



Mellin transforms and asymptotics: digital sums

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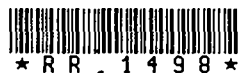
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MELLIN TRANSFORMS AND ASYMPTOTICS : DIGITAL SUMS

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Septembre 1991



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Mellin Transforms and Asymptotics: Digital Sums

PHILIPPE FLAJOLET, PETER GRABNER, PETER KIRSCHENHOFER,
HELMUT PRODINGER, AND ROBERT F. TICHY

Abstract. Arithmetic functions related to number representation systems exhibit various periodicity phenomena. For instance, a well known theorem of Delange expresses the total number of ones in the binary representations of the first n integers in terms of a periodic fractal function.

We show that such periodicity phenomena can be analyzed rather systematically using classical tools from analytic number theory, namely the Mellin–Perron formulæ. This approach yields naturally the Fourier series involved in the expansions of a variety of digital sums related to number representation systems.

Transformation de Mellin et Asymptotique: Les sommes digitales

Résumé. Les fonctions arithmétiques associées aux systèmes de représentation d'entiers mettent en jeu divers phénomènes de périodicité. C'est ainsi qu'un théorème classique de Delange exprime le nombre total de chiffres uns dans la représentation des n premiers entiers au moyen d'une fonction périodique fractale.

L'on montre ici que de telles périodicités peuvent s'analyser assez systématiquement au moyen d'outils usuels de la théorie analytique des nombres, en particulier les formules de Mellin–Perron. Cette approche fournit naturellement les séries de Fourier qui apparaissent dans développements de diverses sommes digitales liées aux systèmes de numérations.

MELLIN TRANSFORMS AND ASYMPTOTICS: DIGITAL SUMS

PHILIPPE FLAJOLET¹, PETER GRABNER², PETER KIRSCHENHOFER¹,
HELMUT PRODINGER¹, AND ROBERT F. TICHY^{1,2}

ABSTRACT. Arithmetic functions related to number representation systems exhibit various periodicity phenomena. For instance, a well known theorem of Delange expresses the total number of ones in the binary representations of the first n integers in terms of a periodic fractal function.

We show that such periodicity phenomena can be analyzed rather systematically using classical tools from analytic number theory, namely the Mellin–Perron formulæ. This approach yields naturally the Fourier series involved in the expansions of a variety of digital sums related to number representation systems.

July 25, 1991

1. INTRODUCTION

Let $S(n)$ represent the total number of 1-digits in the binary representations of the integers $1, 2, \dots, n - 1$. It is not hard to see that

$$(1.1) \quad S(n) = \frac{1}{2}n \log_2 n + o(n \log n),$$

since asymptotically the binary representations contain roughly as many 0's as 1's. The Trollope–Delange formula is more surprising. It expresses $S(n)$ by an exact formula [De75]

$$(1.2) \quad S(n) = \frac{1}{2}n \log_2 n + nF_0(\log_2 n),$$

where $F_0(u)$, a ‘fractal function’, is a continuous, periodic, nowhere differentiable function; $F_0(u)$ has an explicit Fourier expansion that involves the Riemann zeta function, its Fourier coefficient of order k , $k \neq 0$, being

$$f_k = -\frac{1}{\log 2} \frac{\zeta(\chi_k)}{\chi_k(\chi_k + 1)} \quad \text{for } \chi_k = \frac{2\pi ik}{\log 2}.$$

The argument given by Delange relies on a combinatorial decomposition of binary representations of integers, followed by a computation of the Fourier coefficients of the fractal

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function. Our approach instead is more direct and in line with classical methods from analytic number theory. It is based on an integral representation (see Eq. (3.2)), here

$$\frac{1}{n}S(n) - \frac{n-1}{2} = \frac{1}{2\pi i} \int_{2-i\infty}^{2+i\infty} \frac{\zeta(s)}{2^s-1} n^s \frac{ds}{s(s+1)},$$

itself closely related to Mellin transforms and the classical Perron formula. In this context, the periodicity present in $S(n)$ simply arises, by the residue theorem, from poles of the integrand at the regularly spaced poles $s = 2ik\pi/\log 2$.

In other words, as is customary from standard analytic number theory (e.g., the prime number theorem), fluctuations in a number-theoretic function appear to be directly related to singularities of an associated Dirichlet series.

The Mellin Perron formulæ are reviewed briefly in Section 2. In general, they provide *asymptotic* rather than *exact* summation formulæ. A simple additional argument is then needed in order to establish an exact representation like (1.1). Similar exact formulæ are established for the standard sum-of-digit function (Section 3), for the more general case of the number of blocks in binary representations and Gray codes (Section 4), and for a function related to the Cantor set (Section 5),

$$h\left(\sum_i 2^{e_i}\right) = \sum_i 3^{e_i},$$

where the exponents e_i are strictly increasing.

Sections 6 and 7 deal primarily with asymptotic summation formulæ. Section 6 is concerned with the asymptotic evaluation of the function

$$(1.3) \quad \Phi(n) = \sum_{k=0}^{n-1} 2^{\nu(k)},$$

where $\nu(k)$ denotes the binary sum-of-digits function. The value $\Phi(n)$ is also equal to the number of odd binomial coefficients in the first n rows of Pascal's triangle. Stolarsky [Sto77] earlier gave upper and lower bounds for this expression. Applying the Mellin-Perron formula and a pseudo-Tauberian argument the Fourier coefficients of the corresponding fractal function are computed. (Further estimates were given by Harborth [Ha77] and in the q -ary case by Stein [Ste89].) It is found that $\frac{\Phi(N)}{N^\rho}$ is a periodic function of $\log_2 N$, with $\rho = \log_2 3$, see Figure 1 for a graphical rendering. Section 7 is concerned with the asymptotic evaluation of

$$S_3(n) = \sum_{k=0}^{n-1} (-1)^{\nu(3k)},$$

a function obviously related to the distribution of 1-digits in the multiples of three which was first studied by Newman [Ne68]. J. Coquet [Co83] established a Delange type theorem in this case.

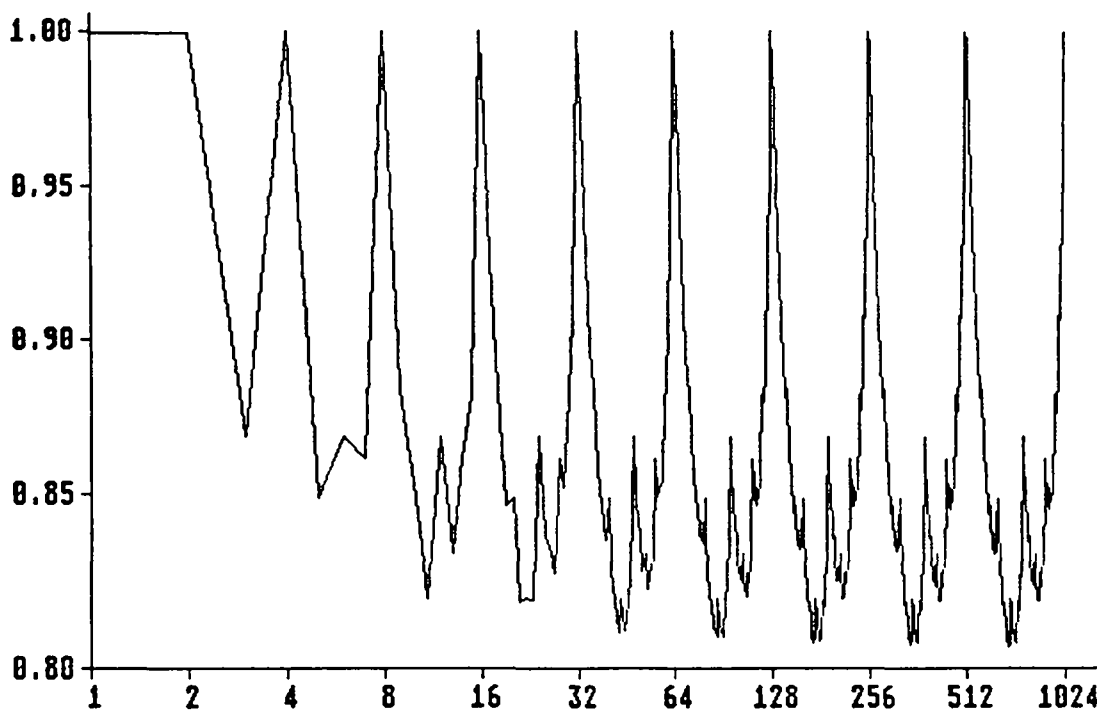


Figure 1. The representation of $\frac{\Phi(n)}{n^\rho}$ plotted against $\log_2 n$; $\Phi(n)$ represents the number of odd binomial coefficients in the first n rows of Pascal's triangle and is also defined by Eq. (1.3); $\rho = \log_2 3$.

The asymptotic formulæ obtained in connection with $\Phi(n)$ and $S_3(n)$ when matched against exact formulæ obtained by direct combinatorial reasoning, lead to new Fourier expansions. This mixed combinatorial-analytic process constitutes another source for summatory formulæ with explicit Fourier coefficients developed in Sections 6 and 7.

Exact summation formulæ related to number representations arise at various places in elementary (combinatorial) number theory as well as in the average case analysis of algorithms. For instance, Delange's formula was used to analyze register allocation algorithms, or equivalently the order of random 'channel networks' in [F-R-V79]. It was later extended to some non-standard digital representations of integers, like Gray code [F-R80], for the purpose of analyzing sorting networks, as well as to occurrences of blocks of digits in standard q -ary representations [Ki83] and subblock occurrences in Gray code representation [K-P84] or to the classical Rudin-Shapiro sequence [Co83]. In a recent paper [G-T91], Delange's result was extended to digit expansions with respect to linear recurrences.

As general references, we refer to Stolarsky's survey [Sto77] and to [K-P-T85]. An especially important paper by the spectrum of its analysis techniques is Brillhart, Erdős and Morton's work [B-E-M83]. It concerns the Rudin-Shapiro sequence, $r(n)$, which gives the parity of the number of blocks 11 in the binary representation of n .

In another direction, the summatory formulæ considered here are also closely related to number-theoretic functions arising in the context of iterated substitutions and so-

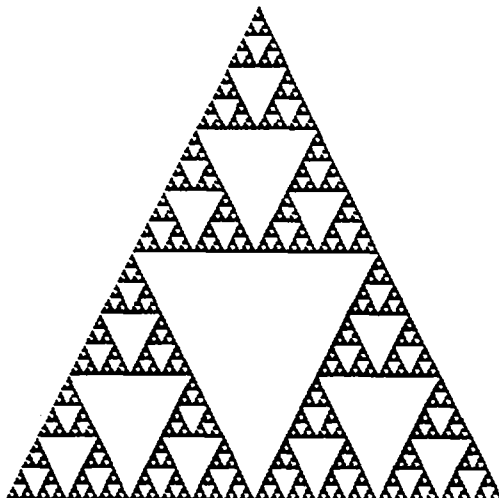


Figure 2. Pascal's triangle modulo 2. The odd numbers are represented by black squares, the even numbers by white squares.

called 'automatic sequences' (see Allouche's paper for a survey [A187]), which constitute a natural framework in which several of our analyses could have been cast. In that framework Dumont and Thomas [D-T89] have used elementary methods to derive for linear functionals of iterated substitution sequences a whole class of asymptotic forms of the type

$$n^\beta (\log_\theta n)^\gamma F(\log n) + o(n^\beta (\log_\theta n)^\gamma),$$

involving some fluctuating function F . Closer to our objectives, Allouche and Cohen have shown that Dirichlet generating functions associated to automatic sequences have meromorphic continuations (see [A187, p.261] and [A-C85]); techniques developed in this paper could then be used in order to provide alternative derivations for some of the results of Dumont and Thomas.

The subject of this paper is finally also close to the classical divide-and-conquer recurrences of theoretical computer science. In a forthcoming paper Flajolet and Golin will discuss several examples for such recurrences that appear in mergesort, heapsort, Karatsuba multiplication, and similar algorithms.

Graphics. As an illustration of the fractal phenomena at stake, we have displayed in Figure 1 the ratio $\frac{\Phi(N)}{N^\rho}$ plotted against $\log_2 N$. When considering successive intervals $[2^{k-1}, 2^k]$, we see the function $\frac{\Phi(N)}{N^\rho}$ which gets refined in a stepwise manner. The picture clearly illustrates the fractal nature of the graph. Figure 2 shows Pascal's triangle mod 2. The picture reveals another aspect of the fractal structure underlying the problem. (Performing an easy transformation one obtains the famous Sierpinski triangle [Fa85]; a popular source for similar graphics is [Wo84].)

2. MELLIN-PERRON FORMULÆ

For completeness, we give a brief outline of the Perron formula by relating it to the

Mellin transform. The resulting summation formulæ are essentially classical, so that we content ourselves with a sketchy description of the analysis involved.

The major reference for Mellin transforms is Doetsch's book [Do50]. Mellin summation is briefly surveyed in [F-R-S84] which is directed towards applications in the average case analysis of algorithms, while in the context of integrals (rather than sums), a useful reference is [Wo89, Chap. III]. The classical Perron formula is discussed at length in Apostol's book [Ap84], and a higher order version is for instance given by Schwarz (see Chapter IV of [Sch69]).

Let $f(x)$ be a function defined over $[0, +\infty)$. Its Mellin transform $f^*(s) = \mathcal{M}[f(x); s]$ is defined by

$$(2.1) \quad f^*(s) = \int_0^{\infty} f(x)x^{s-1} dx.$$

By linearity and the rescaling property we have

$$(2.2) \quad F(x) = \sum_k \lambda_k f(\mu_k x) \Rightarrow F^*(s) = \left(\sum_k \lambda_k \mu_k^{-s} \right) f^*(s).$$

The condition is for s to belong to a 'fundamental strip' defined by the property that the integral giving $f^*(s)$ and the sum $\sum_k \lambda_k \mu_k^{-s}$ are both absolutely convergent.

Similar to the Laplace transform there is an inversion theorem (cf. [Do50]). When applied to (2.2), it provides

$$(2.3) \quad \sum_k \lambda_k f(\mu_k x) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \left(\sum_k \lambda_k \mu_k^{-s} \right) f^*(s) x^{-s} ds,$$

with c in the fundamental strip.

Formula (2.3) could be called Mellin's summation formula. It is especially useful when the integral can be computed by residues, and in that case each residue contributes a term in an asymptotic expansion of $F(x)$.

This formula lends itself to various number theoretic applications, e.g., proofs of the prime number theorem. Introduce the step function $H_0(x)$ defined by

$$H_0(x) = \begin{cases} 1 & \text{if } x \in [0, 1] \\ 0 & \text{if } x > 1, \end{cases}$$

together with the functions $H_m(x) = (1-x)^m H_0(x)$. In the interesting case where $\mu_k \equiv k$, we obtain from (2.3) formulæ of the Perron type that provide integral representations for the iterated summations of arithmetic functions in terms of their Dirichlet generating function:

Theorem 0. *Let $c > 0$ lie in the half-plane of absolute convergence of $\sum_k \lambda_k k^{-s}$. Then for any $m \geq 1$, we have*

$$(2.4) \quad \frac{1}{m!} \sum_{1 \leq k < n} \lambda_k \left(1 - \frac{k}{n}\right)^m = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \left(\sum_{k \geq 1} \frac{\lambda_k}{k^s}\right) n^s \frac{ds}{s(s+1)\cdots(s+m)}.$$

For $m = 0$,

$$\sum_{1 \leq k < n} \lambda_k + \frac{\lambda_n}{2} = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \left(\sum_{k \geq 1} \frac{\lambda_k}{k^s}\right) n^s \frac{ds}{s}.$$

Formula (2.4) is obtained from (2.3) by setting $x \equiv n^{-1}$, $f(x) \equiv H_m(x)$, and observing that $H_m^*(s) = m!(s(s+1)\cdots(s+m))^{-1}$. For $m = 0$ the formula has to be modified slightly by taking a principal value for the sum, since $H_0(x)$ is discontinuous at $x = 1$. See also [Ap84, p.245] for a direct proof of the $m = 0$ case.

For instance, if we use $\lambda_k \equiv 1$ and $m = 1$, we get

$$(2.5) \quad \frac{n-1}{2} = \frac{1}{2\pi i} \int_{2-i\infty}^{2+i\infty} \zeta(s) n^s \frac{ds}{s(s+1)}.$$

Shifting the line of integration to the left and taking residues into account we obtain

$$(2.6) \quad 0 = \int_{-\frac{1}{4}-i\infty}^{-\frac{1}{4}+i\infty} \zeta(s) n^s \frac{ds}{s(s+1)}.$$

Identity (2.6) is the basis for the existence of several exact rather than plainly asymptotic summation formulæ.

3. SUM-OF-DIGITS FUNCTIONS

We apply the Mellin–Perron technique described in the preceding section in order to derive a new proof of Delange’s theorem.

Theorem 1. [Delange] *The sum-of-digits function $S(n)$ satisfies*

$$S(n) = \frac{1}{2} n \log_2 n + n F_0(\log_2 n),$$

where $F_0(u)$ is representable by the Fourier series $F_0(u) = \sum_{k \in \mathbb{Z}} f_k e^{2\pi i k u}$ and

$$f_0 = \frac{\log_2 \pi}{2} - \frac{1}{2 \log 2} - \frac{1}{4}$$

$$f_k = -\frac{1}{\log 2} \frac{\zeta(\chi_k)}{\chi_k(\chi_k + 1)} \quad \text{for } \chi_k = \frac{2\pi i k}{\log 2}, \quad k \neq 0.$$

Theorem 3. Let $(n; w)$ denote the number of occurrences of the 0-1-string w as a contiguous subblock in the binary representation of the integer n . (If w starts with 0, we also count occurrences that overhang to the left of the most significant digit of n ; we only exclude strings w consisting solely of 0's².) Then the mean number of occurrences, $\frac{1}{n} \sum_{k < n} (k; w)$, is given by

$$\frac{1}{n} \sum_{k < n} (k; w) = \frac{\log_2 n}{2^{|w|}} + H_w(\log_2 n) + \frac{E_w(n)}{n}.$$

There $|w|$ denotes the length of the string w , $H_w(x)$ is a continuous periodic function of period 1 with Fourier expansion $\sum_{k \in \mathbb{Z}} h_k e^{2k\pi i x}$,

$$h_0 = \log_2 \left(\frac{\Gamma((0.w)_2)}{\Gamma((0.w)_2 + 2^{-|w|})} \right) - \frac{1}{2^{|w|}} \left(|w| - \frac{1}{2} + \frac{1}{\log 2} \right)$$

$$h_k = \frac{\zeta(\chi_k, (0.w)_2) - \zeta(\chi_k, (0.w)_2 + 2^{-|w|})}{(\log 2)\chi_k(\chi_k + 1)},$$

$\zeta(z, a)$ the Hurwitz ζ -function, $(x)_2$ denotes the real number with binary representation x and $E_w(n)$ is a dyadic rational with denominator $2^{|w|}$ which is described explicitly below in (4.9).

Proof. As in Section 3 we start with summation by parts to find

$$(4.1) \quad \sum_{k < n} (k; w) = \sum_{k < n} \Delta_w(k)(n - k),$$

where $\Delta_w(k) = (k; w) - (k - 1; w)$. The differences $\Delta_w(k)$ obey the following recurrence relation: If $n = 2^{|w|}k + r$ is even, we have

$$(4.2) \quad \Delta_w(n) = \Delta_w\left(\frac{n}{2}\right) + \begin{cases} 1 & \text{if } (w)_2 = r \\ -1 & \text{if } (w)_2 = r - 1 \\ 0 & \text{otherwise} \end{cases}$$

If $n = 2^{|w|}k + r$ is odd, we simply have

$$(4.3) \quad \Delta_w(n) = \begin{cases} 1 & \text{if } (w)_2 = r \\ -1 & \text{if } (w)_2 = r - 1 \\ 0 & \text{otherwise.} \end{cases}$$

From the recurrences (4.2) and (4.3) it easily follows that the Dirichlet series $A_w(s)$ of the differences $\Delta_w(n)$ satisfies

$$A_w(s) = \sum_{n \geq 1} \frac{\Delta_w(n)}{n^s} = \sum_{k \geq 1} \frac{\Delta_w(k)}{(2k)^s} + \sum_{k \geq 0} \frac{1}{(2^{|w|}k + (w)_2)^s} - \sum_{k \geq 0} \frac{1}{(2^{|w|}k + (w)_2 + 1)^s},$$

²A formula for this case exists, but is very unpleasant to formulate

so that

$$(4.4) \quad \left(1 - \frac{1}{2^s}\right) A_w(s) = \frac{1}{2^{|w|s}} \left(\zeta(s, (0.w)_2) - \zeta(s, (0.w)_2 + 2^{-|w|}) \right),$$

where

$$\zeta(s, a) = \sum_{n \geq 0} \frac{1}{(n+a)^s}$$

is the Hurwitz ζ -function [W-W27] and $(0.w)_2$ denotes the rational number $(w)_2 2^{-|w|}$.

From (4.1) and (4.4) we find using Perron's formula for $m = 1$:

$$\begin{aligned} \frac{1}{n} \sum_{k < n} (k; w) &= \frac{1}{2\pi i} \int_{\frac{3}{2}-i\infty}^{\frac{3}{2}+i\infty} \frac{1}{2^{|w|s}} \left(\zeta(s, (0.w)_2) - \zeta(s, (0.w)_2 + 2^{-|w|}) \right) \times \\ &\quad \times \frac{(2n)^s ds}{(2^s - 1) s(s+1)} \end{aligned}$$

Shifting the contour of integration to the left, we observe that the first order poles of the Hurwitz ζ -functions cancel since both have residue 1, so that the main contribution comes from the second order pole $s = 0$. The residue is

$$(4.5) \quad C_1 \log_2 n + \frac{C_2}{\log 2} + C_1 \left(\frac{1}{2} - |w| - \frac{1}{\log 2} \right),$$

with

$$C_1 = \zeta(0, (0.w)_2) - \zeta(0, (0.w)_2 + 2^{-|w|}) = 2^{-|w|},$$

since $\zeta(0, a) = \frac{1}{2} - a$ and

$$C_2 = \zeta'(0, (0.w)_2) - \zeta'(0, (0.w)_2 + 2^{-|w|}) = \log \frac{\Gamma((0.w)_2)}{\Gamma((0.w)_2 + 2^{-|w|})},$$

since $\zeta'(0, a) = \log \Gamma(a) - \frac{1}{2} \log(2\pi)$ [W-W27].

Thus the main term and the mean h_0 of the fluctuating term are established. The other Fourier coefficients h_k are easily derived from the residues at the simple poles $\chi_k = \frac{2\pi i k}{\log 2}$, $k \neq 0$.

We still have to analyze the remainder term

$$(4.6) \quad \begin{aligned} R_n &= \frac{1}{2\pi i} \int_{-\frac{1}{4}-i\infty}^{-\frac{1}{4}+i\infty} \frac{1}{2^{|w|s}} \left(\zeta(s, (0.w)_2) - \zeta(s, (0.w)_2 + 2^{-|w|}) \right) \times \\ &\quad \times \frac{(2n)^s ds}{(2^s - 1) s(s+1)} = - \sum_{k \geq 0} R'_{2^k n}, \end{aligned}$$

where

$$R'_n = \frac{1}{2\pi i} \int_{-\frac{1}{4}-i\infty}^{-\frac{1}{4}+i\infty} \frac{1}{2^{|w|s}} \left(\zeta(s, (0.w)_2) - \zeta(s, (0.w)_2 + 2^{-|w|}) \right) \frac{(2n)^s ds}{s(s+1)}.$$

After shifting the contour back to the right, we find, by taking into account the residues at $s = 0$,

$$(4.7) \quad R'_n = -\frac{1}{2^{|w|}} + \sum_{k < 2n} \lambda_k \left(1 - \frac{k}{2n} \right),$$

where

$$\lambda_k = \begin{cases} 1 & \text{if } k \equiv (w)_2 \pmod{2^{|w|}} \\ -1 & \text{if } k \equiv (w)_2 + 1 \pmod{2^{|w|}} \\ 0 & \text{otherwise.} \end{cases}$$

The sum in (4.7) can be computed explicitly to give

$$(4.8) \quad R'_n = \begin{cases} \frac{2^{|w|}-1-r}{2^{|w|+1}} \frac{1}{n} & \text{if } r = 2n - 1 \pmod{2^{|w|}} \geq (w)_2 \\ \frac{-1-r}{2^{|w|+1}} \frac{1}{n} & \text{if } r = 2n - 1 \pmod{2^{|w|}} < (w)_2. \end{cases}$$

From (4.8) we see that $R'_{2^k n}$ will be zero for $k \geq |w| - 1$, so that in fact (4.6) reduces to

$$(4.9) \quad \frac{E_w(n)}{n} = - \sum_{k=0}^{|w|-2} R'_{2^k n},$$

and this completes the proof. \square

Theorem 3 and formula (4.9) have a number of consequences of special interest:

Corollary 1. *If $|w| = 1$ and w is the 1-digit, i.e. in the case of the sum-of-digits function, we have $E_w(n) = 0$, as stated already in Theorem 1.*

Corollary 2. *If $|w| = 2$, the remainder terms $\tilde{E}_w(n)$ are given by the scheme*

	$E_{01}(n)$	$E_{10}(n)$	$E_{11}(n)$
n even	0	0	0
n odd	$-\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{4}$

Corollary 3. [Flajolet, Ramshaw] *The mean value of the sum of digits in the Gray Code of n is given by*

$$\frac{\log_2 n}{2} + F_1(\log_2 n),$$

with F_1 as described in Theorem 2.

Proof. An alternative proof to that of Theorem 2 runs as follows. The k -th bit in the Gray Code $GC(n)$ of n is given by the sum mod 2 of the k -th and $k + 1$ -st digit in the binary

representation of n . Thus the number of 1's in $GC(n)$ is just $(n; 01) + (n; 10)$, where we have to count the one occurrence of 01 overhanging to the left of the most significant 1 in the binary representation of n . It follows that the mean is given by

$$\begin{aligned} & 2 \frac{\log_2 n}{4} + H_{01}(\log_2 n) + H_{10}(\log_2 n) + \frac{E_{01}(n)}{n} + \frac{E_{10}(n)}{n} \\ &= \frac{\log_2 n}{2} + F_1(\log_2 n), \end{aligned}$$

which also relates F_1 to H_{01} and H_{10} . \square

Remark 2. All results in Sections 3 and 4 are easily generalized to base q representations. As an application of the special instance $q = 4$ we could get an alternative proof of a result due to Osbaldestin and Shiu [O-S89] concerning the number of integers $\leq n$ that are representable as a sum of three squares.

5. TRIADIC BINARY NUMBERS

Let $h(n)$ be the number that results from interpreting in base 3 the binary representation of n , i.e.,

$$h\left(\sum_i 2^{e_i}\right) = \sum_i 3^{e_i},$$

where the exponents e_i are strictly increasing. It is known that $h(1) < h(2) < \dots < h(n)$ is the 'minimal' sequence of n positive integers not containing an arithmetic progression. The sequence is also an analog of Cantor's triadic set. An exact formula for the summatory function H of h is established.

Theorem 4. For the summatory function $H(n) = \sum_{k < n} h(k)$ we have

$$H(n) = n^{\rho+1} F_3(\log_2 n) - \frac{1}{4}n,$$

where $\rho = \log_2 3$ and $F_3(u)$ is given by its Fourier representation

$$F_3(u) = \frac{1}{3 \log 2} \sum_{k \in \mathbb{Z}} \zeta(\rho + \chi_k) \frac{e^{2\pi i k u}}{(\rho + \chi_k)(\rho + \chi_k + 1)},$$

with $\chi_k = \frac{2k\pi i}{\log 2}$.

Proof. Using $h(n) - h(n-1) = \frac{1}{2}(3^{\nu_2(n)} + 1)$ we obtain

$$\begin{aligned} \sum_{n=1}^{\infty} \frac{h(n) - h(n-1)}{n^s} &= \frac{1}{2} \sum_{n=1}^{\infty} \frac{3^{\nu_2(n)}}{n^s} + \frac{1}{2} \zeta(s) \\ &= \frac{1}{2} \sum_{k=0}^{\infty} \frac{3^k}{2^{ks}} \sum_{n=1}^{\infty} \frac{1}{(2n-1)^s} + \frac{1}{2} \zeta(s) \\ &= \frac{2^s - 2}{2^s - 3} \zeta(s). \end{aligned}$$

Applying the Mellin-Perron summation formula (2.4) with $c = 3$ and shifting the line of integration yields

$$H(n) = \frac{n}{2\pi i} \int_{-\frac{1}{4}-i\infty}^{-\frac{1}{4}+i\infty} \frac{2^s - 2}{2^s - 3} \zeta(s) n^s \frac{ds}{s(s+1)} \\ + n^{\rho+1} \sum_{k \in \mathbb{Z}} \frac{1}{3 \log 2} \zeta(\rho + \chi_k) \frac{n^{\chi_k}}{(\rho + \chi_k)(\rho + \chi_k + 1)} + \frac{1}{2} \zeta(0)n.$$

The remainder integral is 0 by the same argument as for the sum-of-digits function, and the proof is complete. \square

Remark 3. Obviously the base 3 can be replaced by an arbitrary number $\alpha > 1$ and $\alpha \neq 2$. The corresponding exact formula is

$$\sum_{k < n} h_\alpha(k) = n^{\rho+1} F_\alpha(\log_2 n) - \frac{1}{2(\alpha-1)} n,$$

where F_α has a Fourier expansion similar to F_3 and $\rho = \log_2 \alpha$.

Remark 4. The function g defined by $h(n) = n^\rho g(\log_2 n)$ is periodic with period 1 but not continuous.

6. ODD NUMBERS IN PASCAL'S TRIANGLE

In this section we establish an exact formula for the summatory function

$$\Phi(n) = \sum_{0 \leq k < n} 2^{\nu(k)}.$$

As pointed out in the introduction, $\Phi(n)$ is the number of odd binomial coefficients in the first n rows of Pascal's triangle.

The application of Mellin-Perron techniques requires convergence of the complex integral of Theorem 0. For an m -fold summation, the 'kernel' in the integral involves $1/(s(s+1)\cdots(s+m-1))$, which decreases at infinity like $|s|^{-m}$. Thus, higher summations lead to better converging (inverse Mellin) integrals.

For the problem of $\Phi(n)$, we thus start with the double summatory function

$$\Psi(N) = \sum_{1 \leq n < N} (\Phi(n) - 1),$$

where Mellin-Perron is easier to apply since the Fourier expansion converges absolutely (We have to subtract 1, because the summation in (2.4) starts at $n = 1$.) The formula that we get in this way is asymptotic.

Theorem 5. *The arithmetic function $\Psi(N)$ satisfies the asymptotic estimate*

$$\Psi(N) = N^{\rho+1}G(\log_2 N) + O(N^{2+\epsilon}),$$

for arbitrary $\epsilon > 0$, where G is a continuous periodic function with period 1. G admits an absolutely convergent Fourier expansion

$$G(u) = \sum_{k \in \mathbb{Z}} g_k e^{2\pi i k u}$$

with

$$g_k = \frac{2}{\log 2} \frac{(1 - B(\rho_k))}{\rho_k(\rho_k + 1)},$$

where $\rho_k = \rho + \frac{2k\pi i}{\log 2}$ and

$$(6.1) \quad B(s) = \frac{1}{\Gamma(s)} \int_0^\infty \left((1 - e^{-t}) \left(\prod_{k=1}^\infty (1 + 2e^{-t2^k}) \right) - 1 \right) t^{s-1} dt.$$

In order to come back to the Fourier expansion of F we need an external argument in order to convert the expansion of $\Psi(N)$ into an expansion for $\Phi(n)$. One ingredient is a direct combinatorial proof of existence for the fluctuating part of $\Phi(n)$; this induces corresponding periodicities for $\Psi(N)$, and by identification, we *indirectly* derive the Fourier expansion relative to $\Phi(n)$. (This process is a pseudo-Tauberian argument!)

Theorem 6. *The summatory function $\Phi(n)$ satisfies the exact formula*

$$(6.2) \quad \Phi(n) = n^\rho F(\log_2 n),$$

where $\rho = \log_2 3$ and F is a continuous function of period 1. The Fourier coefficients of $F(u)$ are given by

$$(6.3) \quad f_k = \frac{2}{\log 2} \frac{(1 - B(\rho_k))}{\rho_k},$$

and in particular the mean value of $F(u)$ is approximately

$$f_0 \approx 0.86360\ 49963\ 99079\ 60496\ 05033\ 61308\ 09499.$$

$F(u)$ is represented by its Fourier series in the sense of standard $(C, 1)$ Cesàro averages.

Observe that from [Ha77] it is already known that

$$0.812 \leq F(u) \leq 1.$$

Proof. (Theorem 5) Let

$$(6.4) \quad A(s) = \sum_{n=1}^\infty \frac{2^{\nu(n)}}{n^s}$$

be the Dirichlet generating function of $2^{\nu(n)}$. Since $\Psi(N)$ is a double summation of $2^{\nu(k)}$ we have an integral representation by means of the iterated Mellin-Perron formula (2.4). We get

$$(6.5) \quad \Psi(N) = \frac{N}{2\pi i} \int_{3-i\infty}^{3+i\infty} A(s)N^s \frac{ds}{s(s+1)},$$

where the abscissa $\Re(s) = 3$ has been chosen, since $A(3)$ converges absolutely. We need to locate the singularities of $A(s)$. From the recurrences

$$\nu(2k+1) = \nu(k) + 1 \quad \text{and} \quad \nu(2k) = \nu(k)$$

we get

$$\begin{aligned} A(s) &= \sum_{k \equiv 0 \pmod{2}} \frac{2^{\nu(k)}}{n^s} + \sum_{k \equiv 1 \pmod{2}} \frac{2^{\nu(k)}}{k^s} = \\ &= \frac{1}{2^s} A(s) + 2 \sum_{l=0}^{\infty} \frac{2^{\nu(l)}}{(2l+1)^s} = \\ &= \frac{3}{2^s} A(s) + 2 - 2B(s) \end{aligned}$$

with

$$(6.6) \quad B(s) = \sum_{k=1}^{\infty} 2^{\nu(k)} \left(\frac{1}{(2k)^s} - \frac{1}{(2k+1)^s} \right).$$

Using summation by parts and Stolarsky's estimate (cf. [Sto77])

$$(6.7) \quad \frac{1}{3} < \frac{\Phi(n)}{n^\rho} < 3$$

we know that $B(s)$ converges for $\Re(s) > \rho - 1$ and by

$$(6.7) \quad A(s) = \frac{2^s}{2^s - 3} (2 - 2B(s))$$

$A(s)$ has abscissa of convergence equal to ρ . This expression also provides us with the analytic continuation of $A(s)$ for $\Re(s) > \rho - 1$. We see that $A(s)$ is meromorphic with simple poles at the points $\rho_k = \rho + \frac{2\pi i k}{\log 2}$.

In order to shift the contour of integration in (6.5) to the left we need that $A(s)$ does not grow too large along vertical lines. For $s = \sigma + it$ $\sigma > 1$ and $|t| > \frac{\sigma}{\sqrt{3}}$ we have

$$\left| 1 - \left(1 - \frac{1}{2k+1} \right)^s \right| \leq \min \left(2, \frac{|t|}{k} \right).$$

Thus we obtain

$$\begin{aligned} |B(s)| &= \left| \sum_{k=1}^{\infty} \frac{2^{\nu(k)}}{(2k)^s} \left(1 - \left(1 - \frac{1}{2k+1} \right)^s \right) \right| \leq \\ &\leq 2 \sum_{1 \leq k \leq 2|t|} \frac{2^{\nu(k)}}{(2k)^{\sigma}} + \sum_{k > 2|t|} \frac{2^{\nu(k)}}{(2k)^{\sigma}} \frac{|t|}{k} \ll |t|^{2-\sigma}. \end{aligned}$$

Shifting the line of integration to $\Re(s) = 1 + \varepsilon < \rho$, noting that $|A(1 + \varepsilon + it)| \ll |t|^{1-\varepsilon}$ and taking the residues at the poles ρ_k into account we get

$$\begin{aligned} \Psi(N) &= \frac{N}{2\pi i} \int_{1+\varepsilon-i\infty}^{1+\varepsilon+i\infty} A(s) N^s \frac{ds}{s(s+1)} + \\ &\quad + N^{\rho+1} \sum_{k \in \mathbb{Z}} \frac{2}{\log 2} \frac{1 - B(\rho_k)}{\rho_k(\rho_k + 1)}. \end{aligned}$$

Estimating the integral trivially we derive the asymptotic formula for $\Psi(N)$.

So far, $B(s)$ is defined in terms of the sequence $\{2^{\nu(k)}\}$ itself. An integral representation derives from an ordinary generating function, setting

$$\varphi(t) = \sum_{k=0}^{\infty} 2^{\nu(k)} e^{-kt} = \prod_{j=0}^{\infty} \left(1 + 2e^{-2^j t} \right).$$

Consider the Mellin transform of $(1 - e^{-t})\varphi(2t) - 1$; by Formula (2.2) (with $\mu_k = k$, $f(t) = e^{-t}$), we get the integral representation for $B(s)$. Thus the proof of Theorem 5 is completed. \square

In order to get information on the number-theoretic function $\Phi(n)$ we first work out a refinement of Stolarsky's elementary approach.

Proposition 1. *There exists an exact summation formula*

$$(6.8) \quad \Phi(n) = n^{\rho} F(\log_2 n)$$

with F continuous and periodic with period 1.

Proof. From [Sto77] there is an alternative formula

$$(6.9) \quad \Phi \left(n = \sum_{i=1}^r 2^{e_i} \right) = \sum_{i=1}^r 2^{i-1} 3^{e_i}$$

with decreasing exponents e_i . Pulling out the main term we get

$$(6.10) \quad \Phi(n) = 3^{e_1} \sum_{i=1}^r 2^{i-1} 3^{e_i - e_1},$$

where $e_1 = \lfloor \log_2 n \rfloor$.

We now define a real function $\psi(x)$ on the interval $[1, 2]$ as follows. Let

$$x = \sum_{j=0}^{\infty} 2^{-d_j}$$

with $0 = d_0 < d_1 < \dots$. Then we set

$$(6.11) \quad \psi(x) = \sum_{j=0}^{\infty} 2^j 3^{-d_j}.$$

Note that ψ is well-defined since the dyadic rationals are written in their infinite representation. Next we show the continuity of ψ . As the representation of dyadic irrationals is unique the continuity at these points follows immediately, since (because of $d_j \geq j$) the expansion (6.11) converges faster than a geometric series with quotient $\frac{2}{3}$. For the proof of continuity at dyadic rationals we have to show

$$\psi\left(\sum_{j=0}^k 2^{-d_j}\right) = \psi\left(\sum_{j=0}^{k-1} 2^{-d_j} + \sum_{l=d_k+1}^{\infty} 2^{-l}\right),$$

which follows immediately by direct computation. Note that $\psi(1) = 1$ and $\psi(2) = 3$ and that ψ satisfies a Lipschitz-condition

$$|\psi(x) - \psi(y)| \leq C|x - y|^{\rho-1}.$$

Using the function ψ we can write

$$(6.12) \quad \Phi(n) = 3^{\lfloor \log_2 n \rfloor} \psi\left(\frac{n}{2^{\lfloor \log_2 n \rfloor}}\right),$$

since $\frac{n}{2^{\lfloor \log_2 n \rfloor}}$ is nothing but n "scaled" in binary to the interval $[1, 2)$. Formula (6.12) thus transforms into

$$\Phi(n) = n^{\rho} F(\{\log_2 n\}),$$

where

$$F(u) = 3^{-u} \psi(2^u)$$

is defined over the interval $[0, 1]$ and $\{x\} = x - \lfloor x \rfloor$ denotes the fractional part. F can be extended into a periodic function since $F(0) = F(1)$. The proof of Proposition 1 is completed. \square

For the computation of the Fourier coefficients of F we make use of Theorem 5 and the following simple pseudo-Tauberian argument.

Proposition 2. *Let f be a continuous function and periodic with period 1 and let τ be a complex number with $\Re(\tau) > 0$. Then there exists a continuously differentiable function g of period 1 such that*

$$(6.13) \quad \frac{1}{N^{\tau+1}} \sum_{n < N} n^{\tau} f(\log_2 n) = g(\log_2 n) + o(1)$$

$$(6.14) \quad \int_0^1 g(u) du = \frac{1}{\tau+1} \int_0^1 f(u) du.$$

Proof. We set

$$(6.15) \quad g(u) = \log 2 \cdot \frac{2^{-(\tau+1)u}}{2^{\tau+1} - 1} \int_0^1 2^{(\tau+1)t} f(t) dt + \log 2 \cdot 2^{-(\tau+1)u} \int_0^u 2^{(\tau+1)t} f(t) dt.$$

Obviously g is continuously differentiable and (6.14) follows by a straightforward application of integration by parts. Further we note that $g(0) = g(1)$.

In order to prove (6.13) we proceed as follows

$$\begin{aligned} & \frac{1}{N^{\tau+1}} \sum_{n < N} n^{\tau} f(\log_2 n) = \\ &= \frac{1}{N^{\tau+1}} \sum_{p=0}^{\lfloor \log_2 N \rfloor - 1} \sum_{2^p \leq n < 2^{p+1}} n^{\tau} f(\log_2 n) + \frac{1}{N^{\tau+1}} \sum_{n=2^{\lfloor \log_2 N \rfloor}}^N n^{\tau} f(\log_2 n) = \\ &= \frac{1}{N^{\tau+1}} \sum_{p < \lfloor \log_2 N \rfloor} 2^{p(\tau+1)} \sum_{1 \leq x < 2} x^{\tau} f(\log_2 x) \Delta x + \\ &+ \frac{1}{N^{\tau+1}} 2^{\lfloor \log_2 N \rfloor (\tau+1)} \sum_{1 \leq x < y} x^{\tau} f(\log_2 x) \Delta x, \end{aligned}$$

where $x = \frac{n}{2^p}$, $y = \frac{N}{2^{\lfloor \log_2 N \rfloor}}$ and x runs through all dyadic rationals with denominator 2^p and $\Delta x = 2^{-p}$, $p = 0, \dots, \lfloor \log_2 N \rfloor$. Now we interpret the sums over x as Riemann sums. Thus we have with remainder terms $\varepsilon(p)$ tending to 0 (for $p \rightarrow \infty$)

$$\begin{aligned} & \frac{1}{N^{\tau+1}} \sum_{n < N} n^{\tau} f(\log_2 n) = \\ &= \sum_{p < \lfloor \log_2 N \rfloor} 2^{(p - \lfloor \log_2 N \rfloor)(\tau+1)} y^{-(\tau+1)} \left(\int_1^2 x^{\tau} f(\log_2 x) dx + \varepsilon(p) \right) + \\ &+ y^{-(\tau+1)} \left(\int_1^y x^{\tau} f(\log_2 x) dx + \varepsilon(\lfloor \log_2 N \rfloor) \right) = \\ &= g(\log_2 y) + y^{-(\tau+1)} \sum_{p=0}^{\lfloor \log_2 N \rfloor} \varepsilon(p) 2^{-(\lfloor \log_2 N \rfloor - p)}. \end{aligned}$$

We note that only $\varepsilon(\lfloor \log_2 N \rfloor)$ depends on y . Since the convergence of Riemann sums is uniform with respect to the upper limit y , the remainder term tends to 0. Thus the proof of Proposition 2 is complete. \square

Proof. (Theorem 6) We can now conclude and determine the Fourier coefficients f_k of the fractal function F in Theorem 6. We set $\tau = \rho_k$ in Proposition 2 and apply (6.14)

$$f_k = \int_0^1 F(u)e^{-2\pi iku} du = (\rho_k + 1) \int_0^1 G(u)e^{-2\pi iku} du = (\rho_k + 1)g_k.$$

Inserting the value of g_k yields

$$f_k = \frac{2}{\log 2} \frac{1 - B(\rho_k)}{\rho_k}.$$

Using $|B(\rho_k)| \ll k^{2-\rho}$ we obtain the L^2 -convergence of the Fourier expansion of F . Also, since we know that $F(u)$ is continuous, its Fourier series converges in the mean by Fejér's theorem [Ko88]. (More information on the convergence of the Fourier series would have to depend on a more detailed knowledge of the analytical behaviour of the function B .) This completes the proof of Theorem 6. \square

7. THE NEWMAN-COQUET FUNCTION

In this section we study the function

$$(7.1) \quad S_3(n) = \sum_{k < n} (-1)^{\nu(3k)}.$$

The motivation for the study of this function goes back to Newman [Ne68] who noted that examination of the multiples of three, 3, 6, 9, 12, 15, 18, 21, 24, 27, ..., written in the base two,

$$11, 110, 1001, 1100, 1111, 10010, 10101, 11000, 11011, \dots,$$

shows a definite preponderance of those containing an even number of one-digits over those containing an odd number. Newman proved that this strange behaviour persists forever. Coquet [Co83] gave an exact formula by Delange type computations. Our method uses this result and allows us to compute the Fourier coefficients (especially the mean value) of the related fractal function.

Theorem 7. *The summatory function S_3 satisfies the exact formula*

$$S_3(n) = n^\alpha \psi(\log_4 n) + \frac{\eta(n)}{3},$$

where ψ is a continuous nowhere differentiable function of period 1, η is given by

$$\eta(n) = \begin{cases} 0 & \text{if } n \text{ is even} \\ (-1)^{\nu(3n-3)} & \text{if } n \text{ is odd} \end{cases}$$

and $\alpha = \frac{\log 3}{\log 4}$. The Fourier expansion $\psi(u) = \sum_{k \in \mathbb{Z}} \psi_k e^{2k\pi i u}$ is given by

$$\psi_k = \frac{3^{\alpha_k} (-1)^k}{\alpha_k 2\sqrt{3} \log 2} \left(3 + (-1)^k \sqrt{3} - \left(1 + (-1)^k \sqrt{3} \right) f_0(\alpha_k) + \right. \\ \left. + \left(2 + (-1)^k \sqrt{3} \right) f_1(\alpha_k) - f_2(\alpha_k) \right),$$

where $\alpha_k = \alpha + \frac{k\pi i}{\log 2}$ and

$$f_0(s) = \frac{1}{3\Gamma(s)} \int_0^\infty \left((F(e^{-t}) + F(\zeta e^{-t}) + F(\zeta^2 e^{-t})) (1 - e^{-\frac{t}{2}}) - 1 \right) t^{s-1} dt$$

$$f_1(s) = \frac{1}{3\Gamma(s)} \int_0^\infty (F(e^{-t}) + \zeta^2 F(\zeta e^{-t}) + \zeta F(\zeta^2 e^{-t})) (1 - e^{-\frac{t}{2}}) t^{s-1} dt$$

$$f_2(s) = \frac{1}{3\Gamma(s)} \int_0^\infty (F(e^{-t}) + \zeta F(e^{-t}) + \zeta^2 F(\zeta^2 e^{-t})) (1 - e^{-\frac{t}{2}}) t^{s-1} dt$$

with $\zeta = e^{\frac{2\pi i}{3}}$ and

$$(7.2) \quad F(z) = \prod_{k=0}^{\infty} (1 - z^{2^k}).$$

In particular, a rough estimate of the mean value ψ_0 is

$$\psi_0 \approx 2.93039 \ 13325 \ 63657 \ 25112 \ 43191 \ 20179.$$

Sketch of the proof. The proof runs along the same lines as the proof of Theorem 5; the only difference is that the computation of the Dirichlet generating function is slightly more involved than in Section 6. We first prove an asymptotic formula for the double summatory function

$$T(N) = \sum_{1 \leq n < N} (S_3(n) - 1).$$

For this purpose we need some information on the function

$$\xi_0(s) = \sum_{n=1}^{\infty} \frac{(-1)^{\nu(3n)}}{(3n)^s}.$$

Using the function F given by (7.2) that satisfies $F(z) = (1 - z)F(z^2)$, and setting

$$\Xi_0(z) = \frac{1}{3} (F(z) + F(\zeta z) + F(\zeta^2 z))$$

$$\Xi_1(z) = \frac{1}{3} (F(z) + \zeta^2 F(\zeta z) + \zeta F(\zeta^2 z))$$

$$\Xi_2(s) = \frac{1}{3} (F(z) + \zeta F(\zeta z) + \zeta^2 F(\zeta^2 z))$$

we obtain the functional equations

$$(7.3) \quad \begin{aligned} \Xi_0(z) &= \Xi_0(z^2) - z\Xi_1(z^2) \\ \Xi_1(z) &= \Xi_2(z^2) - z\Xi_0(z^2) \\ \Xi_2(z) &= \Xi_1(z^2) - z\Xi_2(z^2). \end{aligned}$$

Consider now the companion Dirichlet series ξ_k , for $k = 1, 2$, defined in a way similar to ξ_0 , where summation runs through the other residue classes $\pmod 3$,

$$\xi_k(s) = \sum_{n=0}^{\infty} \frac{(-1)^{\nu(3n+k)}}{(3n+k)^s} \quad \text{for } k = 1, 2.$$

By Mellin transforms again, Eq. (2.2), we derive the alternative expressions,

$$(7.4) \quad \begin{aligned} \xi_0(s) &= \frac{1}{\Gamma(s)} \int_0^{\infty} (\Xi_0(e^{-t}) - 1) t^{s-1} dt \\ \xi_k(s) &= \frac{1}{\Gamma(s)} \int_0^{\infty} \Xi_k(e^{-t}) t^{s-1} dt \quad \text{for } k = 1, 2. \end{aligned}$$

The image of the collection of functional equations (7.3) is then the system of equations

$$(7.5) \quad \begin{array}{rcl} (1 - 2^{-s}) \xi_0(s) + 2^{-s} \xi_1(s) & & = 2^{-s} f_1(s) \\ 2^{-s} \xi_0(s) + \xi_1(s) - 2^{-s} \xi_2(s) & & = 2^{-s} f_0(s) - 1 \\ - 2^{-s} \xi_1(s) + (1 + 2^{-s}) \xi_2(s) & & = 2^{-s} f_2(s), \end{array}$$

where the functions f_k are given by

$$\begin{aligned} f_0(s) &= \sum_{n=1}^{\infty} (-1)^{\nu(3n)} \left(\frac{1}{(3n)^s} - \frac{1}{(3n + \frac{1}{2})^s} \right) \\ f_k(s) &= \sum_{n=0}^{\infty} (-1)^{\nu(3n+k)} \left(\frac{1}{(3n+k)^s} - \frac{1}{(3n+k + \frac{1}{2})^s} \right) \quad \text{for } k = 1, 2. \end{aligned}$$

These functions are defined for $\Re(s) > 0$ and satisfy $|f_k(\sigma + it)| \ll |t|^{1-\sigma}$ for $0 < \sigma < 1$, which can be shown using the same arguments as in Section 6.

Solving (7.5) yields

$$(7.6) \quad \xi_0(s) = \frac{1}{2^s(4^s - 3)} (4^s + 2^s - (2^s + 1) f_0(s) + (4^s + 2^s - 1) f_1(s) - f_2(s)).$$

This equation provides us with the analytic continuation of ξ_0 and shows that all poles of this function have to satisfy the equation $4^s = 3$.

After these preparations we can write using (2.4)

$$(7.7) \quad T(N) = \frac{N}{2\pi i} \int_{2-i\infty}^{2+i\infty} 3^s \xi_0(s) N^s \frac{ds}{s(s+1)}.$$

Shifting the line of integration to the left and taking residues into account yields

$$(7.8) \quad T(N) = N^{\alpha+1} \sum_{k \in \mathbb{Z}} \frac{\psi_k}{\alpha_k + 1} e^{2k\pi i \log_4 N} + O(N^{1+\varepsilon}),$$

where the term $O(N^{1+\varepsilon})$ is obtained by trivial estimation of the integral from $\varepsilon - i\infty$ to $\varepsilon + i\infty$ over the same integrand as in (7.7).

Now using the exact formula due to Coquet and a slightly modified version of Proposition 2 yields the Fourier expansion of the function ψ . \square

Remark 5. The method used above can also be used to gather information on the summatory functions

$$\sum_{k < n} (-1)^{\nu(3k+1)}, \quad \sum_{k < n} (-1)^{\nu(5k)}$$

and other functions of this type.

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