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Phase diagram of a long bar under a higher-order active contour energy: application to hydrographic network extraction from VHR satellite images

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Abstract

The segmentation of networks is important in several imaging domains, and models incorporating prior shape knowledge are often essential for the automatic performance of this task. Higher-order active contours provide a way to include such knowledge, but their behaviour can vary significantly with parameter values: e.g. the same energy can model networks or a ‘gas of circles’. In this paper, we present a stability analysis of a HOAC energy leading to the phase diagram of a long bar. The results, which are confirmed by numerical experiments, enable the selection of parameter values for the modelling of network shapes using the energy. We apply the resulting model to the problem of hydrographic network extraction from VHR satellite images.

1. Introduction

Water is a critical resource, and likely to become more so in the future. Consequently, the ability to monitor water resources is of great importance. One way to perform such monitoring is through the use of remote sensing imagery, and in particular, very high resolution (VHR) satellite images. Many organizations that deal with such images are therefore interested in automatic processing that allows water resources to be measured and evaluated. In this paper, we focus on the segmentation of hydrographic networks from VHR satellite images. This is a difficult problem, because such networks are frequently not distinguishable from the background by local image measurements alone. Rather, it is the ‘shape’ of the region in the image domain occupied by the network that distinguishes them: networks occupy regions in the image domain composed of a set of ‘arms’ that join together at junctions. There has been a significant amount of work on shape modelling for segmentation, but most of this work models shape by comparing the region sought to a reference regions or regions [1, 7]. This works well for many applications, but is not appropriate when the region sought has an arbitrary topology, as is the case for networks. To model such regions, ‘higher-order active contours’ (HOACs) were introduced [5]. Unlike the original active contours [4] and their successors, this new generation of active contours incorporates non-trivial shape information about the region being modelled via long-range interactions between region boundary points.

The HOAC prior energy defined in [5] was used to model network-shaped regions. However, it was soon discovered that it could be used to model other families of regions too, notably a ‘gas of circles’ (regions composed of an arbitrary number of approximate circles [3]), simply by varying the model parameters. Figure 1 illustrates these different behaviours. Thus, in order to use the model for network segmentation, for example, it is important to know which parameter ranges lead to stable network regions, as opposed, for example, to stable ‘gas of circles’ regions.

This paper addresses this issue by performing a stability analysis of the prior HOAC energy introduced...
2 Theory

The prior HOAC energy introduced in [5] is
\[ E_G(R) = \lambda_C L(R) + \alpha_C A(R) - \frac{\beta_C}{2} \int dt \int dt' \gamma(t) \cdot \gamma(t') \Phi(\sigma(t, t')/d), \]
where \( R \) is a region in the image domain; \( L \) is region boundary length; \( A \) is region area; \( \gamma \) is an embedding representing the region boundary, parameterized by \( t \); \( \gamma \) is its derivative, i.e. the tangent vector field; \( \Phi \) defines the interaction between two boundary points separated by a distance \( \sigma(t, t') = |\gamma(t) - \gamma(t')| \); and \( \lambda_C, \alpha_C, \beta_C, \) and \( d \) are real parameters. In particular, \( d \) controls the interaction range. The long-range interaction is responsible for the prior shape information.

Our aim is to find the ranges of \( \alpha_C \) and \( \beta_C \) (without loss of generality we can put \( \lambda_C = d = 1 \)) for which a long bar is stable, i.e. is a local minimum of this energy.

To begin, we Taylor expand the energy around a long bar.

\[ E_G(R) \doteq E_G(\gamma_0) + \int dt \delta \gamma \tfrac{\partial E_G}{\partial \gamma}(\gamma_0) + \tfrac{1}{2} \int dt \int dt' \delta \gamma \cdot \tfrac{\partial^2 E_G}{\partial \gamma^2}(\gamma_0) \delta \gamma, \]
where \( \delta \gamma \) is a perturbation and the derivative \( \tfrac{\partial E_G}{\partial \gamma}(\gamma_0) \) and the second functional derivative should be positive definite (the extremum is a minimum).

The second-order Taylor expansion of \( E_G \) is
\[ E_G^{(2)}(\gamma) = E_G^{(2)}(\gamma_0 + \delta \gamma) \doteq E_G(\gamma_0) + \int dt \delta \gamma \tfrac{\partial E_G}{\partial \gamma}(\gamma_0) + \tfrac{1}{2} \int dt \int dt' \delta \gamma \cdot \tfrac{\partial^2 E_G}{\partial \gamma^2}(\gamma_0) \delta \gamma, \]
where \( \langle \cdot | \cdot \rangle \) is a metric defined over the space of embeddings, and \( \delta \gamma \) is a small perturbation of \( \gamma_0 \). The configuration \( \gamma_0 \) is stable if and only if \( \tfrac{\partial E_G}{\partial \gamma}(\gamma_0) = 0 \), i.e. if \( \gamma_0 \) is an extremum of \( E_G \), and the Hessian matrix \( \tfrac{\partial^2 E_G}{\partial \gamma^2}(\gamma_0) \) is positive definite, i.e. the extremum is a minimum.

A bar, \( \gamma_0 \), of length \( l \) and width \( w_0 \), is given by
\[ \gamma_{0,\mu}(t) = \begin{cases} x_{0,\mu}(t) = \pm \mu t, & t \in [-\frac{1}{2}, \frac{1}{2}] \\ y_{0,\mu}(t) = \pm \frac{w_0}{2} \end{cases} \]
The bar can be viewed as periodic in \( x \) with period \( l \), and has two boundary components, indicated by \( \mu = 1 \) or 2. The symbol \( \pm \mu = 1 \) if \( \mu = 1 \) and \( -1 \) if \( \mu = 2 \).

Perturbations of the bar \( \delta \gamma \) are defined by tangential and normal changes in the boundary components:
\[ \delta \gamma_{\mu}(t) = (\delta x_{\mu}(t), \delta y_{\mu}(t)). \]
Tangential changes do not change the bar’s shape and hence do not change the energy; we can therefore set \( \delta x_{\mu} = 0 \). The Hessian matrix is diagonal in the Fourier basis of the tangent space at \( \gamma_0 \), so it is useful to express \( \delta \gamma \) in this basis:
\[ \delta y_{\mu}(t) = \sum_{k_{\mu}} a_{k_{\mu}} e^{ik_{\mu} t} \]
After some calculation, we find [2]
\[ e_G^{(2)}(\gamma) \doteq \underbrace{E_G^{(2)}(\gamma)}_{e_0} \doteq \underbrace{2 + \alpha_C w_0 - \beta_C G_{00}(w_0)}_{e_1} \]
\[ + \underbrace{[a_{1,0} - a_{2,0}] \alpha_C - \beta_C G_{10}(w_0)}_{e_2} \]
\[ + \underbrace{\frac{1}{2} \sum_k [a_{1,k}^2 + a_{2,k}^2]}_{e_3} \underbrace{\left[ k^2 + \beta_C G_{20}(w_0, k) \right]}_{e_4} \]
\[ + \underbrace{a_{1,k} a_{2,k} + a_{1,-k} a_{2,-k}}_{e_5} \beta_C G_{21}(w_0, k), \]
where \( a_k = (a_{1,k}^*, a_{2,k}); \) and \( \gamma \) and \( * \) indicate respectively Hermitian and complex conjugates; and the \( G(w_0, k) \) depend on \( \Phi \), \( e_2(w_0) \) is the energy per unit length of a long bar of width \( w_0 \), while \( e_1(w_0) = \partial e_0(w_0)/\partial w_0 \) is the change in energy due to a change in width (to first order, non-zero Fourier frequencies do not contribute).

A preliminary version of this work was published in French [2].
matrix whose diagonal and off-diagonal terms, \( e_{20} \) and \( e_{21} \), express the self-energy of perturbations of one side, and the interaction between perturbations of opposite sides of the bar respectively.

### 2.2 Stability analysis of a long bar

The condition that \( \delta E_G / \delta \gamma = 0 \), i.e. that \( e_1 = 0 \), determines \( \beta_C \) in terms of \( \alpha_C \) and \( w_0 \):

\[
\beta_C(\alpha_C, w_0) = \frac{\alpha_C}{G_{10}(w_0)}. \tag{5}
\]

The condition that the Hessian be positive definite, i.e. that the eigenvalues of \( e_2 \) be strictly positive for all frequencies, allows us to find lower and upper bounds on \( \alpha_C \) [2]. The eigenvectors of \( e_2 \) are \( v_{\pm} = (1, \pm 1) \), with eigenvalues \( \lambda_{\pm,k} = e_{20}(k) \pm e_{21}(k) \) respectively. Equation (4) can be written in terms of \( \lambda_{\pm} \):

\[
e^{(2)}_G = e_0 + e_1(a_{1,0} - a_{2,0}) + \frac{1}{4} k \sum \lambda_{+,k} |a^+_k|^2 + \lambda_{-,k} |a^-_k|^2
\]

where \( a^\pm_k = a_{1,k} \pm a_{2,k} \). These eigenmodes are illustrated in Figure 2: \( a^+ \) describes in-phase perturbations of the two sides, while \( a^- \) describes out-of-phase perturbations. For low frequencies, the former cost less energy, as locally the sides of the shape are still parallel. For higher frequencies, the difference is negligible.

A minimum in the energy exists only for \( w_0 \in (0.88, 2) \) due to the form of \( \Phi \) [2]. Using equation (5), we can plot \( \beta_C \) against \( \alpha_C \) for \( w_0 \in (0.88, 2) \), between the upper and lower bounds on \( \alpha_C \). The result is a 2D zone in the \((\alpha, \beta)\) plane, illustrated in Figure 3, in which the parameter values permit a stable bar. Since the total energy is \( E_G^{(2)}(\gamma) = l e_G^{(2)}(\gamma) \), the behaviour of a long bar depends on the sign of \( e_G^{(2)}(\gamma) \). If \( e_G^{(2)} > 0 \), the bar will shorten and either form a stable circle or disappear. If \( e_G^{(2)} < 0 \), the bar will elongate, without limit if the domain is unbounded. The blue and yellow zones in Figure 3 correspond to \( e_G^{(2)} > 0 \) and \( e_G^{(2)} < 0 \) respectively. Below the coloured zone, the dominant instability is change in \( w_0 \): in gradient descent, the bar should get thinner and disappear. Above the coloured zone, the dominant instability is at non-zero frequencies: the bar should develop oscillations and break up. These results will be confirmed numerically in Section 3.1.

### 3 Experiment

We first present gradient descent evolution experiments using \( E_G \) (not \( E_G^{(2)} \)) to confirm the results of the analysis in Section 2. We then add a likelihood term, and test the energy on a real image.

#### 3.1 Geometric evolution

Figure 4 shows five gradient descent evolutions starting from a long bar. The parameter values for the rows of Figure 4 were \((w_0, \alpha_C, \beta_C, d) = (0, 0.052, 0.02, 10), (10, 0.067, 0.053, 7.75), (30, 8.48, 6.28, 23.8), (20, 8.02, 6.56, 15.26), \) and \((25, 0.319, 2.4, 11.36) \). The following points correspond to these rows: 1) parameters chosen from below
the coloured zone; the bar is unstable to changes in its width, and so it thins and disappears; 2) parameters chosen from the blue zone; the bar is stable, but has positive energy per unit length, so it shortens and disappears; 3) parameters again chosen from the blue zone; the bar shortens, but rather than disappear, it evolves to a stable circle [3]; 4) parameters chosen from the yellow zone; the bar has negative energy per unit length, so it elongates, and then buckles due to the finite size of the domain; 5) parameters chosen from above the coloured zone; the bar is unstable to non-zero frequencies, so it develops oscillations and then breaks up; note that the oscillations are in-phase, since these cost less energy and are therefore the first to become unstable; their frequency corresponds quantitatively with the prediction of the stability analysis.

3.2 Results on images

In this Section, we show the results of river network extraction from a piece of a four-band colour infrared Quickbird image (60cm resolution) using $E_G$. We add to $E_G$ a likelihood term $E_I(I, R)$ consisting of multivariate Gaussian models of the image in the region interior and exterior, with parameters learned by example. For implementation, we use the implicit ‘phase field’ framework [6]. This simplifies and speeds up the algorithm. We select HOAC parameter values from the phase diagram and then compute the corresponding phase field parameters [6]. Figure 5 shows a segmentation result. The parameter values were $(w_0, d, \lambda_C, \alpha_C, \beta_C) = (5, 5.67, 80, 16.91, 8.77)$.

4 Conclusion

We have calculated the phase diagram of a bar under a higher-order active contour energy via a stability analysis, and we have confirmed the relevance of the analysis via experiments. Aside from its intrinsic interest, the calculation enables the determination of model parameter values for hydrographic network extraction from VHR remote sensing images. This leads to the successful extraction of hydrographic networks, although further work remains to be done on closing gaps.

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