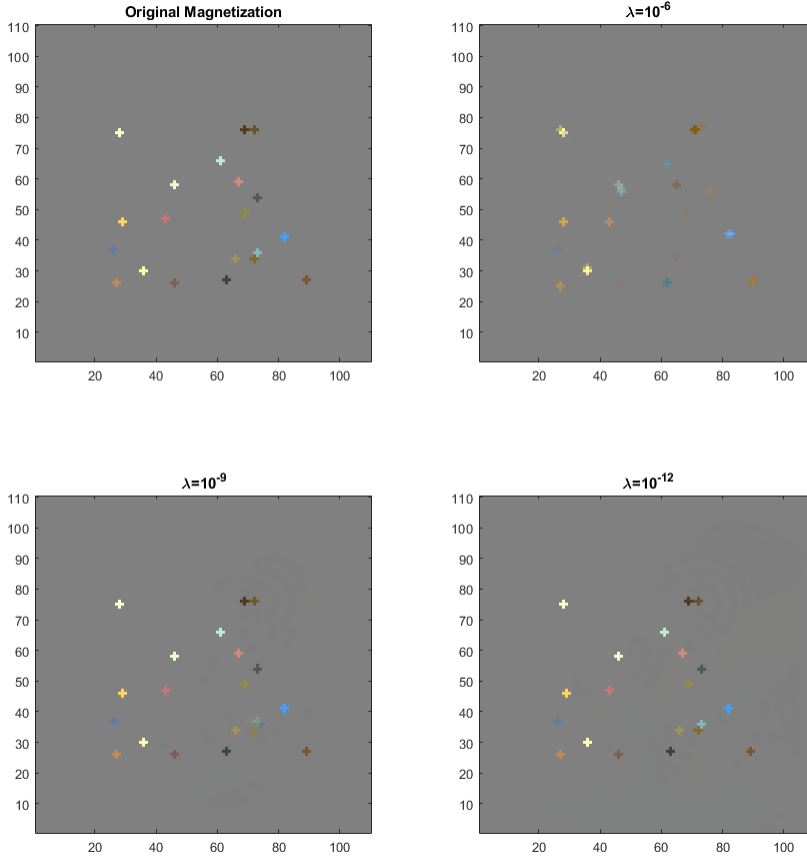


FIGURE 2. Example for a Sparse Magnetization. Each figure shows for its respective magnetization the magnitude and direction of its dipoles by representing different vectors with different colors. The closer a dipole is to zero the closer its assigned color is to grey. The relative distances in total variation to μ_0 are 1.294, 0.207 and 0.004 respectively, confirming convergence as expected. (Convolution with a 3 by 3 cross matrix was used to increase visibility)

fig:Sparse

$z = 0$ plane intersected with S and samples $b_3(\mu)$ evaluated at points in a rectangular grid in the parallel $z = h$ plane intersected with a rectangle Q . Numerical solutions to the discretized problems are then obtained using the Fast Iterative Shrinkage-Thresholding algorithm (FISTA) from [10] together with the In-Crowd algorithm from [29]. We will consider in more detail in forthcoming work the relations between the solutions of such

discretized problems and the associated continuous problems of the type addressed in this paper and we will also address the algorithmic and computational details for obtaining solutions to these discrete problems. The examples provided here are only intended for illustrative purposes.

The continuous problem for the first example is designed to illustrate the recovery of magnetizations with sparse support as in part (b) of Theorem 4.4. In this example S and

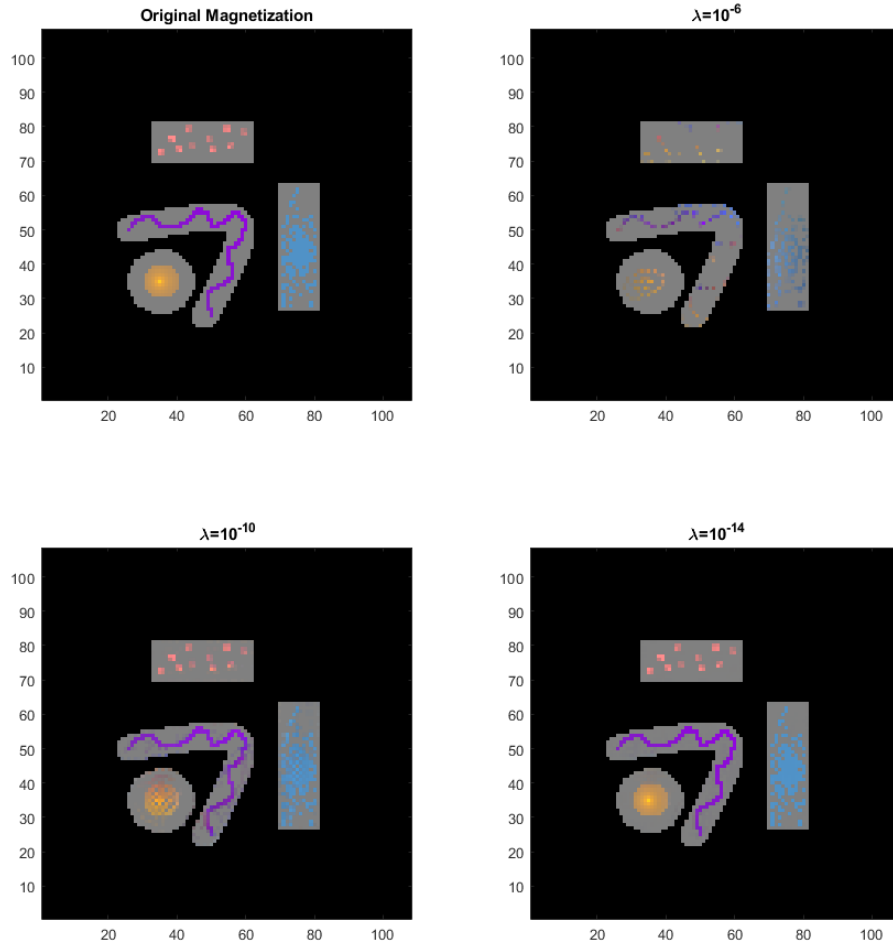


FIGURE 3. Example for a piece-wise Unidirectional Magnetization. The figures are made similarly to the ones of Figure 2 with black representing the complement of S . Here the relative distances are 1.167, 0.247 and 0.015

fig:Uni

Q are squares in the planes $z = 0$ and $z = h = .1$, respectively, and $\boldsymbol{\mu}_0$ consists of 20 dipoles in S with moments of differing directions. The source and measurement grids both have $.0187 \times .0187$ spacing. The source grid consists of 108×108 points and the moments of each dipole are allowed to take on any values in \mathbb{R}^3 . The measurement grid consists of 215×215 points. Reconstructions are computed for three values of λ : 10^{-6} , 10^{-9} and 10^{-12} .

In the second case, S consists of the union of four disjoint compact regions as shown in Figure 3. The restriction of the magnetization $\boldsymbol{\mu}_0$ to each component is unidirectional as in part (a) of Theorem 4.4. The source grid now consists of the grid points that are contained in the set S and again the moments of these dipoles are unconstrained; i.e., there is no uni-directionality assumption when solving for a reconstruction. The $\boldsymbol{\mu}_\lambda$ are taken among all magnetizations supported in those regions for λ equal to 10^{-6} , 10^{-10} and 10^{-14} .

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app1

APPENDIX A. JORDAN-BROUWER SEPARATION THEOREM IN THE NON COMPACT CASE

In this section, we record a proof of the Jordan-Brouwer separation theorem for smooth and connected, complete but not necessarily compact surfaces in \mathbb{R}^3 . The argument applies in any dimension. We are confident the result is known, but we could not find a published reference. More general proofs, valid for non-smooth manifolds as well, could be given using deeper facts from algebraic topology. For instance, one based on Alexander-Lefschetz duality can be modelled after Theorem 14.13 in www.seas.upenn.edu/~jean/sheaves-cohomology.pdf (which deals with compact topological manifolds). More precisely, using Alexander-Spanier cohomology with compact support and appealing to [41, ch. 6, sec. 6 cor. 12 and sec. 9 thm. 10], one can generalize the proof just mentioned to handle the case of non-compact manifolds. Hereafter, we merely deal with smooth surfaces and rely on basic notions from differential-topology, namely intersection theory modulo 2.

Recall that a smooth manifold X of dimension k embedded in \mathbb{R}^n is a subset of the latter, each point of which has a neighborhood V such that $V \cap X = \phi(U)$ where U is an open subset of \mathbb{R}^k and $\phi : U \rightarrow \mathbb{R}^n$ a C^∞ -smooth injective map with injective derivative at every point. The map ϕ is called a parametrization of V with domain U , and the image of its derivative $D\phi(u)$ at u is the tangent space to X at $\phi(u)$, hereafter denoted by $T_{\phi(u)}X$. Then, by the constant rank theorem, there is an open set $W \subset \mathbb{R}^n$ with $W \cap X = V$ and a C^∞ -smooth map $\psi : W \rightarrow U$ such that $\psi \circ \phi = \text{id}$, the identity map of U . The restriction $\psi|_V$ is called a chart with domain V . This allows one to carry over to X local tools from differential calculus, see [30, ch. 1]. We say that X is closed if it is a closed subset of \mathbb{R}^n .

If X, Y are smooth manifolds embedded in \mathbb{R}^m and \mathbb{R}^n respectively, and if $Z \subset Y$ is a smooth embedded submanifold, a smooth map $f : X \rightarrow Y$ is said to be transversal to Z if $\text{Im}Df(x) + T_{f(x)}Z = T_{f(x)}Y$ at every $x \in X$ such that $f(x) \in Z$. If f is transversal to Z , then $f^{-1}(Z)$ is an embedded submanifold of X whose codimension is the same as the codimension of Z in Y . In particular, if X is compact and $\dim X + \dim Z = \dim Y$, then $f^{-1}(Z)$ consists of finitely many points. The residue class modulo 2 of the cardinality of such points is the intersection number of f with Z modulo 2, denoted by $I_2(f, Z)$. If in addition Z is closed in Y , then $I_2(f, Z)$ is invariant under small homotopic deformations of f , and this allows one to define $I_2(f, Z)$ even when f is not transversal to Z , because a suitable but arbitrary small homotopic deformation of f will guarantee transversality, see [30, ch. 2].

JB **Theorem A.1.** *If \mathcal{A} is a C^∞ -smooth complete and connected surface embedded in \mathbb{R}^3 , then $\mathbb{R}^3 \setminus \mathcal{A}$ has two connected components.*

Proof. Let W be a tubular neighborhood of \mathcal{A} in \mathbb{R}^3 [30, Ch. 2, Sec. 3, ex. 3 & 16]. That is, W is an open neighborhood of \mathcal{A} in \mathbb{R}^3 comprised of points y having a unique closest point from X , say x , such that $|y - x| < \varepsilon(x)$ where ε is a suitable smooth and strictly positive function on \mathcal{A} . Thus, we can write $W = \{x + tn(x), x \in \mathcal{A}, |t| < \varepsilon(x)\}$, where $n(x)$ is a normal vector to \mathcal{A} at x of unit length. Note that, for each $x \in \mathcal{A}$, there are two possible (opposite) choices of $n(x)$, but the definition of W makes it irrelevant which one we make. Moreover, if we fix $n(x)$ and $\eta \in (0, 1)$, we can find a neighborhood V of

x in \mathcal{A} such that, to each $y \in V$, there is a unique choice of $n(y)$ with $|n(y) - n(x)| < \eta$. Indeed, if $\phi : U \rightarrow V$ is a parametrization with inverse ψ such that $x \in V$, and if we set $n_y := \partial_{x_1}\phi(\psi(y)) \wedge \partial_{x_2}\phi(\psi(y)) / \|\partial_{x_1}\phi(\psi(y)) \wedge \partial_{x_2}\phi(\psi(y))\|$ where x_1, x_2 are Euclidean coordinates on $U \subset \mathbb{R}^2$ while ∂_{x_j} denotes the partial derivative with respect to x_j and the wedge indicates the vector product, then the two possible choices for $n(y)$ when $y \in V$ are $\pm n_y$. Thus, if we select for instance $n(x) = n_x$ and subsequently set $n(y) = n_y$, we get upon shrinking V if necessary that $|n(x) - n(y)| < \eta$ and $|n(x) + n(y)| > 2 - \eta$ for $y \in V$. As a consequence, if $\Upsilon : [a, b] \rightarrow \mathcal{A}$ is a continuous path, and if n_b is a unit normal vector to \mathcal{A} at $\Upsilon(b)$, there is a continuous choice of $n(\Upsilon(\tau))$ for $\tau \in [a, b]$ such that $n(\Upsilon(b)) = n_b$.

Fix $x_0 \in \mathcal{A}$ and let n_0 be an arbitrary choice for $n(x_0)$. Pick t_0 with $0 < t_0 < \varepsilon(x_0)$, and define two points in W by $x_0^\pm = x_0 \pm t_0 n_0$. We claim that each $y \in \mathbb{R}^3 \setminus \mathcal{A}$ can be joined either to x_0^+ or to x_0^- by a continuous arc contained in $\mathbb{R}^3 \setminus \mathcal{A}$. Indeed, let $\gamma : [0, 1] \rightarrow \mathbb{R}^3$ be a continuous path with $\gamma(0) = y$ and $\gamma(1) = x_0$. Let $\tau_0 \in (0, 1]$ be smallest such that $\gamma(\tau_0) \in \mathcal{A}$; such a τ_0 exists since \mathcal{A} is closed. Pick $0 < \tau_1 < \tau_0$ close enough to τ_0 that $\gamma(\tau_1) \in W$, say $\gamma(\tau_1) = x_1 + t_1 n_1$ where $x_1 \in \mathcal{A}$, $|t_1| < \varepsilon(x_1)$, and n_1 is a unit vector normal to \mathcal{A} at x_1 . Since \mathcal{A} is connected, there is a continuous path $\Upsilon : [\tau_1, 1] \rightarrow \mathcal{A}$ such that $\Upsilon(\tau_1) = x_1$ and $\Upsilon(1) = x_0$. Along the path Υ , there is a continuous choice of $\tau \rightarrow n(\Upsilon(\tau))$ such that $n(x_0) = n_0$; this follows from a previous remark. Changing the sign of t_1 if necessary, we may assume that $n_1 = n(x_1)$. Let $\eta : [\tau_1, 1] \rightarrow \mathbb{R}_+$ be a continuous function such that $0 < |\eta(\tau)| < \varepsilon(\Upsilon(\tau))$ with $\eta(\tau_1) = t_1$ and $\eta(1) = \text{sgn } t_1 |t_0|$. Such an η exists, since ε is continuous and strictly positive while $|t_1| < \varepsilon(x_1)$ and $|t_0| < \varepsilon(x_0)$. Now the concatenation of γ restricted to $[0, \tau_1]$ and $\gamma_1 : [\tau_1, 1] \rightarrow \mathbb{R}^3$ given by $\gamma_1(\tau) = \Upsilon(\tau) + \eta(\tau)n(\Upsilon(\tau))$ is a continuous path from y to either x_0^+ or x_0^- (depending on the sign of t_1) which is entirely contained in $\mathbb{R}^3 \setminus \mathcal{A}$. *This proves the claim*, showing that $\mathbb{R}^3 \setminus \mathcal{A}$ has at most two components. To see that it has at least two, it is enough to know that any smooth cycle $\varphi : \mathbb{S}^1 \rightarrow \mathbb{R}^3$ has intersection number $I_2(\varphi, \mathcal{A}) = 0$ modulo 2. Indeed, if this is the case and if x_0^+ and x_0^- could be joined by a continuous arc $\gamma : [0, 1] \rightarrow \mathbb{R}^3$ not intersecting \mathcal{A} , then γ could be chosen C^∞ -smooth (see [30, Ch.1, Sec. 6, Ex. 3]) and we could complete it into a cycle $\varphi : \mathbb{S}^1 \rightarrow \mathbb{R}^3$ by concatenation with the segment $[x_0^-, x_0^+]$ which intersects \mathcal{A} exactly once (at x_0), in a transversal manner. Elementary modifications at x_0^- and x_0^+ will arrange things so that φ becomes C^∞ -smooth, and this would contradict the fact that the number of intersection points with \mathcal{A} must be even. Now, if \mathbb{D} is the unit disk, any smooth map $\varphi : \mathbb{S}^1 \rightarrow \mathbb{R}^3$ extends to a smooth map $f : \overline{\mathbb{D}} \rightarrow \mathbb{R}^3$ (take for example $f(re^{i\theta}) = e^{1-1/r}\varphi(e^{i\theta})$). Thus, by the boundary theorem [30, p. 80], the intersection number modulo 2 of φ with any smooth and complete embedded submanifold of dimension 2 in \mathbb{R}^3 (in particular with \mathcal{A}) must be zero. This achieves the proof. \square

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