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ERROR ANALYSIS OF SOME FINITE ELEMENT METHODS FOR THE STOKES PROBLEM

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Abstract. We prove the optimal order of convergence for some two-dimensional finite element methods for the Stokes equations. First we consider methods of the Taylor-Hood type: the triangular $P_3 - P_2$ element and the $Q_k - Q_{k-1}$, $k \geq 2$, family of quadrilateral elements. Then we introduce two new low-order methods with piecewise constant approximations for the pressure. The analysis is performed using our macroelement technique, which is reviewed in a slightly altered form.

ANALYSE D'ERREUR DE QUELQUES METHODES D'ELEMENTS FINIS POUR LE PROBLEME DE STOKES

Résumé. Nous montrons la convergence optimale pour quelques méthodes d'éléments finis mixtes pour les équations de Stokes en dimension deux. Nous considérons d'abord des méthodes de type Taylor-Hood: l'élément $P_3 - P_2$ et la famille $Q_k - Q_{k-1}$, $k \geq 2$, quadrilatère. Ensuite nous introduisons deux nouvelles méthodes d'ordre inférieur avec une approximation de type constante par morceaux pour la pression. Pour l'analyse nous utilisons notre technique de macroélément et la présentons sous une forme modifiée.

1. Introduction

In this paper we will consider some finite element methods for the Stokes equations: Find the velocity \mathbf{u} and the pressure p such that

$$\begin{aligned} -\nu\Delta\mathbf{u} + \nabla p &= \mathbf{f} & \text{in } \Omega, \\ \operatorname{div} \mathbf{u} &= 0 & \text{in } \Omega, \\ \mathbf{u} &= 0 & \text{on } \partial\Omega, \end{aligned} \tag{1.1}$$

where $\Omega \subset \mathbf{R}^d$, $d = 2, 3$, is a bounded polygonal or polyhedral domain, \mathbf{f} is the given body force and $\nu > 0$ is the dynamic viscosity.

The usual variational formulation of (1.1) is the following. Find $\mathbf{u} \in H_0^1(\Omega)^d$ and $p \in L_0^2(\Omega)$ such that

$$\begin{aligned} \nu(\nabla\mathbf{u}, \nabla\mathbf{v}) - (\operatorname{div} \mathbf{v}, p) &= (\mathbf{f}, \mathbf{v}), & \mathbf{v} \in H_0^1(\Omega)^d, \\ (\operatorname{div} \mathbf{u}, q) &= 0, & q \in L_0^2(\Omega), \end{aligned} \tag{1.2}$$

where (\cdot, \cdot) denotes the inner product in $L^2(\Omega)$, $L^2(\Omega)^d$ or $L^2(\Omega)^{d \times d}$, and $L_0^2(\Omega)$ the space

$$L_0^2(\Omega) = \left\{ p \in L^2(\Omega) \mid \int_{\Omega} p \, d\mathbf{x} = 0 \right\}.$$

For $\mathbf{f} \in H^{-1}(\Omega)^d$ this problem has a unique solution; cf. [10].

The mixed method based on (1.2) reads as follows. Find $\mathbf{u}_h \in \mathbf{V}_h \subset H_0^1(\Omega)^d$ and $p_h \in P_h \subset L_0^2(\Omega)$ such that

$$\begin{aligned} \nu(\nabla\mathbf{u}_h, \nabla\mathbf{v}) - (\operatorname{div} \mathbf{v}, p_h) &= (\mathbf{f}, \mathbf{v}), & \mathbf{v} \in \mathbf{V}_h, \\ (\operatorname{div} \mathbf{u}_h, q) &= 0, & q \in P_h. \end{aligned} \tag{1.3}$$

It is well known that in order to get a working method the spaces \mathbf{V}_h and P_h cannot be chosen arbitrarily. The method can be expected to behave well only if the following "inf-sup condition" is satisfied

$$\inf_{0 \neq p \in P_h} \sup_{0 \neq \mathbf{v} \in \mathbf{V}_h} \frac{(\operatorname{div} \mathbf{v}, p)}{\|\mathbf{v}\|_1 \|p\|_0} \geq C > 0. \quad (1.4)$$

The following fundamental result is classical.

PROPOSITION 1.1. (Babuška [1,2], Brezzi [6].) *Suppose that the finite element spaces \mathbf{V}_h and P_h satisfy (1.4). Then the system (1.3) has a unique solution (\mathbf{u}_h, p_h) satisfying*

$$\|\mathbf{u} - \mathbf{u}_h\|_1 + \|p - p_h\|_0 \leq C \left\{ \inf_{\mathbf{v} \in \mathbf{V}_h} \|\mathbf{u} - \mathbf{v}\|_1 + \inf_{q \in P_h} \|p - q\|_0 \right\},$$

where (\mathbf{u}, p) is the solution to (1.2). ■

By now this field of problems is rather well understood; there exist rather general techniques for verifying the inf-sup condition [5,13] and for the construction of methods satisfying it. As a consequence, the collection of methods which are known to be stable (i.e. satisfying (1.4)) is relatively large; cf. e.g. the recent book by Girault and Raviart [10].

The purpose of this paper is to extend the list of stable methods. In section 3 we consider two families of the "Taylor-Hood" type and in the last section we introduce and analyze two new low-order methods. In the analysis we use the macroelement technique, introduced by us in [13], which we in section 2 recall in a slightly more practical form.

The results of this paper are trivially also valid when the same finite element spaces are used for the equations of (nearly) incompressible elasticity.

Our notation is standard; cf. [7,10].

2. The macroelement technique

For further reference we will here present the technique in a quite general form for both two- and three-dimensional problems, even if all the methods to be analyzed in this paper are two-dimensional. The presentation do not cover all possible cases (such as e.g. a mixing of triangles and quadrilaterals, prismatic elements,...), but the modifications needed for a method not covered are trivial.

Let \mathcal{C}_h be a finite element partitioning of $\bar{\Omega}$ into subdomains which are all assumed to be either triangles or convex quadrilaterals when $\Omega \subset \mathbf{R}^2$, and tetrahedrons or convex hexahedrons when $\Omega \subset \mathbf{R}^3$. The partitioning is assumed to satisfy the standard regularity and compatibility condition [7], but we will not assume \mathcal{C}_h to be quasiuniform. Let \hat{K} denote the reference triangle, tetrahedron, square or cube, respectively, and for $K \in \mathcal{C}_h$ denote by F_K the affine, bilinear or trilinear mapping from \hat{K} onto K . Further, let $\hat{\mathbf{V}}$ and \hat{P} be two polynomial spaces defined on \hat{K} . We now assume that \mathbf{V}_h and P_h are defined as

$$\mathbf{V}_h = \{ \mathbf{v} \in H_0^1(\Omega)^d \mid \mathbf{v}(\mathbf{x}) = \hat{\mathbf{v}}(F_K^{-1}(\mathbf{x})), \hat{\mathbf{v}} \in \hat{\mathbf{V}}, K \in \mathcal{C}_h \}, \quad (2.1)$$

and

$$P_h = \{ p \in L_0^2(\Omega) \mid p(\mathbf{x}) = \hat{p}(F_K^{-1}(\mathbf{x})), \hat{p} \in \hat{P}, K \in \mathcal{C}_h \} \quad (2.2a)$$

or

$$P_h = \{ p \in C(\Omega) \cap L_0^2(\Omega) \mid p(\mathbf{x}) = \hat{p}(F_K^{-1}(\mathbf{x})), \hat{p} \in \hat{P}, K \in \mathcal{C}_h \}. \quad (2.2b)$$

Further, we make the assumption that \mathbf{V}_h contains the piecewise linear, isoparametric bilinear or isoparametric trilinear functions for the triangular and tetrahedral, quadrilateral or hexahedral case, respectively. With the choice (2.2a) P_h is assumed to contain the piecewise constants, whereas for (2.2b) it contains the piecewise linear, isoparametric bilinear or isoparametric trilinear functions.

A method where the pressure space is defined according to (2.2b) is usually referred to as a "Taylor-Hood" method.

By a macroelement M we define a connected set of elements of which the intersection of any two is either empty, a vertex or one edge or face in \mathbf{R}^2 and \mathbf{R}^3 , respectively. Further two macroelements M and \tilde{M} are said to be equivalent if they can be mapped continuously onto each other, or more precisely if one can define a mapping $G : \tilde{M} \rightarrow M$ such that:

(i) $G(\tilde{M}) = M$.

(ii) If $\tilde{M} = \cup_{j=1}^m \tilde{K}_j$, where \tilde{K}_j , $j = 1, 2, \dots, m$, are the elements of \tilde{M} ,
then $K_j = G(\tilde{K}_j)$, $j = 1, 2, \dots, m$, are the elements of M .

(iii) $G|_{\tilde{K}_j} = F_{K_j} \circ F_{\tilde{K}_j}^{-1}$, $j = 1, 2, \dots, m$, where F_{K_j} and $F_{\tilde{K}_j}$ are the mappings from the reference element \hat{K} onto K_j and \tilde{K}_j , respectively.

For a macroelement M we define the spaces

$$\mathbf{V}_{0,M} = \{ \mathbf{v} \in H_0^1(M)^d \mid \mathbf{v}(\mathbf{x}) = \hat{\mathbf{v}}(F_K^{-1}(\mathbf{x})), \hat{\mathbf{v}} \in \hat{\mathbf{V}}, \mathbf{x} \in K, K \subset M \} \quad (2.3)$$

and

$$P_M = \{ p \in L^2(M) \mid p(\mathbf{x}) = \hat{p}(F_K^{-1}(\mathbf{x})), \hat{p} \in \hat{P}, \mathbf{x} \in K, K \subset M \} \quad (2.4a)$$

or

$$P_M = \{ p \in C(M) \mid p(\mathbf{x}) = \hat{p}(F_K^{-1}(\mathbf{x})), \hat{p} \in \hat{P}, \mathbf{x} \in K, K \subset M \} \quad (2.4b)$$

depending on which of the possibilities (2.2ab) is chosen. Further, we denote

$$N_M = \{ p \in P_M \mid (\operatorname{div} \mathbf{v}, p)_M = 0, \mathbf{v} \in \mathbf{V}_{0,M} \}.$$

The collection of edges or faces, of the elements of \mathcal{C}_h , in the interior of Ω is denoted by Γ_h .

The macroelement technique is given by the following

THEOREM 2.1. *Suppose that there is a fixed set of equivalence classes \mathcal{E}_i , $i=1,2,\dots,k$, of macroelements, a positive integer L and a macroelement partitioning \mathcal{M}_h such that:*

(M1) *For each $M \in \mathcal{E}_i$, $i=1,2,\dots,k$, the space N_M is one-dimensional consisting of functions that are constant on M .*

(M2) *Each $M \in \mathcal{M}_h$ belongs to one of the classes \mathcal{E}_i , $i=1,2,\dots,k$.*

(M3) *Each $K \in \mathcal{C}_h$ is contained in at least one and not more than L macroelements of \mathcal{M}_h .*

(M4) *Each $T \in \Gamma_h$ is contained in at least one and not more than L macroelements of \mathcal{M}_h .*

Then the stability inequality (1.4) is valid.

PROOF: The proof will consist of a modification of some arguments given in [11,13,15].

We will need the following norm defined in P_h :

$$\|p\|_h^2 = \sum_{K \in \mathcal{C}_h} h_K^2 \|\nabla p\|_{0,K}^2 + \sum_{T \in \Gamma_h} h_T \int_T |[p]|^2 ds,$$

where $([p])|_T$ denotes the jump in p along the inter element boundary T and h_T stands for the length or area of T . In a macroelement M we similarly define

$$|p|_M^2 = \sum_{K \subset M} h_K^2 \|\nabla p\|_{0,K}^2 + \sum_{T \in \Gamma_M} h_T \int_T |[p]|^2 ds,$$

where Γ_M denotes the inter element boundaries in the interior of M .

Now, since we assume that

- there is only a fixed number of different types of macroelements,
- all the elements in \mathcal{C}_h are regular

and

- the condition (M1) is satisfied,

one can prove (see [13, Lemma 3.1]) that there is a constant C independent of \mathcal{M}_h such that

$$\inf_{0 \neq p \in P_M} \sup_{0 \neq \mathbf{v} \in \mathbf{V}_{0,M}} \frac{(\operatorname{div} \mathbf{v}, p)_M}{|\mathbf{v}|_{1,M} |p|_M} \geq C > 0, \quad (2.5)$$

or alternatively stated: For a given $p \in P_h$ and $M \in \mathcal{M}_h$ there is a $\mathbf{v}_M \in \mathbf{V}_h$, with $v_M = 0$ in $\Omega \setminus M$, such that

$$(\operatorname{div} \mathbf{v}_M, p)_M \geq C |p|_M^2 \quad (2.6)$$

and

$$|\mathbf{v}_M|_{1,M} \leq |p|_M. \quad (2.7)$$

We now define $\mathbf{v} \in \mathbf{V}_h$ through

$$\mathbf{v} = \sum_{M \in \mathcal{M}_h} \mathbf{v}_M.$$

This gives

$$(\operatorname{div} \mathbf{v}, p) = \sum_{M \in \mathcal{M}_h} (\operatorname{div} \mathbf{v}_M, p)_M \geq C \sum_{M \in \mathcal{M}_h} |p|_M^2 \geq C \|p\|_h^2$$

and

$$\|\mathbf{v}\|_1 \leq C |\mathbf{v}|_1 \leq C \sum_{M \in \mathcal{M}_h} |\mathbf{v}_M|_{1,M} \leq C \sum_{M \in \mathcal{M}_h} |p|_M \leq CL \|p\|_h, \quad (2.10)$$

i.e. there is a constant $C_1 > 0$ such that

$$\sup_{0 \neq \mathbf{v} \in \mathbf{V}_h} \frac{(\operatorname{div} \mathbf{v}, p)}{\|\mathbf{v}\|_1} \geq C_1 \|p\|_h. \quad (2.11)$$

Next, let us use an argument due to Verfürth [15, Proposition 3.3] for showing that (2.11) implies the corresponding condition with the L^2 -norm for the pressure, i.e. the inf-sup condition (1.4). We will show that the argument can be used also when P_h consists of discontinuous functions and that the quasiuniformity assumption of [15] can be avoided.

Now, for every $p \in P_h$ there is a $\mathbf{w} \in H_0^1(\Omega)^d$ such that

$$(\operatorname{div} \mathbf{w}, p) \geq C_2 \|p\|_0^2 \quad (2.12)$$

and

$$\|\mathbf{w}\|_1 \leq \|p\|_0. \quad (2.13)$$

Further, one can show (cf. [8], [3, Lemma 3] and [10, pp. 109-111]) that there is an interpolant $\tilde{\mathbf{w}} \in \mathbf{V}_h$ to \mathbf{w} such that

$$\left(\sum_{K \in \mathcal{C}_h} h_K^{-2} \|\mathbf{w} - \tilde{\mathbf{w}}\|_{0,K}^2 + \sum_{T \in \Gamma_h} h_T^{-1} \int_T |\mathbf{w} - \tilde{\mathbf{w}}|^2 ds \right)^{1/2} \leq C_3 |\mathbf{w}|_1 \quad (2.14)$$

and

$$\|\tilde{\mathbf{w}}\|_1 \leq C_4 \|\mathbf{w}\|_1. \quad (2.15)$$

Using (2.12), (2.14) and (2.13) we now get

$$\begin{aligned} (\operatorname{div} \tilde{\mathbf{w}}, p) &= (\operatorname{div} (\tilde{\mathbf{w}} - \mathbf{w}), p) + (\operatorname{div} \mathbf{w}, p) \\ &\geq (\operatorname{div} (\tilde{\mathbf{w}} - \mathbf{w}), p) + C_2 \|p\|_0^2 \\ &= \sum_{K \in \mathcal{C}_h} (\mathbf{w} - \tilde{\mathbf{w}}, \nabla p)_K + \sum_{T \in \Gamma_h} \int_T ((\tilde{\mathbf{w}} - \mathbf{w}) \cdot \mathbf{n})([p]) ds + C_2 \|p\|_0^2 \\ &\geq - \left(\sum_{K \in \mathcal{C}_h} h_K^{-2} \|\mathbf{w} - \tilde{\mathbf{w}}\|_{0,K}^2 + \sum_{T \in \Gamma_h} h_T^{-1} \int_T |\mathbf{w} - \tilde{\mathbf{w}}|^2 ds \right)^{1/2} \cdot \|p\|_h + C_2 \|p\|_0^2 \\ &\geq -C_3 |\mathbf{w}|_1 \|p\|_h + C_2 \|p\|_0^2 \\ &\geq -C_3 \|p\|_0 \|p\|_h + C_2 \|p\|_0^2 \\ &= \|p\|_0^2 \cdot \left(C_2 - C_3 \frac{\|p\|_h}{\|p\|_0} \right), \end{aligned}$$

which combined with (2.13) and (2.15) gives

$$\sup_{0 \neq \mathbf{v} \in \mathbf{V}_h} \frac{(\operatorname{div} \mathbf{v}, p)}{\|\mathbf{v}\|_1} \geq \|p\|_0 \cdot \left(C_5 - C_6 \frac{\|p\|_h}{\|p\|_0} \right). \quad (2.16)$$

On the other hand, (2.11) can be written as

$$\sup_{0 \neq \mathbf{v} \in \mathbf{V}_h} \frac{(\operatorname{div} \mathbf{v}, p)}{\|\mathbf{v}\|_1} \geq C_1 \|p\|_h = \|p\|_0 \cdot C_1 \frac{\|p\|_h}{\|p\|_0}. \quad (2.17)$$

Hence, combining (2.16) and (2.17) we get

$$\inf_{0 \neq p \in P_h} \sup_{0 \neq \mathbf{v} \in \mathbf{V}_h} \frac{(\operatorname{div} \mathbf{v}, p)}{\|\mathbf{v}\|_1 \|p\|_0} \geq \max_{t > 0} \{ C_1 t, C_5 - C_6 t \} \geq \frac{C_1 C_5}{C_1 + C_6}. \quad \blacksquare$$

Let us close this section by giving some remarks on the difference between the formulation of the macroelement technique given in [13] and that of this paper. In [13] the macroelement partitioning \mathcal{M}_h was chosen so that each $K \in \mathcal{C}_h$ was a subset of exactly one $M \in \mathcal{M}_h$. Then the condition (M1) ensures that we can stabilize all pressure components except those which consists of constants on each macroelement. To stabilize the piecewise constant pressure components we clearly need velocity degrees of freedom in the interior of the boundaries between the macroelements. More precisely, we included the condition that for any two neighboring macroelements M_1, M_2 of \mathcal{M}_h with $\int_{M_1 \cap M_2} ds \neq 0$ there is a $\mathbf{v} \in \mathbf{V}_h$ such that

$$\text{supp } \mathbf{v} \subset M_1 \cup M_2 \quad \text{and} \quad \int_{M_1 \cap M_2} \mathbf{v} \cdot \mathbf{n} \, ds \neq 0.$$

The disadvantage of this choice is that in order to find a macroelement partitioning of nonoverlapping macroelements with this property it is often necessary to use many different kinds of macroelements or macroelements that are big. In these cases one could claim that the asymptotical stability inequality so obtained is not valid for most values of the mesh parameter used in practice. In many cases it is not even evident how such a macroelement partitioning should be obtained. A good example showing these problems is the tetrahedral Taylor-Hood method. In [14] we analyzed this method with the technique of this paper. The present technique also shows more clearly that the condition that the space N_M consists of constants is really the only one that has to be verified.

Let us now use the above theorem for the analysis of some mixed methods not earlier analyzed in the literature.

3. Two families of methods

Let $\Omega \subset \mathbf{R}^2$ and let \mathcal{C}_h be a partitioning the elements of which all are assumed to be either triangles or quadrilaterals. For the index $k \geq 2$ we define the families

$$\begin{aligned} \mathbf{V}_h &= \{ \mathbf{v} \in H_0^1(\Omega)^2 \mid \mathbf{v}|_K \in R_k(K)^2, \quad K \in \mathcal{C}_h \}, \\ P_h &= \{ p \in L_0^2(\Omega) \cap C(\Omega) \mid p|_K \in R_{k-1}(K), \quad K \in \mathcal{C}_h \}, \end{aligned} \tag{3.1}$$

where $R_m(K) = P_m(K)$ when K is a triangle and $R_m(K) = Q_m(K)$ when K is a quadrilateral. $P_m(K)$ and $Q_m(K)$ are the usual polynomial spaces on K ; cf. [7].

For these families we have the following optimal error estimates.

THEOREM 3.1. *For the solution (\mathbf{u}_h, p_h) of (1.3) with the method (3.1) we have*

$$\|\mathbf{u} - \mathbf{u}_h\|_1 + \|p - p_h\|_0 \leq Ch^k(|\mathbf{u}|_{k+1} + |p|_k).$$

For a convex domain Ω we additionally have

$$\|\mathbf{u} - \mathbf{u}_h\|_0 \leq Ch^{k+1}(|\mathbf{u}|_{k+1} + |p|_k). \quad \blacksquare$$

For the triangular case the verification of the above estimates consists merely of filling in some gaps; for $k \geq 4$ the result is covered by the analysis in [12] and for $k = 2$ the result is well-known [4,15], but under the restrictions that any $K \in \mathcal{C}_h$ do not have two edges on the boundary $\partial\Omega$ and that \mathcal{C}_h is quasiuniform. We will prove the stability for the case $k = 3$. Exactly the same line of reasoning applies to the case $k = 2$ and hence the above mentioned restrictions on the mesh are unnecessary.

REMARK. In [14] we analyzed the corresponding three-dimensional tetrahedral method for $k = 2$. There the restriction on the mesh was that each tetrahedron in \mathcal{C}_h has at least one vertex in the interior of Ω . An inspection shows, however, that this restriction cannot be dropped. \blacksquare

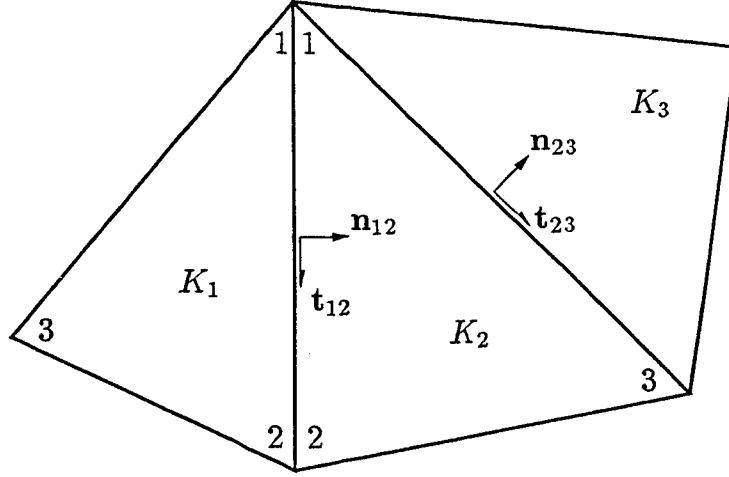
LEMMA 3.1. *Let M be a macroelement consisting of three triangles and define*

$$\mathbf{V}_{0,M} = \{ \mathbf{v} \in H_0^1(M)^2 \mid \mathbf{v}|_K \in P_3(K)^2, \quad K \subset M \},$$

$$P_M = \{ p \in C(M) \mid p|_K \in P_2(K), \quad K \subset M \}.$$

Then the corresponding nullspace N_M is one-dimensional consisting of functions that are constant on M .

PROOF: Let $M = K_1 \cup K_2 \cup K_3$ as in the figure below and let $\mathbf{t}_{12}, \mathbf{t}_{23}$ and $\mathbf{n}_{12}, \mathbf{n}_{23}$ be the tangents and normals, respectively, to the common edges.



Let us first consider $K_1 \cup K_2$ with the "local numbering" of the vertices as in the figure. For $p \in P_M$ we have $\nabla p|_{K_j} \in P_1(K_j)^2$, $j = 1, 2$, and the component $\nabla p \cdot \mathbf{t}_{12}$ is continuous in the whole of $K_1 \cup K_2$. Hence we can write

$$(\nabla p \cdot \mathbf{t}_{12})|_{K_j} = \sum_{i=1}^3 \alpha_{ij} \lambda_{ij}, \quad j = 1, 2,$$

where λ_{ij} , $i = 1, 2, 3$, are the barycentric coordinates of K_j , $j = 1, 2$. Since $\nabla p \cdot \mathbf{t}_{12}$ is continuous in $K_1 \cup K_2$ we have $\alpha_{i1} = \alpha_{i2}$ for $i = 1, 2$. Now, choose $\mathbf{u} \in V_{0, K_1 \cup K_2}$ (i.e. $\mathbf{u} = 0$ in K_3) such that $\mathbf{u} \cdot \mathbf{n}_{12} = 0$ in $K_1 \cup K_2$ and

$$(\mathbf{u} \cdot \mathbf{t}_{12})|_{K_j} = \lambda_{1j} \lambda_{2j} (3\lambda_{1j} - 1), \quad j = 1, 2.$$

Then we have

$$\begin{aligned} (\operatorname{div} \mathbf{u}, p)_M &= -(\mathbf{u}, \nabla p)_M = -\sum_{j=1}^2 \int_{K_j} (\mathbf{u} \cdot \mathbf{t}_{12})(\nabla p \cdot \mathbf{t}_{12}) \, d\mathbf{x} \\ &= -\sum_{j=1}^2 \sum_{i=1}^3 \alpha_{ij} \int_{K_j} \lambda_{1j} \lambda_{2j} (3\lambda_{1j} - 1) \lambda_{ij} \, d\mathbf{x} \\ &= -\frac{\operatorname{area}(K_1)}{60} \alpha_{11} - \frac{\operatorname{area}(K_2)}{60} \alpha_{12} \\ &= -\left(\frac{\operatorname{area}(K_1)}{60} + \frac{\operatorname{area}(K_2)}{60} \right) \alpha_{12} \end{aligned}$$

and the condition $(\operatorname{div} \mathbf{u}, p)_M = 0$ implies that

$$\alpha_{11} = \alpha_{12} = 0.$$

By symmetric reasoning we get

$$\alpha_{21} = \alpha_{22} = 0.$$

Next we choose $\mathbf{u} \in \mathbf{V}_{0, K_j}$, $j = 1, 2$, such that $\mathbf{u} \cdot \mathbf{n}_{12} = 0$ and

$$(\mathbf{u} \cdot \mathbf{t}_{12})|_{K_j} = \lambda_{1j} \lambda_{2j} \lambda_{3j}, \quad j = 1, 2.$$

If $p \in N_M$ then we have

$$\begin{aligned} 0 &= (\operatorname{div} \mathbf{u}, p)_M = -(\mathbf{u}, \nabla p)_M \\ &= \alpha_{3j} \int_{K_j} \lambda_{1j} \lambda_{2j} \lambda_{3j}^2 dx \\ &= \frac{\operatorname{area}(K_j)}{180} \alpha_{3j}, \quad j = 1, 2, \end{aligned}$$

i.e.

$$\alpha_{3j} = 0, \quad j = 1, 2.$$

We have thus shown that

$$\nabla p \cdot \mathbf{t}_{12} = 0 \quad \text{in } K_1 \cup K_2.$$

The same reasoning gives that

$$\nabla p \cdot \mathbf{t}_{23} = 0 \quad \text{in } K_2 \cup K_3,$$

and hence we have

$$\nabla p = 0 \quad \text{in } K_2,$$

i.e. p is a constant in K_2 .

Write next

$$(\nabla p \cdot \mathbf{n}_{12})|_{K_1} = \sum_{i=1}^3 \beta_i \lambda_{i1}$$

and take $\mathbf{u} \in \mathbf{V}_{0, K_1 \cup K_2}$ such that $\mathbf{u} \cdot \mathbf{t}_{12} = 0$ and

$$(\mathbf{u} \cdot \mathbf{n}_{12})|_{K_j} = \lambda_{1j} \lambda_{2j} (3\lambda_{1j} - 1), \quad j = 1, 2.$$

Now since $\nabla p|_{K_2} = 0$ the condition for N_M yields

$$\begin{aligned} 0 &= (\operatorname{div} \mathbf{u}, p)_M = -(\mathbf{u}, \nabla p)_{K_1} \\ &= -\sum_{i=1}^3 \beta_i \int_{K_1} \lambda_{11} \lambda_{22} (3\lambda_{11} - 1) \lambda_{i1} \, d\mathbf{x} \\ &= -\frac{\operatorname{area}(K_1)}{60} \beta_1, \end{aligned}$$

i.e. $\beta_1 = 0$, and analogously we get $\beta_2 = 0$. Upon choosing $\mathbf{u} \in \mathbf{V}_{0,K_1}$ such that $\mathbf{u} \cdot \mathbf{t}_{12} = 0$ and $(\mathbf{u} \cdot \mathbf{n}_{12})|_{K_1} = \lambda_{11} \lambda_{21} \lambda_{31}$ we get

$$0 = (\operatorname{div} \mathbf{u}, p)_M = (\operatorname{div} \mathbf{u}, p)_{K_1} = -\frac{\operatorname{area}(K_1)}{180} \beta_3,$$

i.e. $\beta_3 = 0$ in K_1 . We thus have $\nabla p \cdot \mathbf{n}_{12} = 0$ in K_1 and hence p is a constant in K_1 .

Analogously $p \in N_M$ has to be a constant in K_3 .

Since p is a constant in K_i , $i = 1, 2, 3$, and continuous, it is constant in the whole of M . ■

When applying Theorem 2.1 we now use only one class of macroelements and the partitioning \mathcal{M}_h can be obtained by e.g. taking one macroelement for each element edge in the interior of Ω . Then the assumptions of Theorem 2.1 are valid and standard approximation theory gives the first estimate of Theorem 3.1. The L^2 -estimate for the velocity follows by the standard Aubin-Nitsche trick.

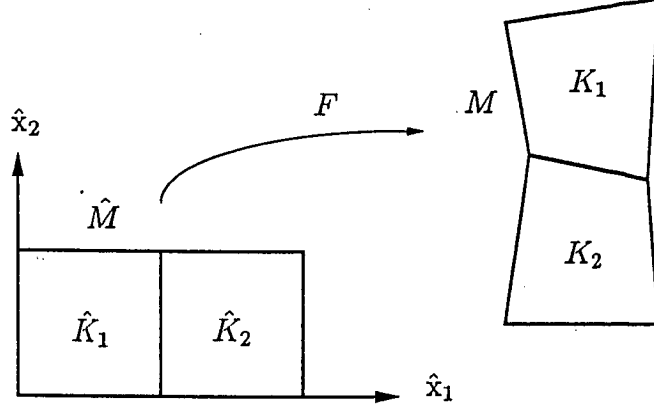
Let us turn to the quadrilateral family. To our knowledge only the lowest order method in this family has earlier been analyzed [4,13].

LEMMA 3.2. *Let M be a macroelement consisting of two quadrilaterals and define*

$$\begin{aligned} \mathbf{V}_{0,M} &= \{ \mathbf{v} \in H_0^1(M)^2 \mid \mathbf{v}|_K \in Q_k(K)^2, \, K \subset M \}, \\ P_M &= \{ p \in C(M) \mid p|_K \in Q_{k-1}(K), \, K \subset M \}. \end{aligned}$$

Then the corresponding nullspace N_M is one-dimensional consisting of functions that are constant on M .

PROOF: Let $M = K_1 \cup K_2$ and $\hat{M} = \hat{K}_1 \cup \hat{K}_2$ be the macroelement and the reference macroelement, respectively, and denote by $F = (F_1, F_2)$ the continuous piecewise bilinear mapping from \hat{M} onto M , see the figure below.



For $\mathbf{u} \in \mathbf{V}_{0,M}$ and $p \in N_M$ we now have

$$\begin{aligned} 0 &= -(\operatorname{div} \mathbf{u}, p)_M = (\mathbf{u}, \nabla p)_M \\ &= \sum_{i=1}^2 \int_{\hat{K}_i} \hat{\mathbf{u}}(\hat{\mathbf{x}})^T J_F^{-T}(\hat{\mathbf{x}}) \nabla \hat{p}(\hat{\mathbf{x}}) |J_F(\hat{\mathbf{x}})| d\hat{\mathbf{x}}, \end{aligned}$$

where J_F is the Jacobian matrix of F , J_F^{-T} is the transpose of J_F^{-1} and $|J_F|$ is the determinant of J_F . Above $\hat{\mathbf{u}}(\hat{\mathbf{x}})$ and $\nabla \hat{p}(\hat{\mathbf{x}})$ are considered as column vectors. We now have

$$|J_F(\hat{\mathbf{x}})| J_F^{-T}(\hat{\mathbf{x}}) = \begin{pmatrix} \partial_2 F_2(\hat{\mathbf{x}}) & -\partial_1 F_2(\hat{\mathbf{x}}) \\ -\partial_2 F_1(\hat{\mathbf{x}}) & \partial_1 F_1(\hat{\mathbf{x}}) \end{pmatrix}. \quad (3.2)$$

Since F_1 and F_2 are bilinear this shows that

$$|J_F(\hat{\mathbf{x}})| J_F^{-T}(\hat{\mathbf{x}}) \nabla \hat{p}(\hat{\mathbf{x}})|_{\hat{K}_i} \in Q_{k-1}(\hat{K}_i)^2, \quad i = 1, 2,$$

and hence

$$\hat{\mathbf{u}}(\hat{\mathbf{x}})^T J_F^{-T}(\hat{\mathbf{x}}) \nabla \hat{p}(\hat{\mathbf{x}}) |J_F(\hat{\mathbf{x}})| \in Q_{2k-1}(\hat{K}_i), \quad i = 1, 2.$$

This means that the integrals

$$\int_{\hat{K}_i} \hat{\mathbf{u}}(\hat{\mathbf{x}})^T J_F^{-T}(\hat{\mathbf{x}}) \nabla \hat{p}(\hat{\mathbf{x}}) |J_F(\hat{\mathbf{x}})| d\hat{\mathbf{x}}, \quad i = 1, 2,$$

are exactly evaluated by the composite $(k + 1)$ -point Gauss-Lobatto formula.

On the other hand we can for the nodal degrees of freedom of \mathbf{u} choose the values at the $(k + 1)^2$ Gauss-Lobatto points in \hat{K}_i , $i = 1, 2$.

Denote by $0 = a_0, a_1, \dots, a_k = 1$ the Gauss-Lobatto points in the unit interval, so that the Gauss-Lobatto points in the two squares are given by

$$\hat{\mathbf{x}}_{ij}^1 = (a_i, a_j), \quad i, j = 0, 1, \dots, k,$$

and

$$\hat{\mathbf{x}}_{ij}^2 = (a_i + 1, a_j), \quad i, j = 0, 1, \dots, k.$$

If we now choose $\hat{\mathbf{u}}$ such that its only non-vanishing degrees of freedom is successively one of the components at the Gauss-Lobatto points in the interior of \hat{K}_1 and \hat{K}_2 we get

$$J_F^{-T}(\hat{\mathbf{x}}_{ij}^m) \nabla \hat{p}(\hat{\mathbf{x}}_{ij}^m) |J_F(\hat{\mathbf{x}}_{ij}^m)| = 0, \quad i, j = 1, 2, \dots, k - 1, \quad m = 1, 2.$$

Since $J_F(\hat{\mathbf{x}}) \neq 0$ for all $\hat{\mathbf{x}}$, this shows that

$$\nabla \hat{p}(\hat{\mathbf{x}}_{ij}^m) = 0, \quad i, j = 1, 2, \dots, k - 1, \quad m = 1, 2. \quad (3.3)$$

Next, by choosing $\hat{\mathbf{u}}$ so that the only non-vanishing degree of freedom is one of the components at $\hat{\mathbf{x}}_{kj}^1 = \hat{\mathbf{x}}_{0j}^2$, $j = 1, 2, \dots, k - 1$, we get the condition

$$J_F^{-T}(\hat{\mathbf{x}}_{kj}^1) \nabla \hat{p}(\hat{\mathbf{x}}_{kj}^1) |J_F(\hat{\mathbf{x}}_{kj}^1)| + J_F^{-T}(\hat{\mathbf{x}}_{0j}^2) \nabla \hat{p}(\hat{\mathbf{x}}_{0j}^2) |J_F(\hat{\mathbf{x}}_{0j}^2)| = 0, \quad (3.4)$$

$$j = 1, 2, \dots, k - 1.$$

Here the values at the points $\hat{\mathbf{x}}_{kj}^1$, $j = 1, 2, \dots, k - 1$, are the limits

$$\lim_{\substack{\hat{\mathbf{x}} \rightarrow \hat{\mathbf{x}}_{kj}^1 \\ \hat{\mathbf{x}} \in \hat{K}_1}} J_F^{-T}(\hat{\mathbf{x}}) \nabla \hat{p}(\hat{\mathbf{x}}) |J_F(\hat{\mathbf{x}})|$$

and analogously for $\hat{\mathbf{x}}_{0j}^2$, $j = 1, 2, \dots, k - 1$. At the common edge $\hat{K}_1 \cap \hat{K}_2$ $\partial_2 \hat{p}$ is continuous, and an inspection shows that this is also the case with $\partial_2 F_i$, $i = 1, 2$. Hence, using (3.2) the condition (3.4) reduces to

$$\begin{pmatrix} \partial_2 F_2(\hat{\mathbf{x}}_{kj}^1) & -\partial_1 F_2(\hat{\mathbf{x}}_{kj}^1) - \partial_1 F_2(\hat{\mathbf{x}}_{0j}^2) \\ -\partial_2 F_1(\hat{\mathbf{x}}_{kj}^1) & \partial_1 F_1(\hat{\mathbf{x}}_{kj}^1) + \partial_1 F_1(\hat{\mathbf{x}}_{0j}^2) \end{pmatrix} \begin{pmatrix} \partial_1 \hat{p}(\hat{\mathbf{x}}_{kj}^1) + \partial_1 \hat{p}(\hat{\mathbf{x}}_{0j}^2) \\ \partial_2 \hat{p}(\hat{\mathbf{x}}_{kj}^1) \end{pmatrix} = 0. \quad (3.5)$$

Further, at a point $\hat{\mathbf{x}} \in \hat{K}_1 \cup \hat{K}_2$ the vectors

$$(\partial_2 F_2(\hat{\mathbf{x}}), -\partial_2 F_1(\hat{\mathbf{x}}))^T \text{ and } (-\partial_1 F_2(\hat{\mathbf{x}}), \partial_1 F_1(\hat{\mathbf{x}}))^T$$

are the normals to $(\partial_2 F_1(\hat{\mathbf{x}}), \partial_2 F_2(\hat{\mathbf{x}}))^T$ and $(\partial_1 F_1(\hat{\mathbf{x}}), \partial_1 F_2(\hat{\mathbf{x}}))^T$, respectively. Hence, the determinant of the coefficient matrix in (3.5) is nonvanishing and we get

$$\partial_1 \hat{p}(\hat{\mathbf{x}}_{kj}^1) + \partial_1 \hat{p}(\hat{\mathbf{x}}_{0j}^2) = 0, \quad j = 1, 2, \dots, k-1,$$

and

$$\partial_2 \hat{p}(\hat{\mathbf{x}}_{kj}^1) = 0, \quad j = 1, 2, \dots, k-1. \quad (3.6)$$

Now (3.3) and (3.6) shows that $\partial_2 \hat{p}$ vanishes at the points

$$\hat{\mathbf{x}}_{ij}^1, \quad i = 1, 2, \dots, k, \quad j = 1, 2, \dots, k-1,$$

and

$$\hat{\mathbf{x}}_{ij}^2, \quad i = 0, 1, \dots, k-1, \quad j = 1, 2, \dots, k-1.$$

Since $\partial_2 \hat{p}|_{\hat{K}_i} \in Q_{k-1, k-2}(\hat{K}_i)$, $i = 1, 2$, (i.e. $\partial_2 \hat{p}|_{\hat{K}_i}$ is of the form $\sum_{i=0}^{k-1} \sum_{j=0}^{k-2} a_{ij} \hat{x}_1^i \hat{x}_2^j$) and the above set of points are unisolvent in $Q_{k-1, k-2}(\hat{K}_i)$, $i = 1, 2$, we conclude that

$$\partial_2 \hat{p} = 0 \text{ in } \hat{K}_1 \cup \hat{K}_2. \quad (3.7)$$

As a consequence, $\hat{p}|_{\hat{K}_i}$, $i = 1, 2$, is a polynomial of degree $k-1$ in the \hat{x}_1 -variable alone.

This means that (3.3) also implies

$$\partial_1 \hat{p} = 0 \text{ in } \hat{K}_1 \cup \hat{K}_2. \quad (3.8)$$

Due to the continuity of \hat{p} , (3.7) and (3.8) implies that \hat{p} is a constant in \hat{M} .

The assertion is thus proved. ■

In the macroelement partitioning we now again choose one macroelement for each element edge in the interior of Ω . Hence Proposition 1.1 and the Aubin-Nitsche trick gives the estimates of Theorem 3.1.

4. Two low-order methods

Let again $\Omega \subset \mathbf{R}^2$ and let \mathcal{C}_h be a regular finite element partitioning of $\bar{\Omega}$ into triangles or quadrilaterals. (It will become clear that one also could use a mesh with both triangles and quadrilaterals). The finite element method is now defined through

$$P_h = \{ p \in L_0^2(\Omega) \mid p|_K \in P_0(K), K \in \mathcal{C}_h \}$$

and

$$\mathbf{V}_h = \{ \mathbf{v} = (v_1, v_2) \in H_0^1(\Omega)^2 \mid v_{1|K} \in R_1(K), v_{2|K} \in S(K), K \in \mathcal{C}_h \},$$

where $R_1(K) = P_1(K)$ if K is a triangle and $R_1(K) = Q_1(K)$ for a quadrilateral. For triangles we let $S(K) = P_2(K)$ and for quadrilaterals we define

$$S(K) = \{ g \mid g(\mathbf{x}) = \hat{g}(F_K^{-1}(\mathbf{x})), \hat{g} \in Q_2'(\hat{K}) \}$$

where F_K is the bilinear mapping from the reference square \hat{K} onto K . and $Q_2'(\hat{K})$ is the space of "reduced biquadratic polynomials" (i.e. the quadratic serendipity element) as defined in e.g. [7,p. 63]. For a macroelement we analogously define

$$\mathbf{V}_{0,M} = \{ \mathbf{v} = (v_1, v_2) \in H_0^1(M)^2 \mid v_{1|K} \in R_1(K), v_{2|K} \in S(K), K \subset M \}$$

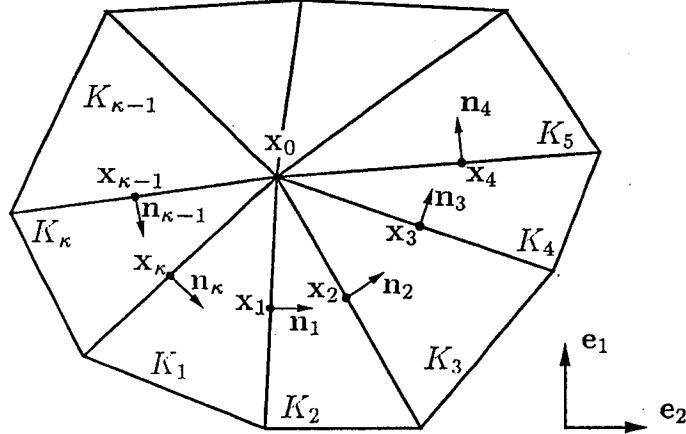
and

$$P_M = \{ p \in L^2(M) \mid p|_K \in P_0(K), K \subset M \}.$$

Let us now verify the following

LEMMA 4.1. Let M be a macroelement consisting of elements which all have one common vertex in the interior of M . Then N_M consists of functions that are constant on M .

PROOF: Let K_i , $i = 1, 2, \dots, \kappa$, be the elements of M and denote $p_i = p|_{K_i}$, $i = 1, 2, \dots, \kappa$. Further, let \mathbf{x}_i and \mathbf{n}_i , $i = 1, 2, \dots, \kappa$, be the the midpoints and normals, respectively, to the edges in M . The common vertex we denote by \mathbf{x}_0 , see the figure below for the triangular case.



The natural degrees of freedom for $\mathbf{u} \in \mathbf{V}_{0,M}$ are the values of both components of \mathbf{u} at \mathbf{x}_0 and the values of u_2 at \mathbf{x}_i , $i = 1, 2, \dots, \kappa$. By taking $\mathbf{u} \in \mathbf{V}_{0,M}$ such that the only nonvanishing degree of freedom is $u_2(\mathbf{x}_i)$, the condition $(\operatorname{div} \mathbf{u}, p)_M = 0$ implies that $p_i = p_{i+1}$ (with $p_{\kappa+1} = p_0$) if $\mathbf{n}_i \cdot \mathbf{e}_2 \neq 0$ where $\mathbf{e}_2 = (0, 1)$. Hence, the space N_M can be at most two-dimensional, and this happens only if two of the edges are parallel to \mathbf{e}_2 . But in this case one chooses \mathbf{u} such that the only non-zero degree of freedom is $u_1(\mathbf{x}_0)$. The condition for N_M then forces p to be constant on the whole of M . ■

If we now impose the restriction on \mathcal{C}_h that every $K \in \mathcal{C}_h$ has at least one vertex in the interior of Ω , then we can easily construct a macroelement partitioning \mathcal{M}_h satisfying (M1)-(M4): For each interior vertex of the mesh one takes one macroelement consisting of all the elements which have this vertex in common.

Hence we get the following

THEOREM 4.1. *Suppose that every $K \in C_h$ has at least one vertex in the interior of Ω . Then we have*

$$\|\mathbf{u} - \mathbf{u}_h\|_1 + \|p - p_h\|_0 \leq Ch(|\mathbf{u}|_2 + |p|_1).$$

For a convex domain Ω we additionally have

$$\|\mathbf{u} - \mathbf{u}_h\|_0 \leq Ch^2(|\mathbf{u}|_2 + |p|_1). \quad \blacksquare$$

REMARK. The elements we have presented here are rather similar to two elements by Fortin [9]. In Fortin's methods piecewise constants are used for the pressure whereas the basic velocity spaces of linear/bilinear functions are augmented with quadratic functions which have as degrees of freedom the normal components of the velocity at the midpoints of the element edges. Hence, our methods are simpler to implement. The calculation of the element matrices is faster and the resulting linear system is more sparse. This property could be an advantage when iterative solution methods are used. \blacksquare

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