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# A regularity result for a solid-fluid system associated to the compressible Navier-Stokes equations

M. Boulakia\* and S. Guerrero\*

## Abstract

In this paper we deal with a fluid-structure interaction problem for a compressible fluid and a rigid structure immersed in a regular bounded domain in dimension 3. The fluid is modelled by the compressible Navier-Stokes system in the barotropic regime with no-slip boundary conditions and the motion of the structure is described by the usual law of balance of linear and angular moment.

The main result of the paper states that, for small initial data, we have the existence and uniqueness of global smooth solutions as long as no collisions occur. This result is proved in two steps; first, we prove the existence and uniqueness of local solution and then we establish some a priori estimates independently of time.

## Résumé

Dans cet article, nous considérons un problème d'interaction fluide-structure entre un fluide compressible et une structure rigide évoluant à l'intérieur d'un domaine borné et régulier en dimension 3. Le fluide est décrit par le système de Navier-Stokes compressible barotrope avec des conditions de non-glissement sur le bord et le mouvement de la structure est régi par les lois de conservation des moments linéaire et angulaire.

Nous montrons, pour des données initiales petites, l'existence et l'unicité de solutions globales régulières tant qu'il n'y a pas de chocs. Ce résultat est obtenu en deux temps; tout d'abord, nous prouvons l'existence et l'unicité de solutions locales puis nous démontrons des estimations a priori indépendamment du temps.

## 1 Introduction

### 1.1 Statement of problem

We consider a rigid structure immersed in a viscous compressible fluid. At time  $t$ , we denote by  $\Omega_S(t)$  the domain occupied by the structure. The structure and the fluid are contained in a fixed bounded domain  $\Omega \subset \mathbb{R}^3$ . We suppose that the boundaries of  $\Omega_S(0)$  and  $\Omega$  are smooth ( $C^4$  for instance) and that

$$\Omega_S(0) \text{ is convex, } d(\Omega_S(0), \partial\Omega) > 0. \quad (1)$$

For any  $t > 0$ , we denote by  $\Omega_F(t) = \Omega \setminus \overline{\Omega_S(t)}$  the region occupied by the fluid. The time evolution of the eulerian velocity  $u$  and the density  $\rho$  in the fluid are governed by the compressible Navier-Stokes equations and the continuity equation:  $\forall t > 0, \forall x \in \Omega_F(t)$

$$\begin{cases} (\rho_t + \nabla \cdot (\rho u))(t, x) = 0, \\ (\rho u_t + \rho(u \cdot \nabla)u)(t, x) - \nabla \cdot (2\mu\epsilon(u) + \mu'(\nabla \cdot u)Id)(t, x) + \nabla p(t, x) = 0, \end{cases} \quad (2)$$

where  $\epsilon(u) = \frac{1}{2}(\nabla u + \nabla u^t)$  denotes the symmetric part of the gradient. The viscosity coefficients  $\mu$  and  $\mu'$  are real constants which are supposed to satisfy

$$\mu > 0, \mu + \mu' \geq 0. \quad (3)$$

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We suppose that we are in a barotropic regime where a constitutive law gives the relation between the pressure of the fluid  $p$  and the density  $\rho$ . Thus, we suppose that

$$p = P(\rho) \text{ where } P \in C^\infty(\mathbb{R}_+^*), P(\rho) > 0 \text{ and } P'(\rho) > 0, \forall \rho > 0. \quad (4)$$

For instance,  $P(\rho) := \rho^\gamma$  with  $\gamma > 0$  is admissible.

Concerning the compressible fluids, a local in time result of existence and uniqueness of a smooth solution was proved in [21]. In [18], the authors proved the existence and uniqueness of a regular solution for small initial data and external forces.

Next, for isentropic fluids ( $P(\rho) = \rho^\gamma$ ,  $\gamma > 0$ ), the global existence of a weak solution for small initial data was proved in [13] (for  $\gamma = 1$ ) and in [14] (for  $\gamma > 1$ ). Also for an isentropic fluid, the first global result for large data was proved in [15] (see also [16]) (with  $\gamma \geq 9/5$  for dimension  $N = 3$  and with  $\gamma > N/2$  for  $N \geq 4$ ). Finally, this last result was improved in [9] (see also [11]) (with  $\gamma > N/2$  for  $N \geq 3$ ).

At time  $t$ , the motion of the rigid structure is given by the position  $a(t) \in \mathbb{R}^3$  of the center of mass and by a rotation (orthogonal) matrix  $Q(t) \in \mathbb{M}_{3 \times 3}(\mathbb{R})$ . Without loss of generality, we can suppose that

$$a(0) = 0 \text{ and } Q(0) = Id. \quad (5)$$

At time  $t$ , the domain occupied by the structure  $\Omega_S(t)$  is defined by

$$\Omega_S(t) = \chi_S(t, \Omega_S(0)), \quad (6)$$

where  $\chi_S$  denotes the flow associated to the motion of the structure:

$$\chi_S(t, y) = a(t) + Q(t)y, \forall y \in \Omega_S(0), \forall t > 0. \quad (7)$$

We notice that, for each  $t > 0$ ,  $\chi_S(t, \cdot) : \Omega_S(0) \rightarrow \Omega_S(t)$  is invertible and

$$\chi_S(t, \cdot)^{-1}(x) = Q(t)^{-1}(x - a(t)), \forall x \in \Omega_S(t).$$

Thus, the eulerian velocity of the structure is given by

$$(\chi_S)_t(t, \cdot) \circ \chi_S(t, \cdot)^{-1}(x) = \dot{a}(t) + \dot{Q}(t)Q(t)^{-1}(x - a(t)), \forall x \in \Omega_S(t).$$

Since  $\dot{Q}(t)Q(t)^{-1}$  is skew-symmetric, for each  $t > 0$ , we can represent this matrix by a unique vector  $\omega(t) \in \mathbb{R}^3$  such that

$$\dot{Q}(t)Q(t)^{-1}y = \omega(t) \wedge y, \forall y \in \mathbb{R}^3.$$

Reciprocally, if  $\omega$  belongs to  $L^2(0, T)$ , then there exists a unique matrix  $Q \in H^1(0, T)$  such that  $Q(0) = Id$  and which satisfies this formula.

Thus, the eulerian velocity  $u_S$  of the structure is given by

$$u_S(t, x) = \dot{a}(t) + \omega(t) \wedge (x - a(t)), \forall x \in \Omega_S(t). \quad (8)$$

For the equations of the structure, we denote by  $m > 0$  the mass of the rigid structure and  $J(t) \in \mathbb{M}_{3 \times 3}(\mathbb{R})$  its tensor of inertia at time  $t$ . This tensor is given by

$$J(t)b \cdot \tilde{b} = \int_{\Omega_S(0)} \rho_{0,S}(y) (b \wedge Q(t)y) \cdot (\tilde{b} \wedge Q(t)y) dy \quad \forall b, \tilde{b} \in \mathbb{R}^3, \quad (9)$$

where  $\rho_{0,S} > 0$  is the initial density of the structure. One can prove that

$$J(t)b \cdot b \geq C_J |b|^2 > 0 \text{ for all } b \in \mathbb{R}^3 \setminus \{0\}, \quad (10)$$

where  $C_J$  is independent of  $t > 0$ . The equations of the structure motion are given by the balance of linear and angular momentum. We have, for all  $t \in (0, T)$

$$\begin{cases} m\ddot{a} = \int_{\partial\Omega_S(t)} (2\mu\epsilon(u) + \mu'(\nabla \cdot u)Id - pId)n d\gamma, \\ J\dot{\omega} = (J\omega) \wedge \omega + \int_{\partial\Omega_S(t)} (x - a) \wedge ((2\mu\epsilon(u) + \mu'(\nabla \cdot u)Id - pId)n) d\gamma. \end{cases} \quad (11)$$

In these equations,  $n$  is the outward unit normal to  $\partial\Omega_S(t)$ . On the boundary of the fluid, the eulerian velocity has to satisfy a no-slip boundary condition. Therefore, we have, for all  $t > 0$

$$\begin{cases} u(t, x) = 0, \forall x \in \partial\Omega, \\ u(t, x) = \dot{a}(t) + \omega(t) \wedge (x - a(t)), \forall x \in \partial\Omega_S(t). \end{cases} \quad (12)$$

The system is completed by the following initial conditions:

$$u(0, \cdot) = u_0 \text{ in } \Omega_F(0), \rho(0, \cdot) = \rho_0 \text{ in } \Omega_F(0), a(0) = 0, \dot{a}(0) = a_0, \omega(0) = \omega_0, \quad (13)$$

which satisfy

$$a_0, \omega_0 \in \mathbb{R}^3, \rho_0, u_0 \in H^3(\Omega_F(0)), \rho_0(x) > 0, \forall x \in \Omega_F(0). \quad (14)$$

Since we will deal with smooth solutions, we will also need some compatibility conditions to be satisfied:

$$u_0 = a_0 + \omega_0 \wedge x \text{ on } \partial\Omega_S(0), \quad u_0 = 0 \text{ on } \partial\Omega \quad (15)$$

and

$$\begin{aligned} & \frac{1}{\rho_0} \nabla \cdot (2\mu\epsilon(u_0) + \mu'(\nabla \cdot u_0)Id) - \frac{1}{\rho_0} \nabla P(\rho_0) = \frac{1}{m} \int_{\partial\Omega_S(0)} (2\mu\epsilon(u_0) + \mu'(\nabla \cdot u_0)Id - P(\rho_0)Id)n \, d\gamma \\ & + (J(0)^{-1}(J(0)\omega_0) \wedge \omega_0) \wedge x + J(0)^{-1} \left( \int_{\partial\Omega_S(0)} x \wedge ((2\mu\epsilon(u_0) + \mu'(\nabla \cdot u_0)Id - P(\rho_0)Id)n) \, d\gamma \right) \wedge x \quad (16) \end{aligned}$$

$$+\omega_0 \wedge (\omega_0 \wedge x) \text{ on } \partial\Omega_S(0),$$

$$\nabla \cdot (2\mu\epsilon(u_0) + \mu'(\nabla \cdot u_0)Id) - \nabla P(\rho_0) = 0 \text{ on } \partial\Omega$$

These two conditions are formally obtained by differentiating system (12) with respect to time and taking  $t = 0$ . To do this, we consider the second equation of (12) on a fix domain by setting  $x = \chi_S(t, y)$ , with  $y \in \partial\Omega_S(0)$ .

Let us now recall some of the most relevant results in problems of fluid-structure interaction. In the below lines, when we refer to a global result, we mean before collision (in the case of a rigid solid) or before interpenetration of the structure (in the case of an elastic solid).

- Incompressible fluids:

As long as rigid solids are concerned, a local result was proved in [12], while the existence of global weak solutions is proved in [5] and [7] (with variable density) and [19] ( $2D$ , with variable density); in this last paper, the existence of a solution is proved even beyond collisions. Later, the existence and uniqueness of strong global solutions in  $2D$  was proved in [20] as well as the local in time existence and uniqueness of strong solutions in  $3D$ .

When talking about elastic solids, a first existence result of weak solution was proved in [8], when the elastic deformation is given by a finite sum of modes. A local existence result of a strong solution for an elastic plate was proved in [1] ( $2D$ ). The local existence of a strong solution is proved in [6]. In [4] and [3] (with variable density), the authors proved the global existence of a weak solution.

- Compressible fluids:

Concerning rigid solids, the global existence of a weak solution was proved in [7] for  $\gamma \geq 2$  and in [10] for  $\gamma > N/2$ . For elastic solids, in [2] the author proved the global existence of a weak solution in  $3D$  for  $\gamma > 3/2$ .

In this paper, we will prove the existence and uniqueness of smooth global solutions for small initial data (Theorem 3). We can also prove the same result for initial data close to a stationary solution  $(\rho, u, a, \omega) = (\rho_e, 0, 0, 0)$  and for special right hand sides (see Remark 6 for more details).

We give a lemma which allows to extend the flow  $\chi_S$  by a flow  $\chi$  defined on the global domain  $\Omega$ .

**Lemma 1** Let  $T \in (0, +\infty)$  and  $(a, \omega) \in (H^3(0, T) \cap W^{2, \infty}(0, T)) \times (H^2(0, T) \cap W^{1, \infty}(0, T))$  be given. We suppose that  $(a, Q)$  satisfies (5) where  $Q \in H^3(0, T)$  is the rotation matrix associated to  $\omega$ . We consider the associated flow  $\chi_S$ , the eulerian velocity  $u_S$  and the domain defined by (6) to (8). We suppose that there exists  $\alpha > 0$  such that

$$\forall t \in [0, T], d(\Omega_S(t), \partial\Omega) \geq \alpha > 0. \quad (17)$$

Then, we can extend the flow  $\chi_S$  by a flow  $\chi \in H^3(0, T; C^\infty(\overline{\Omega}))$  such that

- $\chi(t, y) = a(t) + Q(t)y$  for every  $y$  such that  $\text{dist}(y, \partial\Omega_S(0)) < \alpha/4$  and every  $t \in (0, T)$ .
- $\chi(t, y) = y$  for every  $y$  such that  $\text{dist}(y, \partial\Omega) < \alpha/4$  and for every  $t \in (0, T)$ .
- For all  $t \in (0, T)$ ,  $\chi(t, \cdot)$  is invertible from  $\Omega$  onto  $\Omega$  and from  $\Omega_F(0)$  onto  $\Omega_F(t)$  and

$$(t, x) \in (0, T) \times \Omega \rightarrow \chi(t, \cdot)^{-1}(x)$$

belongs to  $H^3(0, T; C^\infty(\overline{\Omega}))$ .

- The eulerian velocity  $v \in H^2(0, T; C^\infty(\Omega))$  associated to  $\chi$  satisfies

$$\begin{cases} v(t, x) = u_S(t, x) & \forall x \in \Omega_S(t), \forall t \in (0, T), \\ v(t, x) = 0 & \forall x \in \partial\Omega, \forall t \in (0, T). \end{cases}$$

- For  $p \in \mathbb{N}$  there exists a constant  $C_1 > 0$  just depending on  $p$  and  $\Omega$  such that

$$\|\chi_t(t, x)\|_{H^p(\Omega)} \leq C_1(|\dot{a}(t)|^2 + |\omega(t)|^2)^{1/2}, \quad \forall t \in (0, T). \quad (18)$$

**Proof:** We consider a cut-off function which satisfies

$$\xi \in C^\infty(\overline{\Omega}), \quad \xi(y) = 1 \quad \text{dist}(y, \partial\Omega) > \alpha/2, \quad \xi(y) = 0 \quad \text{dist}(y, \partial\Omega) < \alpha/4. \quad (19)$$

Let define the eulerian velocity in the whole  $(0, T) \times \Omega$  by

$$v(t, y) := \xi(y)(\dot{a}(t) + \omega(t) \wedge (y - a(t))) \quad t \in (0, T), \quad y \in \Omega. \quad (20)$$

Next, we can define the flow  $\chi$  associated to  $v$ . For all  $y \in \Omega$ ,  $\chi(\cdot, y)$  is defined on  $(0, T)$  as the solution of the ordinary differential equation

$$\begin{cases} \chi_t(t, y) &= v(t, \chi(t, y)), \\ \chi(0, y) &= y. \end{cases} \quad (21)$$

In  $\{y : \text{dist}(y, \partial\Omega) < \alpha/4\}$ , since  $v = 0$ , we have according to the uniqueness of the flow that  $\chi(t, y) = y$ . Analogously, if we take a point in the set  $\{y : \text{dist}(y, \partial\Omega_S(0)) < \alpha/4\}$ ,  $\chi(t, y) = \chi_S(t, y) = a(t) + Q(t)y$ . The last point comes from (19), (20) and (21).

We define now a flow which transforms the moving domains into the initial one:

**Corollary 2** Under the same hypotheses as in the previous lemma, one can define an inverse flow  $\psi$  extending

$$\chi_S^{-1}(t, x) = Q(t)^{-1}(x - a(t)), \quad t \in (0, T), \quad x \in \Omega_S(t)$$

to the whole domain  $\Omega$  and which satisfies:

- $\psi(t, x) = Q(t)^{-1}(x - a(t)), \forall x \in \Omega_S(t), \forall t \in (0, T)$ .
- $\psi(t, x) = x, \forall x \in \partial\Omega, \forall t \in (0, T)$ .
- The velocity  $w := \psi_t$  associated to  $\psi$  belongs to  $H^2(0, T; C^\infty(\Omega))$  and satisfies

$$\begin{cases} w(t, x) = -Q(t)^{-1}(\dot{a}(t) + \omega \wedge (x - a(t))) & \forall x \in \Omega_S(t), \forall t \in (0, T), \\ w(t, x) = 0 & \forall x \in \partial\Omega, \forall t \in (0, T). \end{cases}$$

- For  $p \in \mathbb{N}$  there exists a constant  $C_2 > 0$  just depending on  $p$  and  $\Omega$  such that

$$\|\psi_t(t, x)\|_{H^p(\Omega)} \leq C_2(|\dot{a}(t)|^2 + |\omega(t)|^2)^{1/2}, \quad \forall t \in (0, T). \quad (22)$$

## 1.2 Main result

Let us now introduce some notation which we will employ all along the paper. Let  $h \in (0, +\infty]$  and  $(a, \omega) \in (H^3(0, h) \cap W^{2, \infty}(0, h)) \times (H^2(0, h) \cap W^{1, \infty}(0, h))$  be given. First we define

$$Z_h := \{(s, x) : s \in (0, h), x \in \Omega_F(s)\}, \quad \Sigma_h := \{(s, x) : s \in (0, h), x \in \partial\Omega_S(s)\}. \quad (23)$$

Then, for  $r, p \geq 0$  natural numbers, we introduce

$$\begin{aligned} L_h^2(L^2) &:= L^2(Z_h) = \left\{ u \text{ measurable} : \int_0^h \int_{\Omega_F(s)} |u|^2 dx ds < +\infty \right\}, \\ L_h^2(H^p) &:= \left\{ u \in L_h^2(L^2) : \int_0^h \|u\|_{H^p(\Omega_F(s))}^2 ds < +\infty \right\}, \\ H_h^r(H^p) &:= \left\{ u \in L_h^2(L^2) : \int_0^h \sum_{\beta=0}^r \|\partial_t^\beta u\|_{H^p(\Omega_F(s))}^2 ds < +\infty \right\}, \end{aligned}$$

with associated norms given by the definition. On the other hand, we define

$$C_h^0(L^2) := \{u \text{ such that } \hat{u}(s, x) := u(s, x)1_{\Omega_F(s)} \in C^0([0, h]; L^2(\Omega))\}$$

and

$$C_h^r(H^p) := \{u : \partial_t^\beta \partial_x^\alpha u \in C_h^0(L^2), \forall 0 \leq \beta \leq r, \forall 0 \leq |\alpha| \leq p\}$$

with associated norms given by

$$\|u\|_{C_h^0(L^2)} := \max_{t \in (0, h)} \|u(t)\|_{L^2(\Omega_F(t))} = \max_{t \in (0, h)} \|\hat{u}(t)\|_{L^2(\Omega)}$$

and

$$\|u\|_{C_h^r(H^p)} := \sum_{\beta=0}^r \max_{t \in (0, h)} \|\partial_t^\beta u(t)\|_{H^p(\Omega_F(t))}.$$

It is clear that for regular functions, the norm  $\|\cdot\|_{C_h^r(H^p)}$  coincides with the norm  $\|\cdot\|_{W_h^{r, \infty}(H^p)}$ . In the sequel, we will always use the second notation.

The main goal of our paper is to prove the following theorem:

**Theorem 3** *Let  $\bar{\rho}$  be the mean-value of  $\rho_0$  in  $\Omega_F(0)$ . We suppose that (1), (3) and (4) are satisfied. We suppose that the initial conditions satisfy (5), (14) and the compatibility conditions (15)-(16). Then there exists a constant  $\delta > 0$  such that, if*

$$\|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + \|u_0\|_{H^3(\Omega_F(0))} + |a_0| + |\omega_0| < \delta, \quad (24)$$

*the system of equations (2), (11), (12) and (13) admits a unique solution  $(\rho, u, a, \omega)$  defined on  $(0, T)$  for all  $T$  such that (17) is satisfied. Moreover, this solution belongs to the following space*

$$\begin{aligned} \rho &\in L_T^2(H^3) \cap C_T^0(H^3) \cap H_T^1(H^2) \cap C_T^1(H^2) \cap H_T^2(L^2), \\ u &\in L_T^2(H^4) \cap C_T^0(H^3) \cap C_T^1(H^1) \cap H_T^2(L^2), \quad a \in H^3(0, T), \quad \omega \in H^2(0, T) \end{aligned}$$

*and there exists a positive constant  $C_1$  independent of  $T$  such that*

$$\begin{aligned} &\|\rho - \bar{\rho}\|_{L_T^2(H^3)} + \|\rho - \bar{\rho}\|_{L_T^\infty(H^3)} + \|\rho - \bar{\rho}\|_{W_T^{1, \infty}(H^2)} + \|\rho - \bar{\rho}\|_{H_T^2(L^2)} + \|\rho - \bar{\rho}\|_{H_T^1(H^2)} \\ &+ \|u\|_{L_T^2(H^4)} + \|u\|_{L_T^\infty(H^3)} + \|u\|_{W_T^{1, \infty}(H^1)} + \|u\|_{H_T^2(L^2)} + \|\dot{a}\|_{W_T^{1, \infty}} + \|\dot{a}\|_{H_T^2} + \|\omega\|_{W_T^{1, \infty}} + \|\omega\|_{H_T^2} \\ &\leq C_1(\|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + \|u_0\|_{H^3(\Omega_F(0))} + |a_0| + |\omega_0|). \end{aligned} \quad (25)$$

**Remark 4** In Theorem 3, we prove the existence and uniqueness of a regular solution of the system of equations (2), (11), (12) and (13) provided that the structure does not touch  $\partial\Omega$  (see (17)). Observe that this condition will always be satisfied on an interval  $(0, T_{\min})$ , where  $T_{\min} > 0$  only depends on  $C_1$ ,  $\delta$  and  $d(\partial\Omega, \bar{\Omega}_S(0))$ , as easily seen from estimate (25).

**Remark 5** This condition on  $T$  is natural for fluid-structure interaction problems (see, for instance, [7], [5] and [20]). For results concerning the existence of weak solutions of this type of systems after a collision has occurred, we refer to [19] and [10].

**Remark 6** As proved in [18], there exists a unique stationary solution  $(\rho_e(x), 0, 0, 0)$  of (2), (12), where, on the fluid, the compressible Navier-Stokes equation is completed with a right hand side of the form  $f_i(x) = \partial_i \phi(x)$  ( $i = 1, 2, 3$ ) satisfying (24) for the  $H^3$ -norm. Then, we can prove that Theorem 3 also holds with  $\bar{\rho}$  replaced by  $\rho_e$  and for a right hand side of the above kind.

## 2 Rewriting of the problem and intermediate results

Let us define  $p^0$  by

$$p^0 := \frac{P'(\bar{\rho})}{\bar{\rho}},$$

and recall that  $\bar{\rho}$  is the mean-value of  $\rho_0$ :

$$\bar{\rho} = \frac{1}{V(\Omega_F(0))} \int_{\Omega_F(0)} \rho_0,$$

where  $V(\Omega_F(0))$  stands for the volume of  $\Omega_F(0)$ . Furthermore, we change  $m/\bar{\rho}$  into  $m$ ,  $J/\bar{\rho}$  into  $J$ ,  $\mu/\bar{\rho}$  into  $\mu$  and  $\mu'/\bar{\rho}$  into  $\mu'$ . Then, defining  $\rho^*(t, x) := \rho(t, x) - \bar{\rho}$ , we find that the system of equations (2), (11), (12) and (13) can be written as follows:

$$\left\{ \begin{array}{ll} \rho_t^* + u \cdot \nabla \rho^* + \bar{\rho} \nabla \cdot u = f_0(\rho^*, u, a, \omega) & \text{in } (0, T) \times \Omega_F(t), \\ u_t - \nabla \cdot \sigma(u, \rho^*) = f_1(\rho^*, u, a, \omega) & \text{in } (0, T) \times \Omega_F(t), \\ m\ddot{a} = \int_{\partial\Omega_S(t)} \sigma(u, \rho^*) n \, d\gamma + f_2(\rho^*, u, a, \omega) & \text{in } (0, T), \\ J\dot{\omega} = \int_{\partial\Omega_S(t)} (x - a) \wedge (\sigma(u, \rho^*) n) \, d\gamma + f_3(\rho^*, u, a, \omega) & \text{in } (0, T), \\ u = 0 & \text{on } (0, T) \times \partial\Omega, \\ u = \dot{a} + \omega \wedge (x - a) & \text{on } (0, T) \times \partial\Omega_S(t), \\ \rho^*(0, \cdot) = \rho_0 - \bar{\rho}, \quad u(0, \cdot) = u_0 & \text{in } \Omega_F(0), \\ a(0) = 0, \quad \dot{a}(0) = a_0, \quad \omega(0) = \omega_0, & \end{array} \right. \quad (26)$$

where

$$\sigma(u, \rho) := 2\mu\epsilon(u) + \mu'(\nabla \cdot u)Id - p^0 \rho Id \quad (27)$$

and

$$\left\{ \begin{array}{l} f_0(\rho^*, u, a, \omega) = -\rho^*(\nabla \cdot u), \\ f_1(\rho^*, u, a, \omega) = -(u \cdot \nabla)u - \frac{\rho^*}{\rho^* + \bar{\rho}} \nabla \cdot (2\mu\epsilon(u) + \mu'(\nabla \cdot u)Id) + \left( p^0 - \frac{P'(\rho^* + \bar{\rho})}{\rho^* + \bar{\rho}} \right) \nabla \rho^*, \\ f_2(\rho^*, u, a, \omega) = \int_{\partial\Omega_S(t)} \left( p^0 \rho^* - \frac{P(\rho^* + \bar{\rho})}{\bar{\rho}} \right) n \, d\gamma, \\ f_3(\rho^*, u, a, \omega) = (J\omega) \wedge \omega + \int_{\partial\Omega_S(t)} (x - a) \wedge \left( \left( p^0 \rho^* - \frac{P(\rho^* + \bar{\rho})}{\bar{\rho}} \right) n \right) \, d\gamma. \end{array} \right. \quad (28)$$

For the integrals in the expressions of  $f_2$  and  $f_3$ , we use that

$$\frac{P(\bar{\rho})}{\bar{\rho}} \int_{\partial\Omega_S(0)} (x-a) \wedge n \, d\gamma = 0 \quad \text{and} \quad \frac{P(\bar{\rho})}{\bar{\rho}} \int_{\partial\Omega_S(0)} n \, d\gamma = 0 \quad (29)$$

and so  $f_2$  and  $f_3 - (J\omega) \wedge \omega$  are quadratic terms of  $\rho^*$ .

Now, for  $0 \leq h \leq +\infty$ , we define the space

$$\begin{aligned} X(0, h) = & \{(\rho, u, a, \omega) : (\rho, u) \in C_h^0(H^3) \cap L_h^2(H^3) \cap C_h^1(H^2) \cap H_h^1(H^2) \cap H_h^2(L^2)\} \\ & \times (L_h^2(H^4) \cap C_h^0(H^3) \cap H_h^1(H^2) \cap C_h^1(H^1) \cap H_h^2(L^2)), (\dot{a}, \omega) \in (C_h^1 \cap H_h^2)\} \end{aligned}$$

endowed with the following norm:

$$\begin{aligned} N_{0,h}(\rho, u, a, \omega) = & \left( \|\rho\|_{L_h^\infty(H^3)}^2 + \|\rho\|_{L_h^2(H^3)}^2 + \|\rho\|_{W_h^{1,\infty}(H^2)}^2 + \|\rho\|_{H_h^1(H^2)}^2 + \|\rho\|_{H_h^2(L^2)}^2 + \|u\|_{L_h^2(H^4)}^2 \right. \\ & \left. + \|u\|_{L_h^\infty(H^3)}^2 + \|u\|_{H_h^1(H^2)}^2 + \|u\|_{W_h^{1,\infty}(H^1)}^2 + \|u\|_{H_h^2(L^2)}^2 + \|\dot{a}\|_{W_h^{1,\infty}}^2 + \|\dot{a}\|_{H_h^2}^2 + \|\omega\|_{W_h^{1,\infty}}^2 + \|\omega\|_{H_h^2}^2 \right)^{1/2}. \end{aligned} \quad (30)$$

The proof of Theorem 3 is divided in two steps: first, a local existence result and next a priori estimates for the system. A suitable combination of both will yield the desired global result.

Thus, we first formulate a local existence and uniqueness result for small initial data.

**Proposition 7** *We suppose that there exists  $h \geq 0$  and  $E_0 > 0$  such that (2),(11), (12) and (13) admits a unique solution  $(\rho, u, a, \omega)$  in  $X(0, h)$  satisfying  $N_{0,h}(\rho - \bar{\rho}, u, a, \omega) \leq E_0$  and  $d(\partial\Omega_S(h), \partial\Omega) > 0$ .*

*Then there exist constants  $0 < \varepsilon_0 \leq E_0$ ,  $\tau > 0$  and  $C_2 > 0$  independent of  $h$  such that, if  $N_{h,h}(\rho - \bar{\rho}, u, a, \omega) \leq \varepsilon_0$ , problem (26) has a unique solution  $(\rho, u, a, \omega)$  defined on  $(h, h + \tau)$  satisfying*

$$(\rho, u, a, \omega) \in X(h, h + \tau), \quad N_{h,h+\tau}(\rho - \bar{\rho}, u, a, \omega) \leq C_2 N_{h,h}(\rho - \bar{\rho}, u, a, \omega).$$

We present now the a priori estimates:

**Proposition 8** *We suppose that there exists  $T > 0$  such that the problem (2),(11), (12) and (13) admits a solution  $(\rho, u, a, \omega)$  in  $X(0, T)$ . Then there exist constants  $0 < \delta_1 \leq \delta_0$  and  $C_1 > 0$  independent of  $t$  such that if  $N_{0,T}(\rho - \bar{\rho}, u, a, \omega) \leq \delta_1$  then*

$$N_{0,T}(\rho - \bar{\rho}, u, a, \omega) \leq C_1 N_{0,0}(\rho - \bar{\rho}, u, a, \omega).$$

These two propositions will be proved in the next sections. They allow to prove theorem 3 in the following way:

**Proof of theorem 3:** We will apply iteratively propositions 7 and 8. We suppose that

$$N_{0,0}(\rho - \bar{\rho}, u, a, \omega) \leq \min \left( \delta_0, \frac{\delta_1}{C_0}, \frac{\delta_1}{C_1 \sqrt{1 + C_0^2}} \right).$$

According to proposition 7 for  $h = 0$ , one can define on  $(0, \tau)$  a solution  $(\rho^*, u, a, \omega) \in X(0, \tau)$  such that

$$N_{0,\tau}(\rho - \bar{\rho}, u, a, \omega) \leq C_0 N_{0,0}(\rho - \bar{\rho}, u, a, \omega) \leq \delta_1 \leq \delta_0.$$

Since  $N_{\tau,\tau}(\rho - \bar{\rho}, u, a, \omega) \leq N_{0,\tau}(\rho - \bar{\rho}, u, a, \omega)$ , we can apply again proposition 7. Our solution can be extended on  $(\tau, 2\tau)$  and  $N_{\tau,2\tau}(\rho - \bar{\rho}, u, a, \omega) \leq C_0 N_{0,\tau}(\rho - \bar{\rho}, u, a, \omega)$ . Thus

$$N_{0,2\tau}^2(\rho - \bar{\rho}, u, a, \omega) = (N_{0,\tau}^2 + N_{\tau,2\tau}^2)(\rho - \bar{\rho}, u, a, \omega) \leq (1 + C_0^2) N_{0,\tau}^2(\rho - \bar{\rho}, u, a, \omega).$$

Thanks to Proposition 8 for  $T = \tau$  and the choice of  $N_{0,0}(\rho - \bar{\rho}, u, a, \omega)$ , we deduce

$$N_{0,2\tau}(\rho - \bar{\rho}, u, a, \omega) \leq C_1 \sqrt{1 + C_0^2} N_{0,0}(\rho - \bar{\rho}, u, a, \omega) \leq \delta_1.$$

This allows to repeat this process and obtain the existence of a regular solution as long as (17) is satisfied.



### 3 A local existence result: proof of Proposition 7

This section is devoted to the proof of Proposition 7. We only consider the case  $h = 0$

#### 3.1 Statement of the problem in a fixed domain

The system of equations can be written on the reference domains  $\Omega_S(0)$  and  $\Omega_F(0)$  with the help of the flow defined by Lemma 1. We consider  $(\rho, u, a, \omega)$  which satisfies the system (2)-(11)-(12) with the hypothesis (3)-(4). We suppose that (17) is satisfied for some  $\alpha > 0$  and that the functions  $\rho, u, a$  and  $\omega$  are regular enough. Let us define the functions  $\tilde{u}, \tilde{\rho}$  and  $\tilde{p}$  on  $V_T := (0, T) \times \Omega_F(0)$  by

$$\tilde{u}(t, y) = u(t, \chi(t, y)), \tilde{\rho}(t, y) = \rho(t, \chi(t, y)) - \bar{\rho}, \tilde{p}(t, y) = P(\tilde{\rho}(t, y) + \bar{\rho}), \forall (t, y) \in V_T. \quad (31)$$

We have the following formulas,  $\forall (t, y) \in V_T$ ,

$$\begin{aligned} \tilde{\rho}_t(t, y) &= \rho_t(t, \chi(t, y)) + (((\nabla\chi)^{-1}\chi_t) \cdot \nabla\tilde{\rho})(t, y) \\ (\nabla\tilde{u})_{ij}(t, y) &= \partial_{y_j}\tilde{u}_i = (\nabla u(t, \chi(t, y))\nabla\chi(t, y))_{ij}. \end{aligned}$$

Thus we get, for the first equation of system (2),

$$\tilde{\rho}_t + ((\nabla\chi)^{-1}(\tilde{u} - \chi_t)) \cdot \nabla\tilde{\rho} + (\tilde{\rho} + \bar{\rho}) \operatorname{tr}(\nabla\tilde{u}(\nabla\chi)^{-1}) = 0, \text{ in } V_T. \quad (32)$$

Next, the second equation of (2) becomes,  $\forall i = 1, 2, 3$ , in  $V_T$

$$\begin{aligned} (\tilde{\rho} + \bar{\rho})(\tilde{u}_i)_t + (\tilde{\rho} + \bar{\rho})(\tilde{u}_j - \chi_{j,t})(\nabla\tilde{u}(\nabla\chi)^{-1})_{ij} - \mu\partial_{y_l}(\partial_{y_k}\tilde{u}_i(\nabla\chi)^{-1}_{kj})(\nabla\chi)^{-1}_{lj} \\ - (\mu + \mu')\partial_{y_l}(\partial_{y_k}\tilde{u}_j(\nabla\chi)^{-1}_{kj})(\nabla\chi)^{-1}_{li} + (\nabla\chi)^{-1}_{ki}\partial_{y_k}\tilde{p} = 0. \end{aligned} \quad (33)$$

In this equation and in what follows, we implicitly sum over repeated indexes. Equations (11) become:

$$\left\{ \begin{aligned} m\ddot{a} &= \int_{\partial\Omega_S(0)} \left( \mu(\nabla\tilde{u}Q^t + Q(\nabla\tilde{u})^t) + \mu' \operatorname{tr}(\nabla\tilde{u}Q^t)Id - \tilde{p}Id \right) Qn \, d\gamma, \\ J\dot{\omega} &= (J\omega) \wedge \omega + \int_{\partial\Omega_S(0)} (Qy) \wedge \left( (\mu(\nabla\tilde{u}Q^t + Q(\nabla\tilde{u})^t) + \mu' \operatorname{tr}(\nabla\tilde{u}Q^t)Id - \tilde{p}Id) Qn \right) d\gamma, \end{aligned} \right. \quad (34)$$

where we have denoted  $A^{-t} = (A^t)^{-1}$ . At last, the boundary conditions (12) become

$$\left\{ \begin{aligned} \tilde{u}(t, y) &= 0, \forall y \in \partial\Omega, \\ \tilde{u}(t, y) &= \dot{a}(t) + \omega(t) \wedge (Qy), \forall y \in \partial\Omega_S(0). \end{aligned} \right. \quad (35)$$

We denote again  $(m, J, \mu, \mu')$  instead of  $(m/\bar{\rho}, J/\bar{\rho}, \mu/\bar{\rho}, \mu'/\bar{\rho})$  and we rewrite the resulting system as follows:

$$\left\{ \begin{aligned} \tilde{\rho}_t + ((\nabla\chi)^{-1}(\tilde{u} - \chi_t)) \cdot \nabla\tilde{\rho} + \bar{\rho} \nabla \cdot \tilde{u} &= g_0(\tilde{\rho}, \tilde{u}, a, \omega) && \text{in } V_T, \\ \tilde{u}_t - \nabla \cdot \sigma(\tilde{u}, \tilde{\rho}) &= g_1(\tilde{\rho}, \tilde{u}, a, \omega) && \text{in } V_T, \\ m\ddot{a} &= \int_{\partial\Omega_S(0)} \sigma(\tilde{u}, \tilde{\rho})n \, d\gamma + g_2(\tilde{\rho}, \tilde{u}, a, \omega) && \text{in } (0, T), \\ J\dot{\omega} &= \int_{\partial\Omega_S(0)} (Qy) \wedge (\sigma(\tilde{u}, \tilde{\rho})n) \, d\gamma + g_3(\tilde{\rho}, \tilde{u}, a, \omega) && \text{in } (0, T), \\ \tilde{u} &= 0 && \text{on } (0, T) \times \partial\Omega, \\ \tilde{u} &= \dot{a} + \omega \wedge (Qy) && \text{on } (0, T) \times \partial\Omega_S(0), \\ \tilde{\rho}(0, \cdot) &= \rho_0 - \bar{\rho}, \quad \tilde{u}(0, \cdot) = u_0 && \text{in } \Omega_F(0), \\ a(0) &= 0, \quad \dot{a}(0) = a_0, \quad \omega(0) = \omega_0, && \end{aligned} \right. \quad (36)$$

with

$$g_0(\tilde{\rho}, \tilde{u}, a, \omega) = \bar{\rho} \operatorname{tr}(\nabla \tilde{u}(Id - (\nabla \chi)^{-1})) - \tilde{\rho} \operatorname{tr}(\nabla \tilde{u}(\nabla \chi)^{-1}), \quad (37)$$

$$\begin{aligned} (g_1)_i(\tilde{\rho}, \tilde{u}, a, \omega) &= -(\tilde{u}_j - \chi_{j,t})(\nabla \tilde{u}(\nabla \chi)^{-1})_{ij} + \mu \left( \frac{\bar{\rho}}{\tilde{\rho} + \bar{\rho}} - 1 \right) \partial_{y_i} (\partial_{y_k} \tilde{u}_i (\nabla \chi)_{kj}^{-1}) (\nabla \chi)_{lj}^{-1} \\ &\quad + \mu [\partial_{y_i} (\partial_{y_k} \tilde{u}_i ((\nabla \chi)_{kj}^{-1} - \delta_{kj})) (\nabla \chi)_{lj}^{-1} + \partial_{y_i y_j}^2 \tilde{u}_i ((\nabla \chi)_{lj}^{-1} - \delta_{lj})] \\ &\quad + (\mu + \mu') \left( \frac{\bar{\rho}}{\tilde{\rho} + \bar{\rho}} - 1 \right) \partial_{y_i} (\partial_{y_k} \tilde{u}_j (\nabla \chi)_{kj}^{-1}) (\nabla \chi)_{li}^{-1} \\ &\quad + (\mu + \mu') [\partial_{y_i} (\partial_{y_k} \tilde{u}_j ((\nabla \chi)_{kj}^{-1} - \delta_{kj})) (\nabla \chi)_{li}^{-1} + \partial_{y_i y_j}^2 \tilde{u}_j ((\nabla \chi)_{li}^{-1} - \delta_{li})] \\ &\quad - ((\nabla \chi)_{ki}^{-1} - \delta_{ki}) \frac{P'(\tilde{\rho} + \bar{\rho}) \partial_{y_k} \tilde{\rho}}{\tilde{\rho} + \bar{\rho}} - \left( \frac{P'(\tilde{\rho} + \bar{\rho})}{\tilde{\rho} + \bar{\rho}} - \frac{P'(\bar{\rho})}{\bar{\rho}} \right) \partial_{y_i} \tilde{\rho}, \end{aligned} \quad (38)$$

$$\begin{aligned} g_2(\tilde{\rho}, \tilde{u}, a, \omega) &= \mu \int_{\partial \Omega_S(0)} [(Q - Id)(\nabla \tilde{u})^t Q + (\nabla \tilde{u})^t (Q - Id)] n \, d\gamma \\ &\quad + \mu' \int_{\partial \Omega_S(0)} [\operatorname{tr}(\nabla \tilde{u}(Q^t - Id))Q + (\nabla \cdot \tilde{u})(Q - Id)] n \, d\gamma \\ &\quad - \int_{\partial \Omega_S(0)} \frac{P(\tilde{\rho} + \bar{\rho}) - P(\bar{\rho})}{\bar{\rho}} (Q - Id) n \, d\gamma - \int_{\partial \Omega_S(0)} (P(\tilde{\rho} + \bar{\rho}) - P(\bar{\rho}) - P'(\bar{\rho})\tilde{\rho}) \frac{n}{\bar{\rho}} \, d\gamma \end{aligned} \quad (39)$$

and

$$\begin{aligned} g_3(\tilde{\rho}, \tilde{u}, a, \omega) &= \mu \int_{\partial \Omega_S(0)} (Qy) \wedge [(Q - Id)(\nabla \tilde{u})^t Q + (\nabla \tilde{u})^t (Q - Id)] n \, d\gamma \\ &\quad + \mu' \int_{\partial \Omega_S(0)} (Qy) \wedge [\operatorname{tr}(\nabla \tilde{u}(Q^t - Id))Q + \operatorname{tr}(\nabla \tilde{u})(Q - Id)] n \, d\gamma \\ &\quad + \int_{\partial \Omega_S(0)} [(Qy) \wedge (p^0 \tilde{\rho}(Id - Q)n)] \, d\gamma \\ &\quad + \int_{\partial \Omega_S(0)} (P(\bar{\rho}) + P'(\bar{\rho})\tilde{\rho} - P(\tilde{\rho} + \bar{\rho}))(Qy) \wedge Q \frac{n}{\bar{\rho}} \, d\gamma + (J\omega) \wedge \omega. \end{aligned} \quad (40)$$

Here, we have used (29). Observe that the functions  $g_i$  are (at least) quadratic functions of the quantities  $\tilde{\rho}, \tilde{u}, (\nabla \chi)^{-1} - Id, Q - Id, \chi_t$ .

### 3.2 Definition of the fixed point mapping

Let  $0 < \tilde{R} < 1$  and  $0 < s < 1$  be small enough. We define the space  $\tilde{Y}((0, s); \tilde{R})$ :

$$\tilde{Y}((0, s); \tilde{R}) = \left\{ (\tilde{\rho}, \tilde{u}_F, a, \omega) \in \tilde{X}(0, s) : \tilde{u}_F = 0 \text{ on } \partial \Omega_F(0), a(0) = 0, Q(0) = Id, \tilde{N}_{0,s}(\tilde{\rho}, \tilde{u}_F, a, \omega) \leq \tilde{R} \right\},$$

where

$$\begin{aligned} \tilde{X}(0, s) &= \{(\tilde{\rho}, \tilde{u}, a, \omega) : \tilde{\rho} \in C_s^0(H^3(\Omega_F(0))) \cap C_s^1(H^2(\Omega_F(0))) \cap H_s^2(L^2(\Omega_F(0))), \\ &\quad \tilde{u} \in L_s^2(H^4(\Omega_F(0))) \cap H_s^2(L^2(\Omega_F(0))), a \in H_s^3, \omega \in H_s^2\} \end{aligned}$$

and  $\tilde{N}_{0,s}(\tilde{\rho}, \tilde{u}_F, a, \omega)$  is given by the corresponding norm. In this definition,  $Q$  is the rotation matrix associated to  $\omega$ .

**Remark 9** *The definition of the space  $\tilde{Y}((0, s); \tilde{R})$  coincides with*

$$Y((0, s); R) = \left\{ (\tilde{\rho} \circ \chi^{-1}, \tilde{u}_F \circ \chi^{-1}, a, \omega) \in X(0, s) / a(0) = 0, Q(0) = Id, N_{0,s}(\tilde{\rho} \circ \chi^{-1}, \tilde{u}_F \circ \chi^{-1}, a, \omega) \leq R \right\}.$$

This comes from the identities, for all  $(t, y) \in (0, s) \times \Omega_F(0)$ , if  $x = \chi(t, y)$ ,

$$\partial_t \tilde{u}(t, y) = \partial_t u(t, x) + (\chi_t(t, y) \cdot \nabla) u(t, x)$$

and

$$\nabla \tilde{u}(t, y) = \nabla u(t, x) \nabla \chi(t, y).$$

In order to prove Proposition 7, we will perform a fixed point argument in the space  $\tilde{Y}((0, s); \tilde{R})$ . Indeed, we consider the mapping

$$\begin{aligned} \Lambda : \tilde{Y}((0, s); \tilde{R}) &\rightarrow \tilde{Y}((0, s); \tilde{R}) \\ (\hat{\rho}, \hat{u}_F, \hat{a}, \hat{\omega}) &\rightarrow (\tilde{\rho}, \tilde{u}_F, a, \omega), \end{aligned}$$

with  $\tilde{u}_F = \tilde{u} - \chi_t$ , where  $(\tilde{\rho}, \tilde{u}, a, \omega)$  is the solution of

$$\left\{ \begin{array}{ll} \tilde{\rho}_t + ((\nabla \hat{\chi})^{-1}(\hat{u} - \hat{\chi}_t)) \cdot \nabla \tilde{\rho} = g_0(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega}) - \bar{\rho} \nabla \cdot \tilde{u} & \text{in } V_T, \\ \tilde{u}_t - \nabla \cdot (2\mu \epsilon(\tilde{u}) + \mu'(\nabla \cdot \tilde{u}) Id) = g_1(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega}) - p^0 \nabla \tilde{\rho} & \text{in } V_T, \\ m\ddot{a} = \int_{\partial\Omega_S(0)} \sigma(\tilde{u}, \tilde{\rho}) n \, d\gamma + g_2(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega}) & \text{in } (0, T), \\ \hat{J}\dot{\omega} = \int_{\partial\Omega_S(0)} (\hat{Q}y) \wedge (\sigma(\tilde{u}, \tilde{\rho}) n) \, d\gamma + g_3(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega}) & \text{in } (0, T), \\ \tilde{u} = 0 & \text{on } (0, T) \times \partial\Omega, \\ \tilde{u} = \dot{a} + \omega \wedge (\hat{Q}y) & \text{on } (0, T) \times \partial\Omega_S(0), \\ \tilde{\rho}(0, \cdot) = \rho_0 - \bar{\rho}, \quad \tilde{u}(0, \cdot) = u_0 & \text{in } \Omega_F(0), \\ a(0) = 0, \quad \dot{a}(0) = a_0, \quad \omega(0) = \omega_0. & \end{array} \right. \quad (41)$$

Here,  $\hat{u} = \hat{u}_F + \hat{\chi}_t$  and  $\hat{J}$  is defined by (9) where we replace  $Q$  by  $\hat{Q}$ .

Thanks to Remark 9, if we prove the existence of a fixed point of  $\Lambda$ , the proof of Proposition 7 will be achieved.

### 3.3 $\Lambda$ is well-defined

Let  $(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})$  be given in  $\tilde{Y}((0, s); \tilde{R})$ . Our goal is to prove that the solution of (41)  $(\tilde{\rho}, \tilde{u}, a, \omega)$  belongs to  $\tilde{Y}((0, s); \tilde{R})$ . In order to prove this, we will establish some estimates for  $\tilde{\rho}$  regarded as the solution of a transport equation (with right hand side  $g_0 - \bar{\rho}(\nabla \cdot \tilde{u})$ ) and for  $\tilde{u}$  as solution of a heat equation (with right hand side  $g_1 - p^0 \nabla \tilde{\rho}$ ).

In the sequel, we will use the following estimates, coming from Lemma 1:

$$\|(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})\|_{\tilde{X}(0, s)} + \|(\nabla \hat{\chi})^{-1} - Id\|_{H_s^3(C^\infty(\overline{\Omega_F(0)}))} + \|\hat{Q} - Id\|_{H_s^3} + \|\hat{\chi}_t\|_{H_s^2(C^\infty(\overline{\Omega_F(0)})} \leq C\tilde{R}.$$

Taking  $\tilde{R} > 0$  small enough, we can suppose that  $|\hat{\rho}| < \bar{\rho}/2$  and so every single term of the expression of  $g_1$  (see (38)) makes sense. Then, one can prove that

$$\begin{aligned} &\|g_0(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})\|_{L_s^2(H^3(\Omega_F(0))) \cap H_s^1(H^1(\Omega_F(0)))} + \|g_1(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})\|_{L_s^2(H^2(\Omega_F(0))) \cap H_s^1(L^2(\Omega_F(0)))} \\ &+ \|g_2(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})\|_{H_s^1} + \|g_3(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})\|_{H_s^1} \leq C\tilde{R}^2. \end{aligned} \quad (42)$$

#### Estimates for $\tilde{\rho}$

Using standard arguments (observe that  $\hat{u} - \hat{\chi}_t$  has null trace), we can prove using (42) that

$$\begin{aligned} \|\tilde{\rho}\|_{L_s^\infty(H^3(\Omega_F(0)))} &\leq C(\|g_0(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})\|_{L_s^2(H^3(\Omega_F(0)))} + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + \|\tilde{u}\|_{L_s^2(H^4(\Omega_F(0)))}) \\ &\leq C(\tilde{R}^2 + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + \tilde{N}_{0, s}(0, \tilde{u}, 0, 0)). \end{aligned} \quad (43)$$

Moreover, using the equation of  $\tilde{\rho}$  and (43), we obtain

$$\begin{aligned} & \|\tilde{\rho}\|_{W_s^{1,\infty}(H^2(\Omega_F(0)))} + \|\tilde{\rho}\|_{H_s^2(L^2(\Omega_F(0)))} \leq C(\|g_0(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})\|_{L_s^\infty(H^2(\Omega_F(0))) \cap H_s^1(L^2(\Omega_F(0)))}) \\ & + \|\tilde{u}\|_{L_s^\infty(H^3(\Omega_F(0))) \cap H_s^1(H^1(\Omega_F(0)))} \leq C(\tilde{R}^2 + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + \tilde{N}_{0,s}(0, \tilde{u}, 0, 0)). \end{aligned}$$

Combining both, we deduce

$$\tilde{N}_{0,s}(\tilde{\rho}, 0, 0, 0) \leq C(\tilde{R}^2 + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + \tilde{N}_{0,s}(0, \tilde{u}, 0, 0)). \quad (44)$$

### Estimates for $\tilde{u}$ , $a$ and $\omega$

In the rest of this subsection, we denote  $g_i$  instead of  $g_i(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega})$  for  $i = 1, 2, 3$ .

• First, we multiply the equation of  $\tilde{u}$  by  $\tilde{u}$ , we integrate on  $V_r$  for all  $r \in (0, s)$  and we integrate by parts. Taking the supremum in  $r$ , this gives

$$\begin{aligned} & \frac{1}{2} \sup_{r \in (0, s)} \left( |\dot{a}|^2(r) + |\omega|^2(r) + \int_{\Omega_F(0)} |\tilde{u}|^2(r) dy \right) + \iint_{V_s} (2\mu|\epsilon(\tilde{u})|^2 + \mu'|\nabla \cdot \tilde{u}|^2) dy dr \\ & \leq C \left( \iint_{V_s} |\tilde{u}|^2 dy dr + \int_0^s |\omega|^2 dr + \iint_{V_s} (|g_1|^2 + |\nabla \tilde{\rho}|^2) dy dr + \int_{\Omega_F(0)} |u_0|^2 dx + |a_0|^2 + |\omega_0|^2 \right. \\ & \left. + \left( \sup_{r \in (0, s)} |\dot{a}| \right) \left( \int_0^s |g_2| dr + \int_0^s \int_{\partial\Omega_S(0)} |\tilde{\rho}| d\gamma dr \right) + \left( \sup_{r \in (0, s)} |\omega| \right) \left( \int_0^s |g_3| dr + \int_0^s \int_{\partial\Omega_S(0)} |\tilde{\rho}| d\gamma dr \right) \right). \end{aligned} \quad (45)$$

Here, we have used that  $\hat{J}$  is coercive (see 10) and satisfies  $\|\hat{J}\|_{L^\infty} \leq C$ .

Now, observe that from the fact that  $\Omega_S(0)$  is convex (see (1)) and thanks to  $\tilde{u} = 0$  on  $\partial\Omega \times (0, s)$ , we have that

$$\iint_{V_s} |\epsilon(\tilde{u})|^2 dy dr \geq C\|\tilde{u}\|_{L_s^2(H^1(\Omega_F(0)))}^2, \quad (46)$$

for some  $C > 0$ . Indeed, first we have  $\sum_{i=1}^3 |\partial_{y_i} \tilde{u}_i|^2 \leq |\epsilon(\tilde{u})|^2$ . Then, for instance for  $\tilde{u}_1$  we notice that

$$\tilde{u}_1(y) = \tilde{u}_1(y_1, y_2, y_3) = \int_{y^*}^{y_1} \partial_{s_1} \tilde{u}_1(s_1, y_2, y_3) ds_1 \quad \forall y \in \Omega_F(0),$$

where  $y^* \in \partial\Omega \cap \{y + \lambda e_1 : \lambda \in \mathbb{R}\}$ ; this is possible since  $\Omega_S(0)$  is convex. Of course, the same is also true for  $\tilde{u}_2$  and  $\tilde{u}_3$ , so we deduce that

$$\int_{\Omega_F(0)} |\tilde{u}|^2 dy \leq C \sum_{i=1}^3 \int_{\Omega_F(0)} |\partial_{y_i} \tilde{u}_i|^2 dy \leq C \int_{\Omega_F(0)} |\epsilon(\tilde{u})|^2 dy.$$

Finally, since  $\tilde{u}, \epsilon(\tilde{u}) \in L^2(\Omega_F(0))$  and  $\tilde{u} = 0$  on a part of the boundary, Korn's inequality tells that  $\tilde{u} \in H^1(\Omega_F(0))$  so we obtain (46).

Thus, using estimate (44) on  $\tilde{\rho}$  and (42), we deduce that there exists  $C > 0$  such that

$$\begin{aligned} & \|\tilde{u}\|_{L_s^\infty(L^2(\Omega_F(0)))} + \|\tilde{u}\|_{L_s^2(H^1(\Omega_F(0)))} + \|a\|_{W_s^{1,\infty}} + \|\omega\|_{L_s^\infty} \\ & \leq C(\tilde{R}^2 + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + s^{1/2} \tilde{N}_{0,s}(0, \tilde{u}, 0, 0) + \|u_0\|_{L^2} + |a_0| + |\omega_0|). \end{aligned} \quad (47)$$

• Let us multiply the equation of  $\tilde{u}$  by  $\tilde{u}_t$ . Arguing as before, this yields

$$\begin{aligned} & \frac{1}{2} \sup_{r \in (0, s)} \left( \int_{\Omega_F(0)} (2\mu|\epsilon(\tilde{u})|^2(r) + \mu'|\nabla \cdot \tilde{u}|^2(r)) dy \right) + \iint_{V_s} |\tilde{u}_t|^2 dy dr \\ & + \int_0^s \int_{\partial\Omega_S(0)} ((2\mu\epsilon(\tilde{u}) + \mu'(\nabla \cdot \tilde{u})Id)n) \cdot \tilde{u}_t d\gamma dr \leq C \left( \iint_{V_s} (|g_1|^2 + |\nabla \tilde{\rho}|^2) dy dr + \int_{\Omega_F(0)} |\nabla u_0|^2 dy \right) \end{aligned} \quad (48)$$

Using the boundary condition of  $\tilde{u}$  on  $\partial\Omega_S(0)$ , we obtain the following for the boundary term:

$$\begin{aligned} & \int_0^s \int_{\partial\Omega_S(0)} ((2\mu\epsilon(\tilde{u}) + \mu'(\nabla \cdot \tilde{u})Id)n) \cdot \tilde{u}_t \, d\gamma dr \\ &= \int_0^s \int_{\partial\Omega_S(0)} ((2\mu\epsilon(\tilde{u}) + \mu'(\nabla \cdot \tilde{u})Id)n) \cdot (\ddot{u} + \dot{\omega} \wedge (\hat{Q}y) + \omega \wedge (\hat{Q}\dot{y})) \, d\gamma dr. \end{aligned}$$

For the two first terms we use the equations of  $a$  and  $\omega$  in (41). This produces the norms  $\|\ddot{u}\|_{L_s^2}^2$  and  $\|\dot{\omega}\|_{L_s^2}^2$ . For the third term, we have

$$\begin{aligned} \left| \int_0^s \int_{\partial\Omega_S(0)} ((2\mu\epsilon(\tilde{u}) + \mu'(\nabla \cdot \tilde{u})Id)n) \cdot (\omega \wedge \hat{Q}\dot{y}) \, d\gamma dr \right| &\leq C\tilde{R}\|\omega\|_{L_s^2} \|\tilde{u}\|_{L_s^2(H^2(\Omega_F(0)))} \\ &\leq \tilde{R}^2 \|\tilde{u}\|_{L_s^2(H^2(\Omega_F(0)))}^2 + C\|\omega\|_{L_s^2}^2. \end{aligned}$$

Using (47) to estimate  $\|\omega\|_{L_s^2}^2$ , (42) and (44), we obtain from (48)

$$\begin{aligned} & \|\tilde{u}\|_{H_s^1(L^2(\Omega_F(0)))} + \|\tilde{u}\|_{L_s^\infty(H^1(\Omega_F(0)))} + \|a\|_{H_s^2} + \|\omega\|_{H_s^1} \leq \tilde{R}\|\tilde{u}\|_{L_s^2(H^2(\Omega_F(0)))} \\ & + C(\tilde{R}^2 + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + s^{1/2}\tilde{N}_{0,s}(0, \tilde{u}, 0, 0) + \|u_0\|_{H^1(\Omega_F(0))} + |a_0| + |\omega_0|). \end{aligned} \quad (49)$$

Now, we regard the equation of  $\tilde{u}$  as a stationary system:

$$\begin{cases} -\mu\Delta\tilde{u} - (\mu + \mu')\nabla(\nabla \cdot \tilde{u}) = g_1 - p^0\nabla\tilde{\rho} - \tilde{u}_t & \text{in } \Omega_F(0), \\ \tilde{u} = (\dot{a} + \omega \wedge (\hat{Q}y))1_{\partial\Omega_S(0)} & \text{on } \partial\Omega_F(0). \end{cases} \quad (50)$$

The solution of this system belongs to  $H^2$  and we have

$$\|\tilde{u}\|_{H^2(\Omega_F(0))} \leq C(\|g_1\|_{L^2(\Omega_F(0))} + \|\tilde{\rho}\|_{H^1(\Omega_F(0))} + \|\tilde{u}_t\|_{L^2(\Omega_F(0))} + |\dot{a}| + |\omega|).$$

Taking here the  $L_s^2$  norm and using (47) to estimate  $\|\dot{a}\|_{L_s^2} + \|\omega\|_{L_s^2}$  and (43) to estimate  $\|\tilde{\rho}\|_{L_s^2(H^1(\Omega_F(0)))}$ , we obtain from (49)

$$\begin{aligned} & \|\tilde{u}\|_{H_s^1(L^2(\Omega_F(0)))} + \|\tilde{u}\|_{L_s^2(H^2(\Omega_F(0)))} + \|\tilde{u}\|_{L_s^\infty(H^1(\Omega_F(0)))} + \|a\|_{H_s^2} + \|\omega\|_{H_s^1} \\ & \leq C(\tilde{R}^2 + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + s^{1/2}\tilde{N}_{0,s}(0, \tilde{u}, 0, 0) + \|u_0\|_{H^1(\Omega_F(0))} + |a_0| + |\omega_0|) \end{aligned} \quad (51)$$

(recall that  $R > 0$  is small enough).

- If we differentiate with respect to time the system satisfied by  $(\tilde{u}, a, \omega)$ , we obtain

$$\left\{ \begin{array}{ll} \tilde{u}_{tt} - 2\mu\nabla \cdot (\epsilon(\tilde{u}_t)) - \mu'\nabla \cdot ((\nabla \cdot \tilde{u}_t)Id) = g_{1,t} - p^0\nabla\tilde{\rho}_t & \text{in } V_T, \\ m\ddot{a} = \int_{\partial\Omega_S(0)} (2\mu\epsilon(\tilde{u}_t) + \mu'(\nabla \cdot \tilde{u}_t)Id - p^0\tilde{\rho}_t Id)n \, d\gamma + g_{2,t} & \text{in } (0, T), \\ \hat{J}\ddot{\omega} = -\hat{J}\dot{\omega} + \int_{\partial\Omega_S(0)} (\hat{Q}x) \wedge ((2\mu\epsilon(\tilde{u}_t) + (\mu'\nabla \cdot \tilde{u}_t - p^0\tilde{\rho}_t)Id)n) \, d\gamma \\ \quad + \int_{\partial\Omega_S(0)} (\hat{Q}x) \wedge ((2\mu\epsilon(\tilde{u}) + \mu'(\nabla \cdot \tilde{u})Id - p^0\tilde{\rho}Id)n) \, d\gamma + g_{3,t} & \text{in } (0, T), \\ \tilde{u}_t = 0 & \text{on } (0, T) \times \partial\Omega, \\ \tilde{u}_t = \ddot{u} + \dot{\omega} \wedge (\hat{Q}y) + \omega \wedge (\hat{Q}\dot{y}) & \text{on } (0, T) \times \partial\Omega_S(0). \end{array} \right. \quad (52)$$

Let us multiply the first equation by  $\tilde{u}_t$  and integrate in space and in time as before. Doing similar compu-

tations to those of the previous case, we obtain

$$\begin{aligned}
& \sup_{r \in (0,s)} \left( |\ddot{a}|^2(r) + |\dot{\omega}|^2(r) + \int_{\Omega_F(0)} |\tilde{u}_t|^2(r) dy \right) + \iint_{Q_s} (2\mu|\epsilon(\tilde{u}_t)|^2 + \mu'|\nabla \cdot \tilde{u}_t|^2) dy dr \\
& \leq C \left( \iint_{Q_s} (|g_{1,t}|^2 + |\tilde{u}_t|^2 + |\nabla \tilde{\rho}_t|^2) dy dr + \int_{\Omega_F(0)} |\tilde{u}_t(0)|^2 dy + |\ddot{a}(0)|^2 + |\dot{\omega}(0)|^2 \right. \\
& \quad \left. + \int_0^s (|\omega|^2 + |\dot{\omega}|^2) dr + \|u\|_{L_s^2(H^2(\Omega_F(0)))}^2 + \left( \sup_{r \in (0,s)} |\ddot{a}| \right) \left( \int_0^s |g_{2,t}| dr + \int_0^s \int_{\partial\Omega_S(0)} |\tilde{\rho}_t| d\sigma dr \right) \right. \\
& \quad \left. + \left( \sup_{r \in (0,s)} |\dot{\omega}| \right) \left( \int_0^s |g_{3,t}| dr + \int_0^s \int_{\partial\Omega_S(0)} (|\tilde{\rho}| + |\tilde{\rho}_t|) d\sigma dr \right) \right) + \tilde{R}^2 \|\tilde{u}_t\|_{L_s^2(H^2(\Omega_F(0)))}^2.
\end{aligned}$$

The last term can be bounded by  $\tilde{R}^2 \tilde{N}_{0,s}^2(0, \tilde{u}, 0, 0)$ . Using (42), (44) and (51), this implies that

$$\begin{aligned}
& \|\tilde{u}\|_{W_s^{1,\infty}(L^2(\Omega_F(0)))} + \|\tilde{u}\|_{H_s^1(H^1(\Omega_F(0)))} + \|a\|_{W_s^{2,\infty}} + \|\omega\|_{W_s^{1,\infty}} \leq C \left( \tilde{R}^2 + |\ddot{a}(0)| + |\dot{\omega}(0)| \right. \\
& \quad \left. + \|\tilde{u}_t(0)\|_{L^2(\Omega_F(0))} + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + (s^{1/2} + \tilde{R})\tilde{N}_{0,s}(0, \tilde{u}, 0, 0) + \|u_0\|_{H^1(\Omega_F(0))} + |a_0| + |\omega_0| \right). \tag{53}
\end{aligned}$$

• Next, if we multiply the first equation of (52) by  $\tilde{u}_{tt}$ , we obtain that

$$\begin{aligned}
& \iint_{V_s} |\tilde{u}_{tt}|^2 dy dr + \frac{1}{2} \sup_{r \in (0,s)} \int_{\Omega_F(0)} (2\mu|\epsilon(\tilde{u}_t)|^2(r) + \mu'|\nabla \cdot \tilde{u}_t|^2(r)) dy \\
& \quad + \int_0^s \int_{\partial\Omega_S(0)} \left( (2\mu\epsilon(\tilde{u}_t) + \mu'(\nabla \cdot \tilde{u}_t)Id)n \right) \cdot \tilde{u}_{tt} d\gamma dr = \iint_{V_s} g_{1,t} \cdot \tilde{u}_{tt} dy dr \\
& \quad + \frac{1}{2} \int_{\Omega_F(0)} (2\mu|\epsilon(\tilde{u}_t)|^2(0) + \mu'|\nabla \cdot \tilde{u}_t|^2(0)) dy - \iint_{V_s} p^0 \nabla \tilde{\rho}_t \cdot \tilde{u}_{tt} dy dr.
\end{aligned}$$

On  $(0, T) \times \partial\Omega_S(0)$  we have

$$u_{tt} = \ddot{a} + \ddot{\omega} \wedge (\hat{Q}y) + 2\dot{\omega} \wedge (\dot{Q}y) + \omega \wedge (\ddot{Q}y). \tag{54}$$

Thus, we deduce

$$\begin{aligned}
& \int_0^s \int_{\partial\Omega_S(0)} \left( (2\mu\epsilon(\tilde{u}_t) + \mu'(\nabla \cdot \tilde{u}_t)Id)n \right) \cdot \tilde{u}_{tt} d\gamma dr = \int_0^s \ddot{a} \cdot (m\ddot{a} - g_{2,t}) dr + \int_0^s \ddot{a} \cdot \int_{\partial\Omega_S(0)} p^0 \tilde{\rho}_t n d\gamma dr \\
& \quad + \int_0^s \ddot{\omega} \cdot (\hat{J}\ddot{\omega} + \dot{J}\dot{\omega} - g_{3,t}) dr + \int_0^s \ddot{\omega} \cdot \int_{\partial\Omega_S(0)} (\hat{Q}y) \wedge (p^0 \tilde{\rho}_t n) d\gamma dr \\
& \quad - \int_0^s \ddot{\omega} \cdot \int_{\partial\Omega_S(0)} (\dot{Q}y) \wedge \left( (2\mu\epsilon(\tilde{u}) + \mu'(\nabla \cdot \tilde{u})Id - p^0 \tilde{\rho}Id)n \right) d\gamma dr \\
& \quad + \int_0^s \int_{\partial\Omega_S(0)} \left( (2\mu\epsilon(\tilde{u}_t) + \mu'(\nabla \cdot \tilde{u}_t)Id)n \right) \cdot (2\dot{\omega} \wedge (\dot{Q}y) + \omega \wedge (\ddot{Q}y)) d\gamma dr.
\end{aligned}$$

Then, proceeding as before and using (53), (51) and (44), we obtain that there exists  $C$  such that

$$\begin{aligned}
& \|\tilde{u}_{tt}\|_{L_s^2(L^2(\Omega_F(0)))} + \|\tilde{u}_t\|_{L_s^\infty(H^1(\Omega_F(0)))} + \|\ddot{a}\|_{L_s^2} + \|\ddot{\omega}\|_{L_s^2} \leq C \left( |\ddot{a}(0)| + |\dot{\omega}(0)| + \|u_0\|_{H^1(\Omega_F(0))} \right. \\
& \quad \left. + \tilde{R}^2 + \|\tilde{u}_t(0)\|_{H^1(\Omega_F(0))} + \|\rho_0 - \bar{\rho}\|_{H^3(\Omega_F(0))} + (s^{1/2} + \tilde{R})\tilde{N}_{0,s}(0, \tilde{u}, 0, 0) + |a_0| + |\omega_0| \right). \tag{55}
\end{aligned}$$

According to the elliptic equation satisfied by  $\tilde{u}_t$  (see (50)), we have

$$\|\tilde{u}_t\|_{L_s^2(H^2(\Omega_F(0)))} \leq C \left( \|g_{1,t}\|_{L_s^2(L^2(\Omega_F(0)))} + \|\tilde{\rho}_t\|_{L_s^2(H^1(\Omega_F(0)))} + \|\tilde{u}_{tt}\|_{L_s^2(L^2(\Omega_F(0)))} + \|\ddot{a}\|_{L_s^2} + \|\dot{\omega}\|_{H_s^1} \right). \tag{56}$$

At last, if we consider system (50) satisfied by  $\tilde{u}$ , since the right-hand side  $g_1 - p^0 \nabla \tilde{\rho} - \tilde{u}_t$  is estimated in the  $L_s^2(H^2(\Omega_F(0))) \cap L_s^\infty(H^1(\Omega_F(0)))$ -norm (see (55)-(56)) and the boundary condition is estimated in  $L_s^\infty(C^\infty(\partial\Omega_F(0)))$ ,  $\tilde{u}$  is estimated in  $L_s^2(H^4(\Omega_F(0))) \cap L_s^\infty(H^3(\Omega_F(0)))$  by

$$\tilde{R}^2 + \|\rho_0 - \tilde{\rho}\|_{H^3(\Omega_F(0))} + (s^{1/2} + \tilde{R})\tilde{N}_{0,s}(0, \tilde{u}, 0, 0) + \|u_0\|_{H^3(\Omega_F(0))} + |a_0| + |\omega_0|.$$

Thus, we deduce that

$$\tilde{N}_{0,s}(0, \tilde{u}, a, \omega) \leq C(\tilde{R}^2 + (s^{1/2} + \tilde{R})\tilde{N}_{0,s}(0, \tilde{u}, 0, 0) + \tilde{N}_{0,0}(\tilde{\rho}, \tilde{u}, a, \omega)).$$

Using that  $\tilde{N}_{0,s}(0, \tilde{u}, 0, 0) \leq \tilde{N}_{0,s}(\tilde{\rho}, \tilde{u}, a, \omega)$  and since  $s$  and  $\tilde{R}$  are sufficiently small, we have

$$\tilde{N}_{0,s}(0, \tilde{u}, a, \omega) \leq C(\tilde{R}^2 + \tilde{N}_{0,0}(\tilde{\rho}, \tilde{u}, a, \omega)).$$

Combining with (44) and taking into account that  $\tilde{u}_F = \tilde{u} - \chi_t$ , we finally get

$$\tilde{N}_{0,s}(\tilde{\rho}, \tilde{u}_F, a, \omega) \leq C\tilde{N}_{0,s}(\tilde{\rho}, \tilde{u}, a, \omega) \leq C(\tilde{R}^2 + \tilde{N}_{0,0}(\tilde{\rho}, \tilde{u}, a, \omega)).$$

This allows to assert that, if  $\tilde{N}_{0,0}(\tilde{\rho}, \tilde{u}, a, \omega)$  satisfies

$$\tilde{N}_{0,0}(\tilde{\rho}, \tilde{u}, a, \omega) \sim \tilde{R}^2 \tag{57}$$

(which implies that  $\tilde{N}_{0,0}(\tilde{\rho}, \tilde{u}, a, \omega)$  is small enough), then  $(\tilde{\rho}, \tilde{u}_F, a, \omega)$  belongs to  $Y((0, s); \tilde{R})$ .

### 3.4 Fixed point argument

Let us fix  $\tilde{R}$  such that (57) is satisfied. We will prove that  $\Lambda$  is a continuous mapping in  $Y((0, s); \tilde{R})$  for the norm  $\|\cdot\|$  defined by

$$\|(\rho, u_F, a, \omega)\| = \|\rho\|_{L_s^\infty(H^1(\Omega_F(0)))} + \|u_F\|_{L_s^2(H^2(\Omega_F(0)))} + \|a\|_{H_s^2} + \|\omega\|_{H_s^1}. \tag{58}$$

Once this is done, we apply the Schauder's fixed point Theorem (see, for instance, [23]) and we obtain the existence of a fixed point of  $\Lambda$  which solves system (36); then, using the inverse of the flow  $\chi$ , we obtain a solution of our original system (26), which satisfies the properties of Proposition 7.

In order to prove the continuity, let us consider  $(\hat{\rho}_k, \hat{u}_{F,k}, \hat{a}_k, \hat{\omega}_k)$  a sequence converging for the previous norm to some element  $(\hat{\rho}, \hat{u}_F, \hat{a}, \hat{\omega})$  of  $Y((0, s); \tilde{R})$ . We define

$$(\tilde{\rho}_k, \tilde{u}_{F,k}, a_k, \omega_k) := \Lambda(\hat{\rho}_k, \hat{u}_{F,k}, \hat{a}_k, \hat{\omega}_k) \quad \text{and} \quad (\tilde{\rho}, \tilde{u}_F, a, \omega) := \Lambda(\hat{\rho}, \hat{u}_F, \hat{a}, \hat{\omega}).$$

Let us prove that  $(R_k, U_k, A_k, \Omega_k) := (\tilde{\rho}_k - \tilde{\rho}, \tilde{u}_k - \tilde{u}, a_k - a, \omega_k - \omega)$  converges to zero for the norm (58). This readily implies that  $(\tilde{\rho}_k, \tilde{u}_{F,k}, a_k, \omega_k) \rightarrow (\tilde{\rho}, \tilde{u}_F, a, \omega)$ .

#### Estimates for $R_k$

The equation fulfilled by  $R_k$  is the following one:

$$\begin{cases} R_{k,t} + ((\nabla \hat{\chi}_k)^{-1}(\hat{u}_k - \hat{\chi}_{k,t})) \cdot \nabla R_k = \ell_{0,k}, \\ R_k(0, x) = 0, \end{cases}$$

with

$$\ell_{0,k} = g_0(\hat{\rho}_k, \hat{u}_k, \hat{a}_k, \hat{\omega}_k) - g_0(\hat{\rho}, \hat{u}, \hat{a}, \hat{\omega}) - \tilde{\rho} \nabla \cdot U_k + ((\nabla \hat{\chi})^{-1}(\hat{u} - \hat{\chi}_t) - (\nabla \hat{\chi}_k)^{-1}(\hat{u}_k - \hat{\chi}_{k,t})) \cdot \nabla \tilde{\rho}.$$

We directly get that

$$\|R_k\|_{L_s^\infty(H^1(\Omega_F(0)))} \leq C\|\ell_{0,k}\|_{L_s^2(H^1(\Omega_F(0)))}.$$

Now, we use the expression of  $g_0$  (see (37)) and we deduce

$$\begin{aligned} \|R_k\|_{L_s^\infty(H^1(\Omega_F(0)))} &\leq C(\|U_k\|_{L_s^2(H^2(\Omega_F(0)))} + \|(\nabla \hat{\chi}_k)^{-1} - (\nabla \hat{\chi})^{-1}\|_{L_s^2(H^1(\Omega_F(0)))}) \\ &+ \|\hat{u}_k - \hat{u}\|_{L_s^2(H^2(\Omega_F(0)))} + \|\hat{\chi}_{k,t} - \hat{\chi}_t\|_{L_s^2(H^1(\Omega_F(0)))} + \|\hat{\rho}_k - \hat{\rho}\|_{L_s^2(H^1(\Omega_F(0)))}. \end{aligned} \tag{59}$$

**Estimates for  $(U_k, A_k, \Omega_k)$**

We have the equation:

$$\left\{ \begin{array}{ll} U_{k,t} - \nabla \cdot (2\mu\epsilon(U_k) + \mu'\nabla \cdot U_k Id) = \ell_{1,k} & \text{in } (0, T) \times \Omega_F(0), \\ m\ddot{A}_k = \int_{\partial\Omega_S(0)} (2\mu\epsilon(U_k) + \mu'\nabla \cdot U_k Id)n d\gamma + \ell_{2,k} & \text{in } (0, T), \\ \hat{J}\dot{\Omega}_k = \int_{\partial\Omega_S(0)} (\hat{Q}y) \wedge ((2\mu\epsilon(U_k) + \mu'\nabla \cdot U_k Id)n) d\gamma + \ell_{3,k} & \text{in } (0, T), \\ U_k = 0 & \text{on } (0, T) \times \partial\Omega, \\ U_k = \dot{A}_k + \Omega_k \wedge (\hat{Q}y) + \omega_k \wedge ((\hat{Q}_k - \hat{Q})y) & \text{on } (0, T) \times \partial\Omega_S(0), \\ U_k(0, \cdot) = 0, \quad A_k(0) = 0, \quad \dot{A}_k(0) = 0, \quad \Omega_k(0) = 0, & \end{array} \right. \quad (60)$$

with

$$\begin{aligned} \ell_{1,k} &:= g_1(\hat{\rho}_k, \hat{u}_k, a_k, \omega_k) - g_1(\hat{\rho}, \hat{u}, a, \omega) - p^0 \nabla R_k, \\ \ell_{2,k} &= g_2(\hat{\rho}_k, \hat{u}_k, a_k, \omega_k) - g_2(\hat{\rho}, \hat{u}, a, \omega) - p^0 \int_{\partial\Omega_S(0)} R_k n d\gamma, \\ \ell_{3,k} &= g_3(\hat{\rho}_k, \hat{u}_k, a_k, \omega_k) - g_3(\hat{\rho}, \hat{u}, a, \omega) - p^0 \int_{\partial\Omega_S(0)} (\hat{Q}y) \wedge (R_k n) d\gamma \\ &+ \int_{\partial\Omega_S(0)} ((\hat{Q}_k - \hat{Q})y) \wedge (\sigma(\hat{u}_k, \tilde{\rho}_k)n) d\gamma + (\hat{J} - \hat{J}_k)\dot{\omega}_k. \end{aligned}$$

We multiply the equation of  $U_k$  by  $U_{k,t}$ , we integrate in  $(0, r) \times \Omega_F(0)$  and we take the supremum of  $r \in (0, s)$ . Arguing as for the proof of (49) and using

$$U_{k,t} = \ddot{A}_k + \dot{\Omega}_k \wedge (\hat{Q}y) + \ell_{4,k}$$

with

$$\ell_{4,k} = \Omega_k \wedge (\dot{\hat{Q}}y) + \dot{\omega}_k \wedge ((\hat{Q}_k - \hat{Q})y) + \omega_k \wedge ((\dot{\hat{Q}}_k - \dot{\hat{Q}})y),$$

we obtain

$$\begin{aligned} &\|A_k\|_{H_s^2} + \|\Omega_k\|_{H_s^1} + \|U_k\|_{H_s^1(L^2(\Omega_F(0)))} + \|U_k\|_{L_s^\infty(H^1(\Omega_F(0)))} \\ &\leq \delta \|U_k\|_{L_s^2(H^2(\Omega_F(0)))} + C(\|\ell_{1,k}\|_{L_s^2(L^2(\Omega_F(0)))} + \|\ell_{2,k}\|_{L_s^2} + \|\ell_{3,k}\|_{L_s^2} + \|\ell_{4,k}\|_{L_s^2(L^2(\partial\Omega_S(0)))}), \end{aligned}$$

for  $\delta > 0$  small enough. Regarding now the equation of  $U_k$  as an elliptic equation, we deduce that  $U_k \in L_s^2(H^2(\Omega_F(0)))$  and

$$\|U_k\|_{L_s^2(H^2(\Omega_F(0)))} \leq C(\|\ell_{1,k} - U_{k,t}\|_{L_s^2(L^2(\Omega_F(0)))} + \|\dot{A}_k\|_{L_s^2} + \|\Omega_k\|_{L_s^2} + \|\hat{Q}_k - \hat{Q}\|_{L_s^2}).$$

Thus,

$$\begin{aligned} &\|A_k\|_{H_s^2} + \|\Omega_k\|_{H_s^1} + \|U_k\|_{L_s^2(H^2(\Omega_F(0)))} \\ &\leq C(\|\ell_{1,k}\|_{L_s^2(L^2(\Omega_F(0)))} + \|\ell_{2,k}\|_{L_s^2} + \|\ell_{3,k}\|_{L_s^2} + \|\ell_{4,k}\|_{L_s^2(L^2(\partial\Omega_S(0)))} + \|\hat{Q}_k - \hat{Q}\|_{L_s^2}). \end{aligned} \quad (61)$$

Observe that, for the right hand side terms, we have the following estimates:

$$\begin{aligned} \|\ell_{1,k}\|_{L_s^2(L^2(\Omega_F(0)))} &\leq C(\|R_k\|_{L_s^2(H^1(\Omega_F(0)))} + \|\hat{u}_k - \hat{u}\|_{L_s^2(H^2(\Omega_F(0)))} \\ &+ \|\hat{\rho}_k - \hat{\rho}\|_{L_s^2(H^1(\Omega_F(0)))} + \|(\nabla \hat{\chi}_k)^{-1} - (\nabla \hat{\chi})^{-1}\|_{L_s^2(L^2(\Omega_F(0)))} + \|\hat{\chi}_{k,t} - \hat{\chi}_t\|_{L_s^2(L^2(\Omega_F(0)))}), \\ \|\ell_{2,k}\|_{L_s^2} &\leq C(\|R_k\|_{L_s^2(H^1(\Omega_F(0)))} + \|\hat{u}_k - \hat{u}\|_{L_s^2(H^2(\Omega_F(0)))} + \|\hat{\rho}_k - \hat{\rho}\|_{L_s^2(H^1(\Omega_F(0)))} + \|\hat{Q}_k - \hat{Q}\|_{L_s^2}), \\ \|\ell_{3,k}\|_{L_s^2} &\leq C(\|R_k\|_{L_s^2(H^1(\Omega_F(0)))} + \|\hat{u}_k - \hat{u}\|_{L_s^2(H^2(\Omega_F(0)))} \\ &+ \|\hat{\rho}_k - \hat{\rho}\|_{L_s^2(H^1(\Omega_F(0)))} + \|\hat{Q}_k - \hat{Q}\|_{L_s^2} + \|\hat{\omega}_k - \hat{\omega}\|_{L_s^2} + \|\hat{J}_k - \hat{J}\|_{L_s^2}) \end{aligned}$$



and

$$\|\ell_{4,k}\|_{L^2_s(L^2(\partial\Omega_S(0)))} \leq C(\|\Omega_k\|_{L^2_s} + \|\hat{Q}_k - \hat{Q}\|_{H^1_s}).$$

Combining the estimates of  $\ell_{j,k}$  ( $1 \leq j \leq 4$ ) with (61), we get

$$\begin{aligned} & \|A_k\|_{H^2_s} + \|\Omega_k\|_{H^1_s} + \|U_k\|_{L^2_s(H^2(\Omega_F(0)))} \leq C(\|R_k\|_{L^2_s(H^1(\Omega_F(0)))} + \|\hat{u}_k - \hat{u}\|_{L^2_s(H^2(\Omega_F(0)))}) \\ & + \|\hat{\rho}_k - \hat{\rho}\|_{L^2_s(H^1(\Omega_F(0)))} + \|(\nabla\hat{\chi}_k)^{-1} - (\nabla\hat{\chi})^{-1}\|_{L^2_s(H^1(\Omega_F(0)))} + \|\hat{\chi}_{k,t} - \hat{\chi}_t\|_{L^2_s(L^2(\Omega_F(0)))} \\ & + \|\hat{\omega}_k - \hat{\omega}\|_{L^2_s} + \|\hat{J}_k - \hat{J}\|_{L^2_s} + \|\hat{Q}_k - \hat{Q}\|_{L^2_s} \end{aligned}$$

To conclude, we plug this inequality into (59) and use that  $\|R_k\|_{L^2_s(H^1(\Omega_F(0)))} \leq Cs^{1/2}\|R_k\|_{L^\infty_s(H^1(\Omega_F(0)))}$ :

$$\begin{aligned} & \|A_k\|_{H^2_s} + \|\Omega_k\|_{H^1_s} + \|U_k\|_{L^2_s(H^2(\Omega_F(0)))} + \|R_k\|_{L^\infty_s(H^1(\Omega_F(0)))} \leq C(\|\hat{u}_k - \hat{u}\|_{L^2_s(H^2(\Omega_F(0)))}) \\ & + \|(\nabla\hat{\chi}_k)^{-1} - (\nabla\hat{\chi})^{-1}\|_{L^2_s(H^1(\Omega_F(0)))} + \|\hat{\chi}_{k,t} - \hat{\chi}_t\|_{L^2_s(H^1(\Omega_F(0)))} + \|\hat{\rho}_k - \hat{\rho}\|_{L^2_s(H^1(\Omega_F(0)))} \quad (62) \\ & + \|\hat{\omega}_k - \hat{\omega}\|_{L^2_s} + \|\hat{J}_k - \hat{J}\|_{L^2_s} + \|\hat{Q}_k - \hat{Q}\|_{L^2_s}. \end{aligned}$$

Thanks to the estimate

$$\|((\nabla\hat{\chi}_k)^{-1} - (\nabla\hat{\chi})^{-1}, \hat{\chi}_{k,t} - \hat{\chi}_t)\|_{L^2_s(H^1(\Omega_F(0)))} + \|(\hat{J}_k - \hat{J}, \hat{Q}_k - \hat{Q})\|_{L^2_s} \leq C(\|\hat{\omega}_k - \hat{\omega}\|_{L^2_s} + \|\hat{a}_k - \hat{a}\|_{H^1_s}),$$

we find that  $(R_k, U_k, A_k, \Omega_k) \rightarrow 0$  for the norm (58).

### 3.5 Uniqueness

Let us consider two solutions  $(\tilde{\rho}_1, \tilde{u}_1, a_1, \omega_1)$  and  $(\tilde{\rho}_2, \tilde{u}_2, a_2, \omega_2)$  of system (36) such that

$$\tilde{N}_{0,s}(\tilde{\rho}_j, \tilde{u}_j, a_j, \omega_j) \leq \tilde{R}, \quad j = 1, 2.$$

Then, we look at the system fulfilled by the difference, we multiply the equation of  $\tilde{\rho}_1 - \tilde{\rho}_2$  by  $\tilde{\rho}_1 - \tilde{\rho}_2$ , the equation of  $\tilde{u}_1 - \tilde{u}_2$  by  $\tilde{u}_1 - \tilde{u}_2$ , the equation of  $a_1 - a_2$  by  $\dot{a}_1 - \dot{a}_2$  and the equation of  $\omega_1 - \omega_2$  by  $\omega_1 - \omega_2$ .

Integrating by parts in the second order term of the equation of  $\tilde{u}_1 - \tilde{u}_2$  and using classical arguments, we obtain

$$\begin{aligned} & \frac{d}{dt} \left( \int_{\Omega_F(0)} (|\tilde{\rho}_1 - \tilde{\rho}_2|^2 + |\tilde{u}_1 - \tilde{u}_2|^2) dx + |\dot{a}_1 - \dot{a}_2|^2 + |\omega_1 - \omega_2|^2 \right) + \int_{\Omega_F(0)} |\nabla\tilde{u}_1 - \nabla\tilde{u}_2|^2 dx \\ & \leq C \left( \int_{\Omega_F(0)} (|\tilde{\rho}_1 - \tilde{\rho}_2|^2 + |\tilde{u}_1 - \tilde{u}_2|^2) dx + |\dot{a}_1 - \dot{a}_2|^2 + |\omega_1 - \omega_2|^2 \right). \end{aligned}$$

Finally, using Gronwall's Lemma we deduce that  $\tilde{\rho}_1 = \tilde{\rho}_2$ ,  $\tilde{u}_1 = \tilde{u}_2$ ,  $a_1 = a_2$  and  $\omega_1 = \omega_2$ .

## 4 A priori estimates: proof of Proposition 8

The proof of this result is inspired by the works [17] and [18], where the authors dealt with the compressible Navier-Stokes equations.

For the sake of simplicity, in this paragraph we will denote our solution by  $(\rho, u, a, \omega)$  instead of

$(\rho^*, u, a, \omega)$ . The equations are

$$\left\{ \begin{array}{ll} \rho_t + (u \cdot \nabla)\rho + \bar{\rho}(\nabla \cdot u) = f_0 & \text{in } Z_T, \\ u_t - \nabla \cdot \sigma(u, \rho) = f_1 & \text{in } Z_T, \\ m\ddot{a} = \int_{\partial\Omega_S(t)} \sigma(u, \rho)n \, d\gamma + f_2 & \text{in } (0, T), \\ J\dot{\omega} = \int_{\partial\Omega_S(t)} (x - a) \wedge (\sigma(u, \rho)n) \, d\gamma + f_3 & \text{in } (0, T), \\ u = \dot{a} + \omega \wedge (x - a) & \text{in } \Sigma_T, \\ u = 0 & \text{on } (0, T) \times \partial\Omega, \\ \rho(0, \cdot) = \rho_0 - \bar{\rho}, u(0, \cdot) = u_0 & \text{in } \Omega_F(0), \\ a(0) = 0, \dot{a}(0) = a_0, \omega(0) = \omega_0. & \end{array} \right. \quad (63)$$

where  $f_0, f_1, f_2$  and  $f_3$  are given by (28) (with  $\rho^*$  replaced by  $\rho$ ). Recall that the definitions of  $Z_T$  and  $\Sigma_T$  were given in (23).

In the following lines, we will give several lemmas where we will present a priori estimates of different nature:

- An estimate of the solid motion in terms of the fluid velocity (Lemma 10).
- Global estimates associated to an elliptic operator (Lemma 11) and energy-type estimates associated to the compressible Stokes system (Lemmas 12 and 13).
- Interior estimates for the compressible Stokes system (Lemma 14).
- Estimates close to the boundary  $\partial\Omega_F(0)$ . First, we will estimate the tangential derivatives (Lemma 15) and then the normal ones (Lemma 16).
- Global estimates for the Stokes operator (Lemma 17).

In the above lines, the term ‘Global’ refers to estimates on the whole domain  $\Omega_F(t)$ .

All this will be proved under the hypothesis that  $(\rho, u, a, \omega) \in X(0, T)$ . The conclusion of all these Lemmas will be the following inequality:

$$\begin{aligned} N_{0,T}(\rho, u, a, \omega) &\leq C \left( \|f_0\|_{H_T^1(L^2)} + \|f_0\|_{L_T^\infty(H^2)} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} + \|f_1\|_{H_T^1(L^2)} \right. \\ &\quad + \|f_1\|_{L_T^\infty(H^1)} + \|f_2\|_{H_T^1} + \|f_3\|_{H_T^1} + \|\partial_t u(0, \cdot)\|_{H^1} + \|u_0\|_{H^3} + \|\partial_t \rho(0, \cdot)\|_{L^2} \\ &\quad \left. + \|\rho_0\|_{H^3} + |\ddot{a}(0)| + |a_0| + |\dot{\omega}(0)| + |\omega_0| + N_{0,T}^{3/2}(\rho, u, a, \omega) + N_{0,T}^2(\rho, u, a, \omega) \right) \end{aligned} \quad (64)$$

(recall that  $N_{0,T}(\rho, u, a, \omega)$  was defined in (30)). Next, using the equations of  $u, \rho, a$  and  $\omega$ , we see that

$$\begin{aligned} \|\partial_t u(0, \cdot)\|_{H^1} &\leq C(\|u_0\|_{H^3} + \|\rho_0\|_{H^2} + \|f_1\|_{L_T^\infty(H^1)}), \\ \|\partial_t \rho(0, \cdot)\|_{L^2} &\leq C(\|u_0\|_{H^1} + \|f_0\|_{L_T^\infty(L^2)} + N_{0,T}^2(\rho, u, a, \omega)), \\ |\ddot{a}(0)| &\leq C(\|u_0\|_{H^2} + \|\rho_0\|_{H^1} + \|f_2\|_{L_T^\infty}) \end{aligned}$$

and

$$|\dot{\omega}(0)| \leq C(\|u_0\|_{H^2} + \|\rho_0\|_{H^1} + \|f_3\|_{L_T^\infty}).$$

Then, from the expression of  $f_i$  ( $0 \leq i \leq 3$ ), we have

$$\begin{aligned} &\|f_0\|_{H_T^1(L^2)} + \|f_0\|_{L_T^\infty(H^2)} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} + \|f_1\|_{H_T^1(L^2)} + \|f_1\|_{L_T^\infty(H^1)} \\ &\quad + \|f_2\|_{H_T^1} + \|f_3\|_{H_T^1} \leq N_{0,T}^2(\rho, u, a, \omega). \end{aligned}$$

Finally, using the assumption  $N_{0,T}(\rho, u, a, \omega) \leq \delta_1$  with  $\delta_1$  small enough, we conclude the proof of Proposition 8. Consequently, from now on we concentrate in the proof of inequality (64).

## 4.1 Technical results

All along the proof of these Lemmas and for the sake of simplicity we will adopt the notation  $N(0, T)$  instead of  $N_{0,T}(\rho, u, a, \omega)$ .

Furthermore,  $C$  will stand for generic positive constants which do not depend on  $t$  but which may depend on  $\Omega_S(0)$ ,  $\Omega_F(0)$  and  $\Omega$ .

The first result is the following one:

**Lemma 10** *There exists a positive constant  $C > 0$  such that*

$$\|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} \leq C \|u\|_{L_T^2(H^1)}. \quad (65)$$

**Proof:** The value of the velocity vector field on the boundary of the solid is given by (see (63))

$$u(t, x) = \dot{a}(t) + \omega(t) \wedge (x - a(t)) = u_S(t, x), \quad t \in (0, T), \quad x \in \partial\Omega_S(t). \quad (66)$$

Taking the scalar product of this by  $(x - a(t))$  and taking the square, we have

$$|u(t, x) \cdot (x - a(t))|^2 = |\dot{a}(t)|^2 |x - a(t)|^2 \cos^2(\dot{a}(t), x - a(t)), \quad t \in (0, T), \quad x \in \partial\Omega_S(t).$$

Then,

$$\int_{\partial\Omega_S(t)} |\dot{a}(t)|^2 |x - a(t)|^2 |\cos(\dot{a}(t), x - a(t))|^2 dx = |\dot{a}(t)|^2 \int_{\partial\Omega_S(0)} |y|^2 |\cos(\dot{a}(t), Q(t)y)|^2 dy. \quad (67)$$

First, we observe that  $|y| \geq \text{dist}(0, \overline{\partial\Omega_S(0)})$ . Thanks to the fact that  $\partial\Omega_S(0)$  is a regular closed curve, we have that for any  $b \in \mathbb{R}^3$

$$\int_{\partial\Omega_S(0)} |\cos(b, y)|^2 \geq C > 0.$$

Using this in (67), we have (65) for  $\|\dot{a}\|_{L_T^2}$ . Finally, from (66) and similar arguments, we also deduce this inequality for  $\|\omega\|_{L_T^2}$ .

The next result concerns a classical elliptic estimate for  $u$ :

**Lemma 11** *Let  $k = 2, 3$ . Then, we have*

$$\|u\|_{L_T^\infty(H^k)} \leq C (\|u\|_{W_T^{1,\infty}(H^{k-2})} + \|\rho\|_{L_T^\infty(H^{k-1})} + \|\dot{a}\|_{L_T^\infty} + \|\omega\|_{L_T^\infty} + \|f_1\|_{L_T^\infty(H^{k-2})}). \quad (68)$$

**Lemma 12** *Let  $k = 0, 1$ . Then, for every  $\delta > 0$ , there exists a constant  $C > 0$  such that*

$$\begin{aligned} & \|u\|_{W_T^{k,\infty}(L^2)} + \|\rho\|_{W_T^{k,\infty}(L^2)} + \|u\|_{H_T^k(H^1)} + \|\dot{a}\|_{W_T^{k,\infty}} + \|\omega\|_{W_T^{k,\infty}} \leq \delta (\|\rho\|_{H_T^k(L^2)} + \|\dot{a}\|_{H_T^k} + \|\omega\|_{H_T^k}) \\ & + C (\|f_0\|_{H_T^k(L^2)} + \|f_1\|_{H_T^k(L^2)} + \|f_2\|_{H_T^k} + \|f_3\|_{H_T^k} + \|u_0\|_{L^2} + \|\rho_0\|_{L^2} + |a_0| + |\omega_0| \\ & + \|\partial_t^k u(0, \cdot)\|_{L^2} + \|\partial_t^k \rho(0, \cdot)\|_{L^2} + |\partial_t^{k+1} a(0)| + |\partial_t^k \omega(0)| + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (69)$$

**Proof:**

**1)**  $k = 0$ . We multiply the equation of  $\rho$  by  $p^0 \rho / \bar{p}$  and the equation of  $u$  by  $u$ . Integrating in  $Z_s$  for  $s \in (0, T)$  and adding up both expressions, we obtain:

$$\begin{aligned} & \frac{1}{2} \sup_{s \in (0, T)} \left( |\dot{a}|^2(s) + |\omega|^2(s) + \int_{\Omega_F(s)} \left( \frac{p^0}{\bar{p}} |\rho|^2 + |u|^2 \right) (s) dx \right) + \iint_{Z_T} (2\mu |\epsilon(u)|^2 + \mu' |\nabla \cdot u|^2) dx dt \\ & \leq \delta (\|u\|_{L_T^2(L^2)} + \|\rho\|_{L_T^2(L^2)}^2 + \|\dot{a}\|_{L_T^2}^2 + \|\omega\|_{L_T^2}^2) + C \left( \iint_{Z_T} (|\nabla \cdot u| |\rho|^2 + |\nabla \cdot (|u|^2 u)|) dx dt \right. \\ & \left. + \iint_{Z_T} (|f_0|^2 + |f_1|^2) dx dt + \int_{\Omega_F(0)} (|u_0|^2 + |\rho_0|^2) dx + |a_0|^2 + |\omega_0|^2 + \int_0^T (|f_2|^2 + |f_3|^2) dt \right). \end{aligned} \quad (70)$$

Here, we have put together the integrals coming from the third term in the equation of  $\rho$  and the third term in the expression of  $\sigma(u, \rho)$  (see (27)) and we have integrated by parts. We have also used the fact that according to the definition (9) of  $J$ , we have  $(J\omega) \cdot \omega = 0$ .

Thanks to (46), we get the  $L_T^2(H^1)$  norm of  $u$  in the left hand side of (70). Moreover, we can estimate the nonlinear term in (70) in the following way:

$$\iint_{Z_T} (|\nabla \cdot u||\rho|^2 + |\nabla \cdot (|u|^2 u)|) dx dt \leq C(\|u\|_{L_T^\infty(W^{1,\infty})})(\|u\|_{L_T^2(L^2)}^2 + \|\rho\|_{L_T^2(L^2)}^2) \leq CN^3(0, T).$$

The definition of  $N(0, T)$  was given in (30). Thus, we deduce (69) for  $k = 0$

**2)  $k = 1$ .** First, we differentiate (63) with respect to  $t$  and we find:

$$\left\{ \begin{array}{ll} (\rho_t)_t + (u \cdot \nabla)\rho_t + \bar{p}(\nabla \cdot u_t) = G_0 & \text{in } Z_T, \\ (u_t)_t - \nabla \cdot \sigma(u_t, \rho_t) = G_1 & \text{in } Z_T, \\ m\ddot{a} = \int_{\partial\Omega_S(t)} \sigma(u_t, \rho_t)n d\gamma + G_2 & \text{in } (0, T), \\ J\dot{\omega} = G_3 + \int_{\partial\Omega_S(t)} (x-a) \wedge (\sigma(u_t, \rho_t))n d\gamma & \text{in } (0, T), \\ u_t = \dot{a} + \dot{\omega} \wedge (x-a) + G_4 & \text{in } \Sigma_T, \\ u_t = 0 & \text{in } (0, T) \times \partial\Omega, \end{array} \right. \quad (71)$$

with

$$\begin{aligned} G_0 &= f_{0,t} - (u_t \cdot \nabla)\rho, & G_1 &= f_{1,t}, & G_2 &= f_{2,t} + \int_{\partial\Omega_S(t)} u_S \cdot \nabla(\sigma(u, p)n) d\gamma, \\ G_3 &= f_{3,t} + \int_{\partial\Omega_S(t)} u_S \cdot \nabla((x-a) \wedge (\sigma(u, p)n)) d\gamma - J\dot{\omega} - \int_{\partial\Omega_S(t)} \dot{a} \wedge (\sigma(u, \rho))n d\gamma, \\ G_4 &= \omega \wedge (\omega \wedge (x-a)) - (u_S \cdot \nabla)u. \end{aligned}$$

Recall that

$$u_S(t, x) = \dot{a}(t) + \omega(t) \wedge (x - a(t))$$

is the velocity of the solid.

Now, we multiply the equation of  $\rho_t$  by  $(p^0/\bar{p})\rho_t$  and the equation of  $u_t$  by  $u_t$ . Let us see that the boundary terms provide estimates for  $\dot{\omega}$  and  $\ddot{a}$  and the remaining terms are bounded by  $C(N^3(0, T) + N^4(0, T))$  and the data:

$$\begin{aligned} \iint_{\Sigma_s} (\sigma(u_t, \rho_t)n) \cdot u_t d\gamma dt &= \frac{m}{2}(|\ddot{a}(t)|^2 - |\ddot{a}(0)|^2) - \int_0^s \ddot{a}G_2 dt \\ &+ \iint_{\Sigma_s} \dot{\omega} \cdot ((x-a) \wedge (\sigma(u_t, \rho_t)n)) d\gamma dt + \iint_{\Sigma_s} G_4 \cdot (\sigma(u_t, \rho_t)n) d\gamma dt. \end{aligned} \quad (72)$$

In order to develop the third term in the right hand side of this identity, we use that

$$\dot{\omega} \cdot \int_{\partial\Omega_S(t)} (x-a) \wedge (\sigma(u_t, \rho_t)n) d\gamma = (J\dot{\omega}) \cdot \dot{\omega} - G_3 \cdot \dot{\omega}. \quad (73)$$

Observe that we have

$$(J\dot{\omega}) \cdot \dot{\omega} = \frac{1}{2} \frac{d}{dt} ((J\dot{\omega}) \cdot \dot{\omega}) - \frac{1}{2} (J\dot{\omega}) \cdot \dot{\omega}.$$

Now, we integrate between  $t = 0$  and  $t = s$  in (73) and we use that  $J$  is a definite positive matrix (see (10)). We find

$$\begin{aligned} &\iint_{\Sigma_s} \dot{\omega} \cdot ((x-a) \wedge (\sigma(u_t, \rho_t)n)) d\gamma dt \\ &\geq (C_J/2)|\dot{\omega}(s)|^2 - C|\dot{\omega}(0)|^2 - \delta\|\dot{\omega}\|_{L_T^2}^2 - C \int_0^T |G_3|^2 d\tau - CN^3(0, T), \end{aligned} \quad (74)$$

for  $\delta > 0$  small enough. Thanks to the estimate

$$\|G_4\|_{L^2(\Sigma_T)} \leq CN^2(0, T)$$

we have that the last term in identity (72) is bounded by  $N^3(0, T)$ .

For the term in  $G_0$  we act as follows:

$$\iint_{Z_T} G_0 \rho_t dx dt \leq C(\|f_{0,t}\|_{L_T^2(L^2)}^2 + N^4(0, T)) + \delta \|\rho_t\|_{L_T^2(L^2)}^2,$$

for  $\delta > 0$  small enough. The same can be done for the term in  $G_1$ .

Similarly as before, we find

$$\begin{aligned} & \frac{1}{2} \sup_{s \in (0, T)} \left( |\ddot{a}|^2(s) + |\dot{\omega}|^2(s) + \int_{\Omega_F(s)} \left( \frac{p^0}{\bar{\rho}} |\rho_t|^2 + |u_t|^2 \right) (s) dx \right) + \iint_{Z_T} (2\mu |\epsilon(u_t)|^2 + \mu' |\nabla \cdot u_t|^2) dx dt \\ & \leq \delta \left( \|u_t\|_{L_T^2(L^2)}^2 + \|\rho_t\|_{L_T^2(L^2)}^2 + \|\ddot{a}\|_{L_T^2}^2 + \|\dot{\omega}\|_{L_T^2}^2 \right) + C \left( \iint_{Z_T} (|\nabla \cdot u| |\rho_t|^2 + |\nabla \cdot (|u_t|^2 u)|) dx dt \right. \\ & \quad + \int_{\Omega_F(0)} (|\partial_t u(0, \cdot)|^2 + |\partial_t \rho(0, \cdot)|^2) dx + |\ddot{a}(0)|^2 + |\dot{\omega}(0)|^2 + \|f_{0,t}\|_{L_T^2(L^2)}^2 + \|f_{1,t}\|_{L_T^2(L^2)}^2 \\ & \quad \left. + \int_0^T |f_{2,t}|^2 dt + \int_0^T |f_{3,t}|^2 dt + N^3(0, T) + N^4(0, T) \right), \end{aligned}$$

for  $\delta > 0$  small enough. In order to prove this last inequality, we have used that

$$\|G_2 - f_{2,t}\|_{L_T^2} + \|G_3 - f_{3,t}\|_{L_T^2} \leq CN^2(0, T).$$

Thus, we deduce (69) for  $k = 1$ .

**Lemma 13** *Let  $k = 0, 1$ . Then, there exists a constant  $C > 0$  such that*

$$\begin{aligned} & \|u_t\|_{H_T^k(L^2)} + \|\rho_t\|_{H_T^k(L^2)} + \|u\|_{W_T^{k,\infty}(H^1)} + \|\ddot{a}\|_{H_T^k} + \|\dot{\omega}\|_{H_T^k} \\ & \leq C(\|u\|_{H_T^k(H^1)} + \|\rho\|_{W_T^{k,\infty}(L^2)} + \|f_0\|_{H_T^k(L^2)} + \|f_1\|_{H_T^k(L^2)} + \|f_2\|_{H_T^k} + \|f_3\|_{H_T^k} \\ & \quad + \|u_0\|_{H^1} + \|\partial_t^k u(0, \cdot)\|_{H^1} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (75)$$

**Proof:**

**1)**  $k = 0$ . We multiply the equation of  $\rho$  by  $\rho_t$  and the equation of  $u$  by  $u_t$ . Integrating in  $Z_s$  for  $s \in (0, T)$  and adding up both expressions, we obtain:

$$\begin{aligned} & \frac{1}{2} \int_{\Omega_F(s)} (2\mu |\epsilon(u)(s)|^2 + \mu' |\nabla \cdot u(s)|^2) dx + \iint_{Z_s} (|\rho_t|^2 + |u_t|^2) dx dt \\ & + \iint_{\Sigma_s} u_t \cdot (\sigma(u, \rho)n) d\gamma dt + \iint_{Z_s} (\bar{\rho}(\nabla \cdot u)\rho_t - p^0(\nabla \cdot u_t)\rho) dx dt \\ & \leq C \left( \iint_{Z_s} |\nabla \cdot ((2\mu |\epsilon(u)|^2 + \mu' |\nabla \cdot u|^2)u)| dx dt + \int_{\Omega_F(0)} |\nabla u_0|^2 dx \right. \\ & \quad \left. + \iint_{Z_s} (|f_0|^2 + |f_1|^2 + \delta(|\rho_t|^2 + |u_t|^2)) dx dt + N^3(0, T) \right), \end{aligned} \quad (76)$$

for  $\delta > 0$  small enough. Using the fifth equation in (71), the boundary term in (76) yields

$$\begin{aligned} & \iint_{\Sigma_s} u_t \cdot (\sigma(u, \rho)n) d\gamma dt = m \int_0^s |\ddot{a}|^2 dt - \int_0^s \ddot{a} \cdot f_2 dt \\ & + \int_0^s ((J\dot{\omega}) \cdot \dot{\omega} - f_3 \cdot \dot{\omega}) dt + \iint_{\Sigma_s} G_4 \cdot (\sigma(u, \rho)n) d\gamma dt \\ & \geq \int_0^s ((m/2)|\ddot{a}|^2 + (C_J/2)|\dot{\omega}|^2) dt - C \left( \int_0^s (|f_2|^2 + |f_3|^2) dt + N^3(0, T) \right). \end{aligned}$$

The last integral in the left hand side of (76) can be estimated as follows:

$$\begin{aligned} & \iint_{Z_s} (\bar{\rho}(\nabla \cdot u)\rho_t - p^0(\nabla \cdot u_t)\rho) dx dt \leq \frac{1}{2} \iint_{Z_s} |\rho_t|^2 dx dt + C \iint_{Z_s} |\nabla u|^2 dx dt \\ & + \iint_{Z_s} p^0(\nabla \cdot u)\rho_t dx dt - p^0 \int_{\Omega_F(t)} ((\nabla \cdot u)\rho)(t)|_{t=0}^{t=s} dx + \iint_{Z_s} p^0 \nabla \cdot (\rho(\nabla \cdot u)u) dx dt. \end{aligned} \quad (77)$$

Taking the supremum in  $s$  in (76), one can readily deduce (75) for  $k = 0$ .

**2)  $k = 1$ .** As in the proof of Lemma 12, we consider system (71). Let us multiply the equation of  $\rho_t$  by  $\rho_{tt}$  and the equation of  $u_t$  by  $u_{tt}$ . This provides:

$$\begin{aligned} & \frac{1}{2} \int_{\Omega_F(t)} (2\mu|\epsilon(u_t)(s)|^2 + \mu'|\nabla \cdot u_t(s)|^2) dx + \iint_{Z_s} (|\rho_{tt}|^2 + |u_{tt}|^2) dx dt \\ & + \iint_{\Sigma_s} u_{tt} \cdot (\sigma(u_t, \rho_t)n) d\gamma dt + \iint_{Z_s} (\bar{\rho}(\nabla \cdot u_t)\rho_{tt} - p^0(\nabla \cdot u_{tt})\rho_t) dx dt \\ & \leq C \left( \iint_{Z_s} |\nabla \cdot ((2\mu|\epsilon(u_t)|^2 + \mu'|\nabla \cdot u_t|^2)u)| dx dt + \iint_{Z_s} |\rho_{tt}||u||\nabla \rho_t| dx dt \right. \\ & \left. + \int_{\Omega_F(0)} |\partial_t \nabla u(0, \cdot)|^2 dx + \iint_{Z_s} (|G_0|^2 + |G_1|^2 + \delta(|\rho_{tt}|^2 + |u_{tt}|^2)) dx dt \right), \end{aligned} \quad (78)$$

for  $\delta > 0$  small enough. Let us now compute the boundary terms. On  $\Sigma_T$ , we have:

$$u_{tt} = \ddot{a} + \ddot{\omega} \wedge (x - a) + G_5,$$

with

$$G_5 = 2\dot{\omega} \wedge (\omega \wedge (x - a)) + \omega \wedge [(\dot{\omega} \wedge (x - a)) + \omega \wedge (\omega \wedge (x - a))] - (u_{S,t} \cdot \nabla)u - (u_S \cdot \nabla)[2u_t + (u_S \cdot \nabla)u].$$

Thus,

$$\begin{aligned} & \iint_{\Sigma_s} u_{tt} \cdot (\sigma(u_t, \rho_t)n) d\gamma dt = m \int_0^s |\ddot{a}|^2 dt - \int_0^s \ddot{a} \cdot G_2 dt \\ & + \iint_{\Sigma_s} \ddot{\omega} \cdot ((x - a) \wedge (\sigma(u_t, \rho_t)n)) d\gamma dt + \iint_{\Sigma_s} G_5 \cdot (\sigma(u_t, \rho_t)n) d\gamma dt. \end{aligned} \quad (79)$$

In order to develop the third term in the right hand side of (79) we use that

$$\ddot{\omega} \cdot \int_{\partial\Omega_S(t)} (x - a) \wedge (\sigma(u_t, \rho_t)n) d\gamma = (J\ddot{\omega}) \cdot \ddot{\omega} - G_3 \cdot \ddot{\omega}. \quad (80)$$

Now, we integrate between  $t = 0$  and  $t = s$  in (80); we find

$$\iint_{\Sigma_s} \ddot{\omega} \cdot ((x - a) \wedge (\sigma(u_t, \rho_t)n)) d\gamma dt \geq (C_J/2) \int_0^s |\ddot{\omega}|^2 dt - C \left( \int_0^s |f_{3,t}|^2 dt + N^4(0, T) \right), \quad (81)$$

where  $C_J$  is such that (10) is satisfied. The last term in the right hand side of (79) is also bounded by  $N^3(0, T) + N^4(0, T)$ .

Finally, the last term in the left hand side of (78) is estimated as the corresponding term in the previous case (see (77)). Taking the supremum in  $s$  in (78), we conclude inequality (75) for  $k = 1$ .

### Interior Estimates

Before giving the results, we introduce the total time derivative of  $\rho$ :

$$\frac{d\rho}{dt} := \rho_t + (u \cdot \nabla)\rho \quad \text{in } Z_T. \quad (82)$$

Let  $\zeta_{0,0} \in C_0^\infty(\Omega_F(0))$  and

$$\zeta_0(t, x) := \zeta_{0,0}(\psi(t, x)), \quad \forall x \in \Omega_F(t), \quad t \in (0, T),$$

where  $\psi$  is the flow (extending  $\chi_S^{-1}$ ) defined in Corollary 2.

**Lemma 14** *For  $1 \leq k \leq 3$ , we have*

$$\begin{aligned} & \|\zeta_0 \nabla \rho\|_{L_T^\infty(H^{k-1})} + \|\zeta_0 \nabla \rho\|_{L_T^2(H^{k-1})} + \|\zeta_0 \nabla u\|_{L_T^\infty(H^{k-1})} + \|\zeta_0 \nabla u\|_{L_T^2(H^k)} + \|\zeta_0 \nabla \frac{d\rho}{dt}\|_{L_T^2(H^{k-1})} \\ & \leq C(\|u\|_{L_T^\infty(H^{k-1})} + \|u\|_{L_T^2(H^k)} + \|\rho_0\|_{H^k} + \|u_0\|_{H^k} + \|f_0\|_{L_T^2(H^k)} + \|f_1\|_{L_T^2(H^{k-1})} + N^{3/2}(0, T)). \end{aligned} \quad (83)$$

**Proof:** In this proof, we denote by  $D^\ell$  all possible derivatives in  $x$  of order  $\ell$ . We divide the proof in two steps:

• First, we apply the operator  $\zeta_0^2 \partial_{x_j} D^{k-1}$  to the equation of  $\rho$ , we multiply it by  $\partial_{x_j} D^{k-1} \rho$ , we sum up in  $j = 1, 2, 3$  and we integrate in  $Z_s$ :

$$\sum_{j=1}^3 \iint_{Z_s} \zeta_0^2 \partial_{x_j} D^{k-1} (\rho_t + (u \cdot \nabla) \rho + \bar{\rho}(\nabla \cdot u) - f_0) \partial_{x_j} D^{k-1} \rho \, dx \, dt = 0. \quad (84)$$

The first term of this expression coincides with

$$\sum_{j=1}^3 \left( \frac{1}{2} \int_0^s \frac{d}{dt} \int_{\Omega_F(t)} \zeta_0^2 |\partial_{x_j} D^{k-1} \rho|^2 \, dx \, dt - \iint_{Z_s} \zeta_0 (\nabla \zeta_{0,0}(\psi) \cdot \psi_t) |\partial_{x_j} D^{k-1} \rho|^2 \, dx \, dt \right). \quad (85)$$

Here, we have used the fact that  $\zeta_0$  is compactly supported, so

$$\iint_{Z_s} \nabla \cdot (\zeta_0^2 |\partial_{x_j} D^{k-1} \rho|^2 u) \, dx \, dt = 0. \quad (86)$$

Thanks to (22), we have that the second integral in (85) is estimated by  $CN^3(0, T)$ .

As long as the nonlinear term in (84) is concerned, using (86) we have

$$\begin{aligned} & \sum_{j=1}^3 \iint_{Z_s} \zeta_0^2 \partial_{x_j} D^{k-1} [(u \cdot \nabla) \rho] \partial_{x_j} D^{k-1} \rho \, dx \, dt \\ & \leq CN^3(0, T) + \sum_{j=1}^3 \iint_{Z_s} \zeta_0^2 [(u \cdot \nabla) \partial_{x_j} D^{k-1} \rho] \partial_{x_j} D^{k-1} \rho \, dx \, dt \\ & = CN^3(0, T) - \frac{1}{2} \sum_{j=1}^3 \iint_{Z_s} \nabla \cdot (\zeta_0^2 u) |\partial_{x_j} D^{k-1} \rho|^2 \, dx \, dt \leq CN^3(0, T). \end{aligned} \quad (87)$$

Putting all this together with (84), we find

$$\begin{aligned} & \|\zeta_0 \nabla \rho\|_{L_T^\infty(H^{k-1})}^2 + \bar{\rho} \iint_{Z_s} \zeta_0^2 \nabla(D^{k-1}(\nabla \cdot u)) \cdot \nabla(D^{k-1} \rho) \, dx \, dt \\ & \leq \delta \|\zeta_0 D^k \rho\|_{L_T^2(L^2)}^2 + C(\|\rho_0\|_{H^k}^2 + \|f_0\|_{L_T^2(H^k)}^2 + N^3(0, T)), \end{aligned} \quad (88)$$

with  $\delta > 0$  small enough.

Next, we apply the operator  $\frac{\bar{\rho}}{2\mu + \mu'} \zeta_0^2 D^{k-1}$  to the  $j$ -th equation of  $u$ , we multiply it by  $\partial_{x_j} D^{k-1} \rho$ , we sum up in  $j = 1, 2, 3$  and we integrate in  $Z_s$ :

$$\frac{\bar{\rho}}{2\mu + \mu'} \sum_{j=1}^3 \iint_{Z_s} \zeta_0^2 D^{k-1} (u_t - \nabla \cdot \sigma(u, \rho) - f_1)_j \partial_{x_j} D^{k-1} \rho \, dx \, dt = 0. \quad (89)$$

We integrate by parts in  $t$  in the term concerning  $u_t$ :

$$\begin{aligned} \iint_{Z_s} \zeta_0^2 D^{k-1} u_{j,t} \partial_{x_j} D^{k-1} \rho \, dx \, dt &= - \iint_{Z_s} \nabla \cdot ((\zeta_0^2 D^{k-1} u_j \partial_{x_j} D^{k-1} \rho) u) \, dx \, dt \\ &- 2 \iint_{Z_s} \zeta_0 (\nabla \zeta_{0,0}(\psi) \cdot \psi_t) D^{k-1} u_j \partial_{x_j} D^{k-1} \rho \, dx \, dt - \iint_{Z_s} \zeta_0^2 D^{k-1} u_j \partial_{x_j} D^{k-1} \rho_t \, dx \, dt. \end{aligned}$$

The first term in the right hand side is zero thanks to (86) while the second one can be estimated by  $CN^3(0, T)$  due to (22). As long as the third term is concerned, we use the equation of  $\rho$  in (63); arguing as in (87), we obtain

$$\iint_{Z_s} \zeta_0^2 D^{k-1} u_{j,t} \partial_{x_j} D^{k-1} \rho \, dx \, dt \leq \delta \|\zeta_0 \rho\|_{L_T^\infty(H^k)}^2 + C(\|f_0\|_{L_T^2(H^k)}^2 + \|u\|_{L_T^2(H^k)}^2 + \|u\|_{L_T^\infty(H^{k-1})}^2 + N^3(0, T)).$$

As long as the elliptic term  $-2\mu \nabla \cdot (\epsilon(u)) - \mu' \nabla \cdot ((\nabla \cdot u) Id)$  is concerned, we rewrite it as

$$-2\mu \nabla \cdot (\epsilon(u)) - \mu' \nabla \cdot ((\nabla \cdot u) Id) = -\mu \Delta u - (\mu + \mu') \nabla (\nabla \cdot u)$$

and we integrate by parts twice (with respect to  $x$ ) in the laplacian term. Observe that there is no boundary term when we integrate in  $x$  since  $\zeta_0$  has a compact support. Then, from (89) we deduce that

$$\begin{aligned} \frac{\bar{p} p^0}{2\mu + \mu'} \|\zeta_0 D^k \rho\|_{L_T^2(L^2)}^2 - \bar{p} \iint_{Z_s} \zeta_0^2 \nabla (D^{k-1} (\nabla \cdot u)) \cdot \nabla (D^{k-1} \rho) \, dx \, dt &\leq \delta (\|\zeta_0 D^k \rho\|_{L_T^2(L^2)}^2 + \|\zeta_0 \rho\|_{L_T^\infty(H^k)}^2) \\ &+ C(\|u\|_{L_T^\infty(H^{k-1})}^2 + \|u\|_{L_T^2(H^k)}^2 + \|\rho_0\|_{H^k}^2 + \|f_0\|_{L_T^2(H^k)}^2 + \|f_1\|_{L_T^2(H^{k-1})}^2 + N^3(0, T)). \end{aligned}$$

This, together with (88) and taking the supremum in  $s$ , yields

$$\begin{aligned} \|\zeta_0 \nabla \rho\|_{L_T^\infty(H^{k-1})} + \|\zeta_0 \nabla \rho\|_{L_T^2(H^{k-1})} &\leq C(\|u\|_{L_T^\infty(H^{k-1})} + \|u\|_{L_T^2(H^k)} \\ &+ \|\rho_0\|_{H^k} + \|f_0\|_{L_T^2(H^k)} + \|f_1\|_{L_T^2(H^{k-1})} + N^{3/2}(0, T)). \end{aligned} \tag{90}$$

• Let us now apply the operator  $(p^0/\bar{p})\zeta_0^2 D^k$  to the equation of  $\rho$ , we multiply it by  $D^k \rho$  and we integrate in  $Z_s$ :

$$\frac{p^0}{\bar{p}} \iint_{Z_s} \zeta_0^2 D^k (\rho_t + (u \cdot \nabla) \rho + \bar{p} (\nabla \cdot u) - f_0) D^k \rho \, dx \, dt = 0.$$

Similarly as before (using (22) and (86)), this gives

$$\begin{aligned} \|\zeta_0 \nabla \rho\|_{L_s^\infty(H^{k-1})}^2 + p^0 \iint_{Z_s} \zeta_0^2 D^k (\nabla \cdot u) D^k \rho \, dx \, dt \\ \leq \delta \|\zeta_0 D^k \rho\|_{L_T^2(L^2)}^2 + C(\|\rho_0\|_{H^k}^2 + \|f_0\|_{L_T^2(H^k)}^2 + N^3(0, T)). \end{aligned} \tag{91}$$

Next, we apply the operator  $\zeta_0^2 D^k$  to the  $j$ -th equation of  $u$ , we multiply it by  $D^k u_j$ , we sum up in  $j = 1, 2, 3$  and we integrate in  $Z_s$ :

$$\sum_{j=1}^3 \iint_{Z_s} \zeta_0^2 D^k (u_t - \nabla \cdot \sigma(u, \rho) - f_1)_j D^k u_j \, dx \, dt = 0.$$

We integrate by parts in  $x$  in the second and third terms of the last integral. We deduce that

$$\begin{aligned} \|\zeta_0 u\|_{L_T^\infty(H^k)}^2 + \|\zeta_0 u\|_{L_T^2(H^{k+1})}^2 - p^0 \iint_{Z_s} \zeta_0^2 D^k (\nabla \cdot u) D^k \rho \, dx \, dt \\ \leq \delta \|\zeta_0 \rho\|_{L_T^2(H^k)}^2 + C(\|u\|_{L_T^2(H^k)}^2 + \|u_0\|_{H^k}^2 + \|f_1\|_{L_T^2(H^{k-1})}^2 + N^3(0, T)), \end{aligned}$$



for  $\delta > 0$  small enough. This, together with (91), yields

$$\begin{aligned} & \|\zeta_0 \rho\|_{L_T^\infty(H^k)} + \|\zeta_0 u\|_{L_T^\infty(H^k)} + \|\zeta_0 u\|_{L_T^2(H^{k+1})} \\ & \leq \delta \|\zeta_0 \rho\|_{L_T^2(H^k)} + C(\|u\|_{L_T^2(H^k)} + \|u_0\|_{H^k} + \|\rho_0\|_{H^k} + \|f_0\|_{L_T^2(H^k)} + \|f_1\|_{L_T^2(H^{k-1})} + N^{3/2}(0, T)). \end{aligned}$$

Combining this with (90) we obtain the conclusion (83).

### Boundary estimates

Let us now do some estimates close to the boundary. For this, we consider a finite covering  $\{\mathcal{O}_i\}_{i=1}^K$  of  $\partial\Omega_F(0)$  included in the set  $\{y : \text{dist}(y, \partial\Omega_F(0)) < \alpha/4\}$ . For each  $1 \leq i \leq K$ , we consider a  $C^4$  function

$$\begin{aligned} \theta^i : (-M_1^i, M_1^i) \times (-M_2^i, M_2^i) & \rightarrow \partial\Omega_F(0) \cap \mathcal{O}_i \\ (\phi_1, \phi_2) & \rightarrow \theta^i(\phi_1, \phi_2) \end{aligned} \quad (92)$$

satisfying

$$\partial_{\phi_1} \theta^i \cdot \partial_{\phi_2} \theta^i = 0, \quad |\partial_{\phi_1} \theta^i| = 1, \quad |\partial_{\phi_2} \theta^i| \geq C > 0.$$

Then, we define the mapping

$$\tilde{\theta}^i(t, \phi_1, \phi_2) = \chi(t, \theta^i(\phi_1, \phi_2)) \quad \forall (t, \phi_1, \phi_2) \in (0, T) \times (-M_1^i, M_1^i) \times (-M_2^i, M_2^i)$$

which sends  $(0, T) \times (-M_1^i, M_1^i) \times (-M_2^i, M_2^i)$  in the set  $\partial\Omega_F(t) \cap \tilde{\mathcal{O}}_i$ , where  $\tilde{\mathcal{O}}_i := \chi(t, \mathcal{O}_i)$ . Recall that the flow  $\chi$  was defined in Lemma 1.

Furthermore, we perform the change of variables

$$\begin{cases} x(t, \phi_1, \phi_2, r) = r\tilde{n}(t, \phi_1, \phi_2) + \tilde{\theta}^i(t, \phi_1, \phi_2) \in \tilde{\mathcal{O}}_i \cap \Omega_F(t) \\ r \in (0, r_0^i), (\phi_1, \phi_2) \in (-M_1^i, M_1^i) \times (-M_2^i, M_2^i), \end{cases} \quad (93)$$

where  $n(\phi_1, \phi_2)$  (resp.  $\tilde{n}(t, \phi_1, \phi_2) := \nabla\chi(t, \theta^i(\phi_1, \phi_2))n(\phi_1, \phi_2)$ ) denotes the inward unit normal vector to  $\partial\Omega_F(0)$  (resp.  $\partial\Omega_F(t)$ ).

In what follows, we will establish some estimates of  $u$  and  $\rho$  in  $\Omega_F(t) \cap \tilde{\mathcal{O}}_i$  for each  $1 \leq i \leq K$  (see Lemmas 15, 16 and 17 below). For the sake of simplicity, we drop the index  $i$  from now on. In the new variables, we define

$$\begin{cases} e_1(t, \phi_1, \phi_2) = \partial_{\phi_1} \tilde{\theta}(t, \phi_1, \phi_2) \quad \text{and} \quad e_2(t, \phi_1, \phi_2) = \frac{\partial_{\phi_2} \tilde{\theta}}{|\partial_{\phi_2} \tilde{\theta}|}(t, \phi_1, \phi_2), \\ \forall (t, \phi_1, \phi_2) \in (0, T) \times (-M_1, M_1) \times (-M_2, M_2). \end{cases}$$

Observe that  $e_1 \cdot e_2 = \partial_{\phi_1} \theta \cdot \frac{\partial_{\phi_2} \theta}{|\partial_{\phi_2} \theta|} = 0$ ,  $e_1 \cdot \tilde{n} = \partial_{\phi_1} \theta \cdot n = 0$  and  $e_2 \cdot \tilde{n} = \frac{\partial_{\phi_2} \theta}{|\partial_{\phi_2} \theta|} \cdot n = 0$  and so we can write

$$\partial_{\phi_1} \tilde{n} = \alpha e_1 + \beta e_2 \quad \text{and} \quad \partial_{\phi_2} \tilde{n} = \alpha' e_1 + \beta' e_2, \quad (94)$$

for some  $\alpha, \beta, \alpha', \beta' \in H^3(0, T; C^3([-M_1, M_1] \times [-M_2, M_2]))$ . Let us compute the Jacobian of this change of variables:

$$\partial_{\phi_1} x = (1 + \alpha r)e_1 + \beta r e_2, \quad \partial_{\phi_2} x = \alpha' r e_1 + (\beta' r + |\partial_{\phi_2} \tilde{\theta}|)e_2 \quad \text{and} \quad \partial_r x = \tilde{n}, \quad (95)$$

that is to say,

$$\begin{pmatrix} \partial_{\phi_1} x \\ \partial_{\phi_2} x \\ \partial_r x \end{pmatrix} = \begin{pmatrix} 1 + \alpha r & \beta r & 0 \\ \alpha' r & \beta' r + |\partial_{\phi_2} \tilde{\theta}| & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ \tilde{n} \end{pmatrix}.$$

Inversely, we have

$$\begin{pmatrix} \partial_x \phi_1 \\ \partial_x \phi_2 \\ \partial_x r \end{pmatrix} = \frac{1}{\bar{J}} \begin{pmatrix} \beta' r + |\partial_{\phi_2} \tilde{\theta}| & -\beta r & 0 \\ -\alpha' r & 1 + \alpha r & 0 \\ 0 & 0 & \bar{J} \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ \tilde{n} \end{pmatrix},$$

where  $\bar{J} := (1 + \alpha r)(\beta' r + |\partial_{\phi_2} \tilde{\theta}|) - \alpha' \beta r^2$  is the Jacobian.

Consequently, we have

$$\nabla_x = \partial_x = \frac{1}{\bar{J}}(A_1 e_1 + A_2 e_2) \partial_{\phi_1} + \frac{1}{\bar{J}}(A_3 e_1 + A_4 e_2) \partial_{\phi_2} + \tilde{n} \partial_r \quad (96)$$

where

$$A_1 = \beta' r + |\partial_{\phi_2} \tilde{\theta}|, \quad A_2 = -\beta r, \quad A_3 = -\alpha' r \quad \text{and} \quad A_4 = 1 + \alpha r.$$

Now, if we denote

$$\tilde{\rho}(t, \phi, r) := \rho(t, x(t, \phi, r)), \quad \tilde{u}(t, \phi, r) := u(t, x(t, \phi, r)) \quad (97)$$

and

$$\tilde{f}_0(t, \phi, r) := f_0(t, x(t, \phi, r)), \quad \tilde{f}_1(t, \phi, r) := f_1(t, x(t, \phi, r)),$$

we can rewrite the equations of (63) in the new variables  $(\phi_1, \phi_2, r)$  as follows:

$$\left\{ \begin{array}{l} \tilde{\rho}_t + \tilde{u} \cdot \left( \frac{1}{\bar{J}}(A_1 e_1 + A_2 e_2) \partial_{\phi_1} \tilde{\rho} + \frac{1}{\bar{J}}(A_3 e_1 + A_4 e_2) \partial_{\phi_2} \tilde{\rho} + \tilde{n} \partial_r \tilde{\rho} \right) \\ - ((r \nabla \chi_t \circ \theta) n + \chi_t \circ \theta) \left( \frac{1}{\bar{J}}(A_1 e_1 + A_2 e_2) \partial_{\phi_1} \tilde{\rho} + \frac{1}{\bar{J}}(A_3 e_1 + A_4 e_2) \partial_{\phi_2} \tilde{\rho} + \tilde{n} \partial_r \tilde{\rho} \right) \\ + \tilde{\rho} \left( \frac{1}{\bar{J}}(A_1 e_1 + A_2 e_2) \cdot \partial_{\phi_1} \tilde{u} + \frac{1}{\bar{J}}(A_3 e_1 + A_4 e_2) \cdot \partial_{\phi_2} \tilde{u} + \tilde{n} \cdot \partial_r \tilde{u} \right) = \tilde{f}_0 \\ (t, \phi_1, \phi_2, r) \in (0, T) \times (-M_1, M_1) \times (-M_2, M_2) \times (0, r_0) \end{array} \right. \quad (98)$$

and

$$\left\{ \begin{array}{l} \tilde{u}_t - \frac{\mu}{\bar{J}^2} ((A_1^2 + A_2^2) \partial_{\phi_1 \phi_1}^2 \tilde{u} + 2(A_1 A_3 + A_2 A_4) \partial_{\phi_1 \phi_2}^2 \tilde{u} + (A_3^2 + A_4^2) \partial_{\phi_2 \phi_2}^2 \tilde{u} + \bar{J}^2 \partial_{rr}^2 \tilde{u}) + P_1(D) \tilde{u} \\ - ((r \nabla \chi_t \circ \theta) n + \chi_t \circ \theta) \left( \frac{1}{\bar{J}}(A_1 e_1 + A_2 e_2) \partial_{\phi_1} \tilde{u} + \frac{1}{\bar{J}}(A_3 e_1 + A_4 e_2) \partial_{\phi_2} \tilde{u} + \tilde{n} \partial_r \tilde{u} \right) \\ + \frac{1}{\bar{J}} \partial_{\phi_1} \left( \frac{\mu + \mu'}{\bar{\rho}} \frac{d\tilde{\rho}}{dt} + p^0 \tilde{\rho} \right) (A_1 e_1 + A_2 e_2) + \frac{1}{\bar{J}} \partial_{\phi_2} \left( \frac{\mu + \mu'}{\bar{\rho}} \frac{d\tilde{\rho}}{dt} + p^0 \tilde{\rho} \right) (A_3 e_1 + A_4 e_2) \\ + \partial_r \left( \frac{\mu + \mu'}{\bar{\rho}} \frac{d\tilde{\rho}}{dt} + p^0 \tilde{\rho} \right) \tilde{n} = \tilde{f}_1 + \frac{\mu + \mu'}{\bar{\rho}} \nabla_x f_0 \\ (t, \phi_1, \phi_2, r) \in (0, T) \times (-M_1, M_1) \times (-M_2, M_2) \times (0, r_0), \end{array} \right. \quad (99)$$

where  $P_1(D)$  is a first order differential operator in the  $x$  variables. In order to obtain this last equation, observe that we have used that

$$u_t - \mu \Delta u + \frac{\mu + \mu'}{\bar{\rho}} \nabla \left( \frac{d\rho}{dt} - f_0 \right) + p^0 \nabla \rho = f_1$$

and then we have rewritten this in the new variables. In the new variables  $\frac{d\tilde{\rho}}{dt}$ , which is given by the two first lines in (98), satisfies

$$\frac{d\tilde{\rho}}{dt}(t, \phi, r) = \frac{d\rho}{dt}(t, x(t, \phi, r)) \quad \forall (t, \phi, r) \in (0, T) \times (-M_1, M_1) \times (-M_2, M_2) \times (0, r_0), \quad (100)$$

where  $x(t, \phi, r)$  is given by (93).

Let  $\zeta_{1,1}^i \in C_0^\infty(\mathcal{O}_i)$  and

$$\zeta_1^i(t, x) := \zeta_{1,1}^i(\psi(t, x)), \quad \forall x \in \tilde{\mathcal{O}}_i, \quad t \in (0, T),$$

where  $\psi$  is the flow (extending  $\chi_S^{-1}$ ) defined in Corollary 2. We will also drop here the index  $i$ .

We first estimate the tangential derivatives:

**Lemma 15** *Let  $k = 1, 2, 3$ . For any  $\delta > 0$ , there exists a positive constant  $C$  such that*

$$\begin{aligned} & \|\zeta_1 D_\phi^k \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi^k u\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)} + \|\zeta_1 D_\phi^k \frac{d\rho}{dt}\|_{L_T^2(L^2)} \\ & \leq \delta(\|\zeta_1 D_\phi^k \rho\|_{L_T^2(L^2)} + \|u\|_{L_T^2(H^2)} + \|\rho\|_{L_T^2(H^1)}) + C(\|u\|_{L_T^2(H^k)} + \|\omega\|_{L_T^2} + \|\rho_0\|_{H^k} \\ & \quad + \|u_0\|_{H^k} + \|f_0\|_{L_T^2(H^k)} + \|f_1\|_{L_T^2(H^{k-1})} + N^{3/2}(0, T)). \end{aligned} \quad (101)$$

**Proof:** The proof of this lemma is very similar to that of the second part of Lemma 14. In fact, we first apply the operator  $\frac{p^0}{\bar{\rho}} \zeta_1^2 D_\phi^k$  to the equation of  $\rho$  in (63) and multiply it by  $D_\phi^k \rho$ :

$$\frac{p^0}{\bar{\rho}} \iint_{Z_s} \zeta_1^2 D_\phi^k (\rho_t + (u \cdot \nabla_x) \rho + \bar{p}(\nabla_x \cdot u) - f_0) D_\phi^k \rho \, dx \, d\tau = 0. \quad (102)$$

We recall now that (see (96))

$$\nabla_x = B_1 \partial_{\phi_1} + B_2 \partial_{\phi_2} + B_3 \partial_r,$$

where  $B_j$  ( $j = 1, 2, 3$ ) belongs to  $H^3(0, T; C^3([-M_1, M_1] \times [-M_2, M_2]))$ . Then, we notice that

$$D_\phi^k \nabla_x h = \nabla_x D_\phi^k h + \sum_{\ell=1}^k \binom{k}{\ell} (D_\phi^\ell B_1 D_\phi^{k-\ell} \partial_{\phi_1} + D_\phi^\ell B_2 D_\phi^{k-\ell} \partial_{\phi_2} + D_\phi^\ell B_3 D_\phi^{k-\ell} \partial_r) h \quad (103)$$

(the same can be done for  $D_r^\ell \nabla_x h$ ). Using this expression, we have the following for the nonlinear term:

$$\iint_{Z_s} \zeta_1^2 D_\phi^k (u \cdot \nabla_x \rho) D_\phi^k \rho \, dx \, d\tau \geq \iint_{Z_s} \zeta_1^2 (u \cdot \nabla_x D_\phi^k \rho) D_\phi^k \rho \, dx \, d\tau - CN^3(0, T). \quad (104)$$

As long as the time derivative is concerned, we get

$$\begin{aligned} & \iint_{Z_s} \zeta_1^2 D_\phi^k \rho_t D_\phi^k \rho \, dx \, d\tau = \frac{1}{2} (\|\zeta_1 D_\phi^k \rho\|_{L^2}^2(s) - \|\zeta_1(0) D_\phi^k \rho_0\|_{L^2}^2) - \iint_{Z_s} \zeta_1 \zeta_{1,t} |D_\phi^k \rho|^2 \, dx \, d\tau \\ & \quad - \frac{1}{2} \iint_{Z_s} \zeta_1^2 u \cdot \nabla_x (|D_\phi^k \rho|^2) \, dx \, d\tau - \frac{1}{2} \iint_{Z_s} \nabla_x \cdot (\zeta_1^2 u) |D_\phi^k \rho|^2 \, dx \, d\tau \end{aligned}$$

Observe that the last terms in the first and second lines can be estimated by  $N^3(0, T)$  and that the first term in the second line simplifies with the first term in the right hand side of (104). Combining with (102), this implies that

$$\begin{aligned} & \|\zeta_1 D_\phi^k \rho\|_{L^2}^2(s) + p^0 \iint_{Z_s} \zeta_1^2 D_\phi^k (\nabla_x \cdot u) D_\phi^k \rho \, dx \, d\tau \\ & \leq \delta \|\zeta_1 D_\phi^k \rho\|_{L_T^2(L^2)}^2 + C(\|\rho_0\|_{H^k}^2 + \|f_0\|_{L_T^2(H^k)}^2 + N^3(0, T)). \end{aligned} \quad (105)$$

Then, we apply the operator  $\zeta_1^2 D_\phi^k$  to the  $j$ -th equation of  $u$  in (63), multiply it by  $D_\phi^k u_j$  and sum up in  $j$ :

$$\sum_{j=1}^3 \iint_{Z_s} \zeta_1^2 D_\phi^k (u_t - 2\mu \nabla_x \cdot \epsilon(u) - \mu' \nabla_x (\nabla \cdot u) + p^0 \nabla_x \rho - f_1)_j D_\phi^k u_j \, dx \, d\tau = 0. \quad (106)$$

Observe that, since  $\zeta_1$  has compact support, we have

$$- \iint_{Z_s} \zeta_1^2 D_\phi^k \nabla_x \cdot \epsilon(u)_j D_\phi^k u_j \, dx \, d\tau = (-1)^{k+1} \iint_{Z_s} \nabla_x \cdot \epsilon(u)_j D_\phi^k (\zeta_1^2 D_\phi^k u_j) \, dx \, d\tau. \quad (107)$$

Integrating by parts on the  $x$  variable, we deduce:

$$\begin{aligned} & - \iint_{Z_s} \zeta_1^2 D_\phi^k \nabla_x \cdot \epsilon(u)_j D_\phi^k u_j \, dx \, d\tau = (-1)^k \iint_{Z_s} \epsilon(u)_j \cdot \nabla_x D_\phi^k (\zeta_1^2 D_\phi^k u_j) \, dx \, d\tau \\ & \quad + (-1)^k \iint_{\Sigma_s} (\epsilon(u)_j \cdot n) D_\phi^k (\zeta_1^2 D_\phi^k u_j) \, d\sigma \, d\tau \\ & \geq (-1)^k \iint_{Z_s} \epsilon(u)_j \cdot \nabla_x D_\phi^k (\zeta_1^2 D_\phi^k u_j) \, dx \, d\tau - C(\delta \|u\|_{L_T^2(H^2)}^2 + \|\omega\|_{L_T^2}^2), \end{aligned} \quad (108)$$

where we have used that  $|D_\phi^\ell u| \leq C|\omega|$  for all  $|\ell| \geq 1$ .

Let us take  $h = \zeta_1^2 D_\phi^k u_j$  in (103). For each term of the sum, we have

$$\begin{aligned} & \left| \iint_{Z_s} \epsilon(u)_j \cdot (D_\phi^\ell B_1 D_\phi^{k-\ell} \partial_{\phi_1} (\zeta_1^2 D_\phi^k u_j) + D_\phi^\ell B_2 D_\phi^{k-\ell} \partial_{\phi_2} (\zeta_1^2 D_\phi^k u_j) + D_\phi^\ell B_3 D_\phi^{k-\ell} \partial_\tau (\zeta_1^2 D_\phi^k u_j)) dx d\tau \right| \\ & \leq \left| \iint_{Z_s} D_\phi^{k-\ell} (D_\phi^\ell B_1 \epsilon(u)_j) \partial_{\phi_1} (\zeta_1^2 D_\phi^k u_j) + D_\phi^{k-\ell} (D_\phi^\ell B_2 \epsilon(u)_j) \partial_{\phi_2} (\zeta_1^2 D_\phi^k u_j) dx d\tau \right| \\ & + \left| \iint_{Z_s} D_\phi^{k-\ell} (D_\phi^\ell B_3 \epsilon(u)_j) \partial_\tau (\zeta_1^2 D_\phi^k u_j) dx d\tau \right| \leq C(\|u\|_{L_T^2(H^k)}^2 + \delta \|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)}^2). \end{aligned} \quad (109)$$

This implies that

$$\begin{aligned} & (-1)^k \iint_{Z_s} \epsilon(u)_j \cdot \nabla_x D_\phi^k (\zeta_1^2 D_\phi^k u_j) dx d\tau \\ & \geq (-1)^k \iint_{Z_s} \epsilon(u)_j \cdot D_\phi^k \nabla_x (\zeta_1^2 D_\phi^k u_j) dx d\tau - C(\|u\|_{L_T^2(H^k)}^2 + \delta \|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)}^2) \\ & \geq \iint_{Z_s} D_\phi^k \epsilon(u)_j \cdot \nabla_x (\zeta_1^2 D_\phi^k u_j) dx d\tau - C(\|u\|_{L_T^2(H^k)}^2 + \delta \|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)}^2). \end{aligned} \quad (110)$$

Analogously, we can prove that

$$\begin{aligned} & \iint_{Z_s} D_\phi^k \epsilon(u)_j \cdot \nabla_x (\zeta_1^2 D_\phi^k u_j) dx d\tau \\ & \geq \iint_{Z_s} \epsilon(\zeta_1 D_\phi^k u)_j \cdot \nabla_x (\zeta_1 D_\phi^k u_j) dx d\tau - C(\|u\|_{L_T^2(H^k)}^2 + \delta \|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)}^2). \end{aligned} \quad (111)$$

So, from (107)-(111), we have for the second term in (106):

$$\begin{aligned} & - \iint_{Z_s} \zeta_1^2 D_\phi^k \nabla_x \cdot \epsilon(u)_j D_\phi^k u_j dx d\tau \geq \iint_{Z_s} \epsilon(\zeta_1 D_\phi^k u)_j \cdot \nabla_x (\zeta_1 D_\phi^k u_j) dx d\tau \\ & - C(\|u\|_{L_T^2(H^k)}^2 + \delta \|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)}^2 + \delta \|u\|_{L_T^2(H^2)}^2 + \|\omega\|_{L_T^2}^2). \end{aligned} \quad (112)$$

The same can also be done for the third term in (106):

$$\begin{aligned} & - \sum_{j=1}^3 \iint_{Z_s} \zeta_1^2 D_\phi^k \partial_{x_j} (\nabla_x \cdot u) D_\phi^k u_j dx d\tau \geq \iint_{Z_s} |\nabla_x \cdot (\zeta_1 D_\phi^k u)|^2 dx d\tau \\ & - C(\|u\|_{L_T^2(H^k)}^2 + \delta \|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)}^2 + \delta \|u\|_{L_T^2(H^2)}^2 + \|\omega\|_{L_T^2}^2). \end{aligned} \quad (113)$$

Equations (112) and (113) provide the estimate of  $\|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)}^2$  in (101).

Then, using the same arguments, one can prove that

$$\begin{aligned} & p^0 \sum_{j=1}^3 \iint_{Z_s} \zeta_1^2 D_\phi^k \partial_{x_j} \rho D_\phi^k u_j dx d\tau \geq -p^0 \iint_{Z_s} \zeta_1^2 D_\phi^k (\nabla_x \cdot u) D_\phi^k \rho dx d\tau \\ & - \delta(\|\rho\|_{L_T^2(H^1)}^2 + \|\zeta_1 D_\phi^k \rho\|_{L_T^2(L^2)}^2) + C(\|\omega\|_{L_T^2}^2 + \|u\|_{L_T^2(H^k)}^2). \end{aligned} \quad (114)$$

Then, inequality (101) is obtained by reassembling (105) (for  $\|\zeta_1 D_\phi^k \rho\|_{L_T^2(L^2)}$ ), (106) (for  $\|\zeta_1 D_\phi^k u\|_{L^\infty(L^2)}$ ), (112)-(113) (for  $\|\zeta_1 D_\phi^k u\|_{L_T^2(H^1)}^2$ ) and, consequently,  $\|\zeta_1 D_\phi^k \frac{d\rho}{dt}\|_{L_T^2(L^2)}$  and (114).

Let us now take the derivative with respect to  $r$  of equation (98), multiply it by  $\mu/\bar{\rho}$  and sum it with equation (99) multiplied by  $\tilde{n}$  (which we can symbolically write as  $(\mu/\bar{\rho})(98)_r + \tilde{n} \cdot (99)$ ):

$$\left\{ \begin{aligned} & \frac{2\mu + \mu'}{\bar{\rho}} \left( \frac{d\tilde{\rho}}{dt} \right)_r + p^0 \tilde{\rho}_r = -\frac{\mu}{\bar{J}^2} [(A_1^2 + A_2^2)\tilde{n} \cdot \tilde{u}_{\phi_1\phi_1} + 2(A_1A_3 + A_2A_4)\tilde{n} \cdot \tilde{u}_{\phi_1\phi_2} \\ & + (A_3^2 + A_4^2)\tilde{n} \cdot \tilde{u}_{\phi_2\phi_2} - \bar{J}(A_1e_1 + A_2e_2) \cdot \tilde{u}_{\phi_1r} - \bar{J}(A_3e_1 + A_4e_2) \cdot \tilde{u}_{\phi_2r}] \\ & - \tilde{n} \cdot \tilde{u}_t - ((r\nabla\chi_t \circ \theta)n + \chi_t \circ \theta) \left( \frac{1}{\bar{J}}(A_1e_1 + A_2e_2)\partial_{\phi_1}\tilde{u} + \frac{1}{\bar{J}}(A_3e_1 + A_4e_2)\partial_{\phi_2}\tilde{u} + \tilde{n}\partial_r\tilde{u} \right) \cdot \tilde{n} \\ & - P_1(D)\tilde{u} \cdot \tilde{n} + \tilde{f}_1 \cdot \tilde{n} + \frac{2\mu + \mu'}{\bar{\rho}} \tilde{f}_{0,r}, \quad \phi \in (-M_1, M_1) \times (-M_2, M_2), r \in (0, r_0). \end{aligned} \right. \quad (115)$$

Recall that the third line in this expression corresponds to  $-u_t(t, x) \cdot \tilde{n}$ .

**Lemma 16** For  $0 \leq k + \ell \leq 2$ , we have

$$\begin{aligned} & \|\zeta_1 D_\phi^k D_r^{\ell+1} \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi^k D_r^{\ell+1} \rho\|_{L_T^2(L^2)} + \|\zeta_1 D_\phi^k D_r^{\ell+1} \frac{d\rho}{dt}\|_{L_T^2(L^2)} \\ & \leq C(\|\zeta_1 D_\phi^{k+1} D_r^\ell u\|_{L_T^2(H^1)} + \|u_t\|_{L_T^2(H^{k+\ell})} + \|u\|_{L_T^2(H^{k+\ell+1})} + \|\rho_0\|_{H^{k+\ell+1}} \\ & + \|f_0\|_{L_T^2(H^{k+\ell+1})} + \|f_1\|_{L_T^2(H^{k+\ell})} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (116)$$

**Proof:** We take the operator  $D_\phi^k D_r^\ell$  in equation (115), multiply this by  $\zeta_1$  and square the whole expression. Using the definition (97) and (100), the left hand side term is

$$\begin{aligned} & \iint_{Z_s} \zeta_1^2 \left( \left| \frac{2\mu + \mu'}{\bar{\rho}} \right|^2 \left| D_\phi^k D_r^{\ell+1} \frac{d\rho}{dt} \right|^2 + (p^0)^2 |D_\phi^k D_r^{\ell+1} \rho|^2 + \frac{4\mu + 2\mu'}{\bar{\rho}} p^0 D_\phi^k D_r^{\ell+1} \frac{d\rho}{dt} D_\phi^k D_r^{\ell+1} \rho \right) dx d\tau. \\ & = \left\| \zeta_1 \frac{2\mu + \mu'}{\bar{\rho}} D_\phi^k D_r^{\ell+1} \frac{d\rho}{dt} \right\|_{L_T^2(L^2)}^2 + \|\zeta_1 p^0 D_\phi^k D_r^{\ell+1} \rho\|_{L_T^2(L^2)}^2 + \frac{4\mu + 2\mu'}{\bar{\rho}} p^0 \left( \zeta_1 D_\phi^k D_r^{\ell+1} \frac{d\rho}{dt}, \zeta_1 D_\phi^k D_r^{\ell+1} \rho \right)_{L_T^2(L^2)}. \end{aligned}$$

Let us see that the last term readily provides the  $L_T^\infty(L^2)$  norm of  $\zeta_1 D_\phi^k D_r^{\ell+1} \rho$  in terms of  $\|\rho_0\|_{H^{k+\ell+1}}$  and  $N^{3/2}(0, T)$ . First, thanks to the definition of  $\frac{d\rho}{dt}$  (see (82)), we have

$$\iint_{Z_s} \zeta_1 D_\phi^k D_r^{\ell+1} (\rho_t + (u \cdot \nabla)\rho) \zeta_1 D_\phi^k D_r^{\ell+1} \rho dx d\tau \quad (117)$$

As long as the first term is concerned, we use (95) and the fact that  $|\tilde{n}_t| + |\tilde{\theta}_t| + \|\zeta_{1,t}\|_{L_T^\infty(L^\infty)} \leq CN(0, T)$  (see (22)) and we obtain

$$(\zeta_1 D_\phi^k D_r^{\ell+1} \rho_t)(t, x) = (\zeta_1 D_\phi^k D_r^{\ell+1} \rho)_t(t, x) + h(t, x),$$

where  $\|h\|_{L_T^2(L^2)} \leq CN^2(0, T)$ . Consequently, the first term in (117) is given by

$$\begin{aligned} & \frac{1}{2} \int_{\Omega_F(s)} |\zeta_1 D_\phi^k D_r^{\ell+1} \rho|^2(s) dx - \frac{1}{2} \int_{\Omega_F(0)} |\zeta_1(0) D_\phi^k D_r^{\ell+1} \rho_0|^2 dx + \iint_{Z_s} h \zeta_1 D_\phi^k D_r^{\ell+1} \rho dx d\tau \\ & - \frac{1}{2} \iint_{Z_s} \nabla \cdot (|\zeta_1 D_\phi^k D_r^{\ell+1} \rho|^2 u) dx d\tau \geq \|\zeta_1 D_\phi^k D_r^{\ell+1} \rho\|_{L^2(s)}^2 - \|\zeta_1(0) D_\phi^k D_r^{\ell+1} \rho_0\|_{L^2}^2 \\ & - \delta \|\zeta_1 D_\phi^k D_r^{\ell+1} \rho\|_{L_T^2(L^2)}^2 - CN^4(0, T) - \frac{1}{2} \iint_{Z_s} \nabla \cdot (|\zeta_1 D_\phi^k D_r^{\ell+1} \rho|^2 u) dx d\tau, \end{aligned} \quad (118)$$

for  $\delta > 0$  small enough. Using (103), the second term in (117) is estimated from below by

$$\frac{1}{2} \iint_{Z_s} \nabla \cdot (|\zeta_1 D_\phi^k D_r^{\ell+1} \rho|^2 u) dx d\tau - CN^3(0, T).$$

Putting this together with (118), we obtain (116).

**Lemma 17** For  $0 \leq k + \ell \leq 2$ , we have

$$\begin{aligned} \|\zeta_1 D_\phi^k u\|_{L_T^2(H^{2+\ell})} + \|\zeta_1 D_\phi^k \rho\|_{L_T^2(H^{1+\ell})} &\leq C \left( \left\| \zeta_1 D_\phi^k \frac{d\rho}{dt} \right\|_{L_T^2(H^{1+\ell})} + \|u_t\|_{L_T^2(H^{k+\ell})} + \|u\|_{L_T^2(H^{k+\ell+1})} \right. \\ &\quad \left. + \|\rho\|_{L_T^2(H^{k+\ell})} + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} + \|f_0\|_{L_T^2(H^{1+k+\ell})} + \|f_1\|_{L_T^2(H^{k+\ell})} \right). \end{aligned} \quad (119)$$

**Proof:** We regard the equation of the fluid in (63) as a stationary Stokes system for each  $s \in (0, T)$ :

$$\begin{cases} -2\mu \nabla \cdot (\epsilon(u)) - \mu' \nabla \cdot ((\nabla \cdot u)Id) + p^0 \nabla \rho = f_1 - u_t & \text{in } \mathcal{O}_i \cap \Omega_F(t), \\ \nabla \cdot u = \frac{1}{\rho} \left( f_0 - \frac{d\rho}{dt} \right) & \text{in } \mathcal{O}_i \cap \Omega_F(t), \\ u = (\dot{a} + \omega \wedge (x - a)) \mathbf{1}_{\partial\Omega_S(t)} & \text{on } \mathcal{O}_i \cap \partial\Omega_F(t). \end{cases}$$

Recall that  $\frac{d\rho}{dt}$  was given in (82). Now, we differentiate the equation of  $u$  with respect to  $\phi = (\phi_1, \phi_2)$   $k$  times and we multiply by  $\zeta_1$ . Denoting  $(u_k, \rho_k) := \zeta_1 D_\phi^k(u, \rho)$ , we have

$$\begin{cases} -2\mu \nabla \cdot (\epsilon(u_k)) - \mu' \nabla \cdot ((\nabla \cdot u_k)Id) + p^0 \nabla (\rho_k) = f_{1,k} & \text{in } \mathcal{O}_i \cap \Omega_F(t), \\ \nabla \cdot u_k = \frac{\zeta_1}{\rho} \left( D_\phi^k f_0 - D_\phi^k \frac{d\rho}{dt} \right) + P_k(D)u & \text{in } \mathcal{O}_i \cap \Omega_F(t), \\ u_k = (\zeta_1 D_\phi^k (\dot{a} + \omega \wedge (x - a))) \mathbf{1}_{\partial\Omega_S(t)} & \text{on } \mathcal{O}_i \cap \partial\Omega_F(t), \end{cases}$$

where

$$f_{1,k} = \zeta_1 D_\phi^k f_1 - \zeta_1 D_\phi^k u_t + P_{k+1}(D)u + P_k(D)\rho,$$

with  $P_j(D)$  a differential operator in the  $x$  variable of order  $j$ . Observe that the  $L_T^2(H^4(\mathcal{O}_i \cap \Omega_F(t)))$  norm of the boundary condition can be estimated by  $C(\|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2})$ . Finally, we use the following classical regularity result for the stationary Stokes problem (see, for instance, [22]):

Let  $U$  be a  $C^{m+2}$  domain,  $m > 0$ ,  $\alpha > 0$ ,  $f \in H^m(U)$ ,  $g_0 \in H^{m+1}(U)$  and  $g_1 \in H^{m+3/2}(\partial U)$ . Then, the solution  $(h, \pi)$  of

$$\begin{cases} -\alpha \Delta h + \nabla \pi = f & \text{in } U, \\ \nabla \cdot h = g_0 & \text{in } U, \\ h = g_1 & \text{on } \partial U \end{cases}$$

belongs to  $H^{m+2}(U) \times H^{m+1}(U)$  and there exists a constant  $C > 0$  such that

$$\|(h, \pi)\|_{H^{m+2}(U) \times H^{m+1}(U)} \leq C(\|f\|_{H^m(U)} + \|g_0\|_{H^{m+1}(U)} + \|g_1\|_{H^{m+3/2}(\partial U)}).$$

## 4.2 Combination of Lemmas and conclusion

In this paragraph we will gather all the technical results previously stated and conclude the proof of Proposition 8.

1) We apply Lemma 12 for  $k = 0$ :

$$\begin{aligned} &\|u\|_{L_T^\infty(L^2)} + \|\rho\|_{L_T^\infty(L^2)} + \|u\|_{L_T^2(H^1)} + \|\dot{a}\|_{L_T^\infty} + \|\omega\|_{L_T^\infty} \\ &\leq \delta(\|\rho\|_{L_T^2(L^2)} + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2}) + C(\|f_0\|_{L_T^2(L^2)} + \|f_1\|_{L_T^2(L^2)} + \|f_2\|_{L_T^2} + \|f_3\|_{L_T^2} \\ &\quad + \|u_0\|_{L^2} + \|\rho_0\|_{L^2} + |a_0| + |\omega_0| + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (120)$$

2) We apply Lemma 13 for  $k = 0$ :

$$\begin{aligned}
& \|u_t\|_{L_T^2(L^2)} + \|\rho_t\|_{L_T^2(L^2)} + \|u\|_{L_T^\infty(H^1)} + \|\dot{a}\|_{L_T^2} + \|\dot{\omega}\|_{L_T^2} \\
& \leq C(\|u\|_{L_T^2(H^1)} + \|\rho\|_{L_T^\infty(L^2)} + \|f_0\|_{L_T^2(L^2)} + \|f_1\|_{L_T^2(L^2)} + \|f_2\|_{L_T^2} + \|f_3\|_{L_T^2} \\
& \quad + \|u_0\|_{H^1} + N^{3/2}(0, T) + N^2(0, T)).
\end{aligned} \tag{121}$$

3) We apply Lemma 12 for  $k = 1$ :

$$\begin{aligned}
& \|u\|_{W_T^{1,\infty}(L^2)} + \|\rho\|_{W_T^{1,\infty}(L^2)} + \|u\|_{H_T^1(H^1)} + \|\dot{a}\|_{W_T^{1,\infty}} + \|\omega\|_{W_T^{1,\infty}} \leq \delta(\|\rho\|_{H_T^1(L^2)} + \|\dot{a}\|_{H_T^1} \\
& + \|\omega\|_{H_T^1}) + C(\|f_0\|_{H_T^1(L^2)} + \|f_1\|_{H_T^1(L^2)} + \|f_2\|_{H_T^1} + \|f_3\|_{H_T^1} + \|u_0\|_{L^2} + \|\rho_0\|_{L^2} + |\dot{a}(0)| + |\omega_0| \\
& + \|\partial_t u(0, \cdot)\|_{L^2} + \|\partial_t \rho(0, \cdot)\|_{L^2} + |\partial_t^2 a(0)| + |\partial_t \omega(0)| + N^{3/2}(0, T) + N^2(0, T)).
\end{aligned} \tag{122}$$

4) We apply Lemma 10:

$$\|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} \leq C\|u\|_{L_T^2(H^1)}. \tag{123}$$

Putting together estimates (120), (121), (122) and (123) and taking  $\delta > 0$  sufficiently small, we have:

$$\begin{aligned}
& \|u\|_{W_T^{1,\infty}(L^2)} + \|u\|_{L_T^\infty(H^1)} + \|u\|_{H_T^1(H^1)} + \|\rho\|_{W_T^{1,\infty}(L^2)} + \|\rho_t\|_{L_T^2(L^2)} \\
& + \|\dot{a}\|_{W_T^{1,\infty}} + \|\dot{a}\|_{H_T^1} + \|\omega\|_{W_T^{1,\infty}} + \|\omega\|_{H_T^1} \leq \delta\|\rho\|_{L_T^2(L^2)} + C(\|f_0\|_{H_T^1(L^2)} + \|f_1\|_{H_T^1(L^2)} + \|f_2\|_{H_T^1} \\
& + \|f_3\|_{H_T^1} + \|\partial_t u(0, \cdot)\|_{L^2} + \|u_0\|_{H^1} + \|\partial_t \rho(0, \cdot)\|_{L^2} + \|\rho_0\|_{L^2} + |\partial_t^2 a(0)| + |a_0| \\
& + |\partial_t \omega(0)| + |\omega_0| + N^{3/2}(0, T) + N^2(0, T)).
\end{aligned} \tag{124}$$

5) We apply Lemma 14 for  $k = 1$ :

$$\begin{aligned}
& \|\zeta_0 \nabla \rho\|_{L_T^\infty(L^2)} + \|\zeta_0 \nabla \rho\|_{L_T^2(L^2)} + \|\zeta_0 \nabla u\|_{L_T^\infty(L^2)} + \|\zeta_0 \nabla u\|_{L_T^2(H^1)} + \|\zeta_0 \nabla \frac{d\rho}{dt}\|_{L_T^2(L^2)} \\
& \leq C(\|u\|_{L_T^\infty(L^2)} + \|u\|_{L_T^2(H^1)} + \|\rho_0\|_{H^1} + \|u_0\|_{H^1} + \|f_0\|_{L_T^2(H^1)} + \|f_1\|_{L_T^2(L^2)} + N^{3/2}(0, T)).
\end{aligned} \tag{125}$$

6) We apply Lemma 15 for  $k = 1$ :

$$\begin{aligned}
& \|\zeta_1 D_\phi \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi u\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi u\|_{L_T^2(H^1)} + \|\zeta_1 D_\phi \frac{d\rho}{dt}\|_{L_T^2(L^2)} \\
& \leq \delta(\|\zeta_1 D_\phi \rho\|_{L_T^2(L^2)} + \|u\|_{L_T^2(H^2)} + \|\rho\|_{L_T^2(H^1)}) + C(\|u\|_{L_T^2(H^1)} + \|\omega\|_{L_T^2} + \|\rho_0\|_{H^1} + \|u_0\|_{H^1} \\
& + \|f_0\|_{L_T^2(H^1)} + \|f_1\|_{L_T^2(L^2)} + N^{3/2}(0, T)).
\end{aligned} \tag{126}$$

7) We apply Lemma 16 for  $k = \ell = 0$ :

$$\begin{aligned}
& \|\zeta_1 D_r \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_r \rho\|_{L_T^2(L^2)} + \|\zeta_1 D_r \frac{d\rho}{dt}\|_{L_T^2(L^2)} \leq C(\|\zeta_1 D_\phi u\|_{L_T^2(H^1)} \\
& + \|u_t\|_{L_T^2(L^2)} + \|u\|_{L_T^2(H^1)} + \|\rho_0\|_{H^1} + \|f_0\|_{L_T^2(H^1)} + \|f_1\|_{L_T^2(L^2)} + N^{3/2}(0, T) + N^2(0, T)).
\end{aligned} \tag{127}$$

From (125), (126) and (127), we deduce in particular that we can estimate  $u$  in  $L_T^2(H^2)$  in the interior and  $\rho$  in  $L_T^\infty(H^1)$ . This last fact comes from the estimate  $\|\rho\|_{H^1} \leq C\|\nabla \rho\|_{L^2}$ , which holds thanks to

$$\frac{d}{dt} \int_{\Omega_F(t)} \rho(t, x) dx = 0.$$

Precisely, we have

$$\begin{aligned}
& \|\zeta_0 \nabla u\|_{L_T^2(H^1)} + \|\rho\|_{L_T^\infty(H^1)} + \|\zeta_0 \nabla \rho\|_{L_T^2(L^2)} + \left\| \nabla \frac{d\rho}{dt} \right\|_{L_T^2(L^2)} \\
& \leq \delta (\|\zeta_1 D_\phi \rho\|_{L_T^2(L^2)} + \|u\|_{L_T^2(H^2)} + \|\rho\|_{L_T^2(H^1)}) + C (\|u\|_{L_T^\infty(L^2)} + \|u\|_{L_T^2(H^1)} + \|u\|_{H_T^1(L^2)} \\
& \quad + \|\omega\|_{L_T^2} + \|u_0\|_{H^1} + \|\rho_0\|_{H^1} + \|f_0\|_{L_T^2(H^1)} + \|f_1\|_{L_T^2(L^2)} + N^{3/2}(0, T) + N^2(0, T)).
\end{aligned} \tag{128}$$

8) We apply Lemma 11 for  $k = 2$ :

$$\|u\|_{L_T^\infty(H^2)} \leq C (\|u\|_{W_T^{1,\infty}(L^2)} + \|\rho\|_{L_T^\infty(H^1)} + \|\dot{a}\|_{L_T^\infty} + \|\omega\|_{L_T^\infty} + \|f_1\|_{L_T^\infty(L^2)}). \tag{129}$$

9) We apply Lemma 17 for  $k = \ell = 0$

$$\begin{aligned}
\|\zeta_1 u\|_{L_T^2(H^2)} + \|\zeta_1 \rho\|_{L_T^2(H^1)} & \leq C \left( \left\| \zeta_1 \frac{d\rho}{dt} \right\|_{L_T^2(H^1)} + \|u_t\|_{L_T^2(L^2)} + \|u\|_{L_T^2(H^1)} + \|\rho\|_{L_T^2(L^2)} \right. \\
& \quad \left. + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} + \|f_0\|_{L_T^2(H^1)} + \|f_1\|_{L_T^2(L^2)} \right).
\end{aligned} \tag{130}$$

Using (128) in (129) and (130), we obtain estimates for  $u$  in  $L_T^\infty(H^2) \cap L_T^2(H^2)$  and for  $\rho$  in  $L_T^2(H^1)$ :

$$\begin{aligned}
& \|u\|_{L_T^\infty(H^2)} + \|u\|_{L_T^2(H^2)} + \|\rho\|_{L_T^\infty(H^1)} + \|\rho\|_{L_T^2(H^1)} + \left\| \nabla \frac{d\rho}{dt} \right\|_{L_T^2(L^2)} \\
& \leq C (\|u\|_{W_T^{1,\infty}(L^2)} + \|u\|_{L_T^2(H^1)} + \|u\|_{H_T^1(L^2)} + \left\| \frac{d\rho}{dt} \right\|_{L_T^2(L^2)} \\
& \quad + \|\dot{a}\|_{L_T^\infty} + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^\infty} + \|\omega\|_{L_T^2} + \|u_0\|_{H^1} + \|\rho_0\|_{H^1} + \|f_0\|_{L_T^2(H^1)} \\
& \quad + \|f_1\|_{L_T^\infty(L^2)} + \|f_1\|_{L_T^2(L^2)} + N^{3/2}(0, T) + N^2(0, T)).
\end{aligned} \tag{131}$$

Therefore, putting this together with (124), we get

$$\begin{aligned}
& \|u\|_{W_T^{1,\infty}(L^2)} + \|u\|_{L_T^\infty(H^2)} + \|u\|_{L_T^2(H^2)} + \|u\|_{H_T^1(H^1)} + \|\rho\|_{W_T^{1,\infty}(L^2)} + \|\rho\|_{L_T^\infty(H^1)} \\
& + \|\rho\|_{H_T^1(H^1)} + \left\| \frac{d\rho}{dt} \right\|_{L_T^2(H^1)} + \|\dot{a}\|_{W_T^{1,\infty}} + \|\dot{a}\|_{H_T^1} + \|\omega\|_{W_T^{1,\infty}} + \|\omega\|_{H_T^1} \\
& \leq C (\|f_0\|_{H_T^1(L^2)} + \|f_0\|_{L_T^2(H^1)} + \|f_1\|_{H_T^1(L^2)} + \|f_1\|_{L_T^\infty(L^2)} + \|f_2\|_{H_T^1} \\
& \quad + \|f_3\|_{H_T^1} + \|\partial_t u(0, \cdot)\|_{L^2} + \|u_0\|_{H^1} + \|\partial_t \rho(0, \cdot)\|_{L^2} + \|\rho_0\|_{H^1} \\
& \quad + |\partial_t^2 a(0)| + |a_0| + |\partial_t \omega(0)| + |\omega_0| + N^{3/2}(0, T) + N^2(0, T)).
\end{aligned} \tag{132}$$

Observe that we have added the term  $\|\rho_t\|_{L_T^2(H^1)}$  in the left hand side of this inequality. This term comes from the estimate of  $\frac{d\rho}{dt}$  in  $L_T^2(H^1)$  since the nonlinear term goes to  $N^2(0, T)$ .

10) We apply Lemma 14 for  $k = 2$

$$\begin{aligned}
& \|\zeta_0 \nabla \rho\|_{L_T^\infty(H^1)} + \|\zeta_0 \nabla \rho\|_{L_T^2(H^1)} + \|\zeta_0 \nabla u\|_{L_T^\infty(H^1)} + \|\zeta_0 \nabla u\|_{L_T^2(H^2)} + \left\| \zeta_0 \nabla \frac{d\rho}{dt} \right\|_{L_T^2(H^1)} \\
& \leq C (\|u\|_{L_T^\infty(H^1)} + \|u\|_{L_T^2(H^2)} + \|\rho_0\|_{H^2} + \|u_0\|_{H^2} + \|f_0\|_{L_T^2(H^2)} + \|f_1\|_{L_T^2(H^1)} + N^{3/2}(0, T)).
\end{aligned} \tag{133}$$

11) We apply Lemma 15 for  $k = 2$ :

$$\begin{aligned}
& \|\zeta_1 D_\phi^2 \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi^2 u\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi^2 u\|_{L_T^2(H^1)} + \left\| \zeta_1 D_\phi^2 \frac{d\rho}{dt} \right\|_{L_T^2(L^2)} \\
& \leq \delta (\|\zeta_1 D_\phi^2 \rho\|_{L_T^2(L^2)} + \|\rho\|_{L_T^2(H^1)}) + C (\|u\|_{L_T^2(H^2)} + \|\omega\|_{L_T^2} + \|\rho_0\|_{H^2} + \|u_0\|_{H^2} \\
& \quad + \|f_0\|_{L_T^2(H^2)} + \|f_1\|_{L_T^2(H^1)} + N^{3/2}(0, T)).
\end{aligned} \tag{134}$$



**12)** We apply Lemma 16 for  $k = 1, \ell = 0$ :

$$\begin{aligned} & \|\zeta_1 D_\phi D_r \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi D_r \rho\|_{L_T^2(L^2)} + \|\zeta_1 D_\phi D_r \frac{d\rho}{dt}\|_{L_T^2(L^2)} \leq C(\|\zeta_1 D_\phi^2 u\|_{L_T^2(H^1)} \\ & + \|u_t\|_{L_T^2(H^1)} + \|u\|_{L_T^2(H^2)} + \|\rho_0\|_{H^2} + \|f_0\|_{L_T^2(H^2)} + \|f_1\|_{L_T^2(H^1)} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (135)$$

Observe that the first term in the right hand side in (135) can be estimated by the third term in the left hand side of (134).

**13)** We apply Lemma 17 for  $k = 1, \ell = 0$ :

$$\begin{aligned} \|\zeta_1 D_\phi u\|_{L_T^2(H^2)} + \|\zeta_1 D_\phi \rho\|_{L_T^2(H^1)} & \leq C \left( \left\| \zeta_1 D_\phi \frac{d\rho}{dt} \right\|_{L_T^2(H^1)} + \|u_t\|_{L_T^2(H^1)} + \|u\|_{L_T^2(H^2)} + \|\rho\|_{L_T^2(H^1)} \right. \\ & \left. + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} + \|f_0\|_{L_T^2(H^2)} + \|f_1\|_{L_T^2(H^1)} \right). \end{aligned} \quad (136)$$

The term  $\|\zeta_1 \nabla D_\phi \frac{d\rho}{dt}\|_{L_T^2(L^2)}$ , which is contained in the first norm of the right hand side of (136), can be estimated with the help of the third term in the left hand side of (135) and the fourth term in the left hand side of (134).

**14)** We apply Lemma 16 for  $k = 0, \ell = 1$ :

$$\begin{aligned} & \|\zeta_1 D_r^2 \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_r^2 \rho\|_{L_T^2(L^2)} + \|\zeta_1 D_r^2 \frac{d\rho}{dt}\|_{L_T^2(L^2)} \leq C(\|\zeta_1 D_\phi D_r u\|_{L_T^2(H^1)} \\ & + \|u_t\|_{L_T^2(H^1)} + \|u\|_{L_T^2(H^2)} + \|\rho_0\|_{H^2} + \|f_0\|_{L_T^2(H^2)} + \|f_1\|_{L_T^2(H^1)} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (137)$$

The norm  $\|\zeta_1 D_\phi D_r u\|_{L_T^2(H^1)}$  will be bounded thanks to the first term in the left hand side of (136).

**15)** We apply Lemma 17 for  $k = 0, \ell = 1$ :

$$\begin{aligned} \|\zeta_1 u\|_{L_T^2(H^3)} + \|\zeta_1 \rho\|_{L_T^2(H^2)} & \leq C \left( \left\| \zeta_1 \frac{d\rho}{dt} \right\|_{L_T^2(H^2)} + \|u_t\|_{L_T^2(H^1)} + \|u\|_{L_T^2(H^2)} + \|\rho\|_{L_T^2(H^1)} \right. \\ & \left. + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} + \|f_0\|_{L_T^2(H^2)} + \|f_1\|_{L_T^2(H^1)} \right). \end{aligned} \quad (138)$$

The term  $\|\zeta_1 D^2 \frac{d\rho}{dt}\|_{L_T^2(L^2)}$ , which is contained in the first norm of the right hand side of (138), can be estimated with the help of third term in the right hand side of (137), the third in the left of (135) and the fourth in the left of (134).

Collecting expressions (133)-(138), we deduce

$$\begin{aligned} & \|u\|_{L_T^2(H^3)} + \|D^2 \rho\|_{L_T^\infty(L^2)} + \|\rho\|_{L_T^2(H^2)} + \left\| \frac{d\rho}{dt} \right\|_{L_T^2(H^2)} \\ & \leq C(\|u\|_{L_T^\infty(H^1)} + \|u\|_{L_T^2(H^2)} + \|u\|_{H_T^1(H^1)} + \|\rho\|_{L_T^2(H^1)} + \left\| \frac{d\rho}{dt} \right\|_{L_T^2(H^1)} + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} \\ & + \|u_0\|_{H^2} + \|\rho_0\|_{H^2} + \|f_0\|_{L_T^2(H^2)} + \|f_1\|_{L_T^2(H^1)} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (139)$$

Using (132), for the moment we have

$$\begin{aligned} & \|u\|_{W_T^{1,\infty}(L^2)} + \|u\|_{L_T^\infty(H^2)} + \|u\|_{L_T^2(H^3)} + \|u\|_{H_T^1(H^1)} + \|\rho\|_{W_T^{1,\infty}(H^1)} + \|\rho\|_{L_T^\infty(H^2)} \\ & + \|\rho\|_{H_T^1(H^2)} + \left\| \frac{d\rho}{dt} \right\|_{L_T^2(H^2)} + \|\dot{a}\|_{W_T^{1,\infty}} + \|\dot{a}\|_{H_T^1} + \|\omega\|_{W_T^{1,\infty}} + \|\omega\|_{H_T^1} \leq C(\|f_0\|_{H_T^1(L^2)} \\ & + \|f_0\|_{L_T^2(H^2)} + \|f_1\|_{L_T^2(H^1)} + \|f_1\|_{H_T^1(L^2)} + \|f_1\|_{L_T^\infty(L^2)} + \|f_2\|_{H_T^1} + \|f_3\|_{H_T^1} + \|\partial_t u(0, \cdot)\|_{L^2} \\ & + \|u_0\|_{H^2} + \|\partial_t \rho(0, \cdot)\|_{L^2} + \|\rho_0\|_{H^2} + |\ddot{a}(0)| + |a_0| + |\dot{\omega}(0)| + |\omega_0| + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (140)$$

In order to estimate the norm  $\|\rho\|_{W_T^{1,\infty}(H^1)}$  we have used the equation of  $\rho$  and  $\|f_0\|_{L_T^\infty(H^1)} \leq C(\|f_0\|_{L_T^2(H^2)} + \|f_0\|_{H_T^1(L^2)})$ .

**16)** We apply Lemma 13 for  $k = 1$ :

$$\begin{aligned} & \|u_t\|_{H_T^1(L^2)} + \|\rho_t\|_{H_T^1(L^2)} + \|u\|_{W_T^{1,\infty}(H^1)} + \|\ddot{a}\|_{H_T^1} + \|\dot{\omega}\|_{H_T^1} \\ & \leq C(\|u\|_{H_T^1(H^1)} + \|\rho\|_{W_T^{1,\infty}(L^2)} + \|f_0\|_{H_T^1(L^2)} + \|f_1\|_{H_T^1(L^2)} + \|f_2\|_{H_T^1} + \|f_3\|_{H_T^1} \\ & \quad + \|\partial_t u(0, \cdot)\|_{H^1} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (141)$$

Now, we regard the equation satisfied by  $u_t$  as a stationary elliptic equation (see system (71)):

$$\begin{cases} -2\mu\nabla \cdot (\epsilon(u_t)) - \mu'\nabla \cdot ((\nabla \cdot u_t)Id) = -u_{tt} - p^0\nabla\rho_t + f_{1,t} & \text{in } (0, T) \times \Omega_F(t), \\ u_t = (\ddot{a} + \dot{\omega} \wedge (x - a) + \omega \wedge (\omega \wedge (x - a)) - (u_S \cdot \nabla)u)\mathbf{1}_{\partial\Omega_S(t)} & \text{in } (0, T) \times \partial\Omega_F(t). \end{cases}$$

Then, we have

$$\|u\|_{H_T^1(H^2)} \leq C(\|u\|_{H_T^2(L^2)} + \|\rho\|_{H_T^1(H^1)} + \|f_1\|_{H_T^1(L^2)} + \|\dot{a}\|_{H_T^1} + \|\omega\|_{H_T^1} + N^2(0, T)). \quad (142)$$

**17)** We apply Lemma 11 for  $k = 3$ :

$$\|u\|_{L_T^\infty(H^3)} \leq C(\|u\|_{W_T^{1,\infty}(H^1)} + \|\rho\|_{L_T^\infty(H^2)} + \|\dot{a}\|_{L_T^\infty} + \|\omega\|_{L_T^\infty} + \|f_1\|_{L_T^\infty(H^1)}). \quad (143)$$

Combining the three estimates (141)-(143), we get

$$\begin{aligned} & \|u\|_{W_T^{1,\infty}(H^1)} + \|u\|_{L_T^\infty(H^3)} + \|u\|_{H_T^2(L^2)} + \|u\|_{H_T^1(H^2)} + \|\rho_t\|_{H_T^1(L^2)} + \|\ddot{a}\|_{H_T^1} + \|\dot{\omega}\|_{H_T^1} \\ & \leq C(\|u\|_{H_T^1(H^1)} + \|\rho\|_{W_T^{1,\infty}(L^2)} + \|\rho\|_{L_T^\infty(H^2)} + \|\rho\|_{H_T^1(H^1)} + \|\dot{a}\|_{L_T^\infty} + \|\dot{a}\|_{H_T^1} \\ & \quad + \|\omega\|_{H_T^1} + \|\omega\|_{L_T^\infty} + \|f_0\|_{H_T^1(L^2)} + \|f_1\|_{H_T^1(L^2)} + \|f_1\|_{L_T^\infty(H^1)} \\ & \quad + \|f_2\|_{H_T^1} + \|f_3\|_{H_T^1} + \|\partial_t u(0, \cdot)\|_{H^1} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (144)$$

**18)** We apply Lemma 14 for  $k = 3$ :

$$\begin{aligned} & \|\zeta_0\nabla\rho\|_{L_T^\infty(H^2)} + \|\zeta_0\nabla\rho\|_{L_T^2(H^2)} + \|\zeta_0\nabla u\|_{L_T^\infty(H^2)} + \|\zeta_0\nabla u\|_{L_T^2(H^3)} + \|\zeta_0\frac{d\rho}{dt}\|_{L_T^2(H^2)} \\ & \leq C(\|u\|_{L_T^\infty(H^2)} + \|u\|_{L_T^2(H^3)} + \|\rho_0\|_{H^3} + \|u_0\|_{H^3} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} + N^{3/2}(0, T)). \end{aligned} \quad (145)$$

**19)** We apply Lemma 15 for  $k = 3$ :

$$\begin{aligned} & \|\zeta_1 D_\phi^3 \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi^3 u\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi^3 u\|_{L_T^2(H^1)} + \|\zeta_1 D_\phi^3 \frac{d\rho}{dt}\|_{L_T^2(L^2)} \\ & \leq \delta(\|\zeta_1 D_\phi^3 \rho\|_{L_T^2(L^2)} + \|\rho\|_{L_T^2(H^1)}) + C(\|u\|_{L_T^2(H^3)} + \|\omega\|_{L_T^2} + \|\rho_0\|_{H^3} + \|u_0\|_{H^3} \\ & \quad + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} + N^{3/2}(0, T)). \end{aligned} \quad (146)$$

**20)** We apply Lemma 16 for  $k = 2, \ell = 0$ :

$$\begin{aligned} & \|\zeta_1 D_\phi^2 D_r \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi^2 D_r \rho\|_{L_T^2(L^2)} + \|\zeta_1 D_\phi^2 D_r \frac{d\rho}{dt}\|_{L_T^2(L^2)} \leq C(\|\zeta_1 D_\phi^3 u\|_{L_T^2(H^1)} \\ & \quad + \|u_t\|_{L_T^2(H^2)} + \|u\|_{L_T^2(H^3)} + \|\rho_0\|_{H^3} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (147)$$

Observe that the norm  $\|\zeta_1 D_\phi^3 u\|_{L_T^2(H^1)}$  is estimated with the third term in the left of (146).

**21)** We apply Lemma 17 for  $k = 2, \ell = 0$ :

$$\begin{aligned} \|\zeta_1 D_\phi^2 u\|_{L_T^2(H^2)} + \|\zeta_1 D_\phi^2 \rho\|_{L_T^2(H^1)} &\leq C \left( \left\| \zeta_1 D_\phi^2 \frac{d\rho}{dt} \right\|_{L_T^2(H^1)} + \|u_t\|_{L_T^2(H^2)} + \|u\|_{L_T^2(H^3)} + \|\rho\|_{L_T^2(H^2)} \right. \\ &\quad \left. + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} \right). \end{aligned} \quad (148)$$

As for (138), the term  $\|\zeta_1 \nabla D_\phi^2 \frac{d\rho}{dt}\|_{L_T^2(L^2)}$  (which is contained in the first norm of the right hand side of (148)) is bounded by the third term in the left of (147) and the fourth term in the left of (146).

**22)** We apply Lemma 16 for  $k = 1, \ell = 1$ :

$$\begin{aligned} \|\zeta_1 D_\phi D_r^2 \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_\phi D_r^2 \rho\|_{L_T^2(L^2)} + \|\zeta_1 D_\phi D_r^2 \frac{d\rho}{dt}\|_{L_T^2(L^2)} &\leq C(\|\zeta_1 D_\phi^2 D_r u\|_{L_T^2(H^1)} \\ &\quad + \|u_t\|_{L_T^2(H^2)} + \|u\|_{L_T^2(H^3)} + \|\rho_0\|_{H^3} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (149)$$

The first term in the left of (148) will serve to absorb the first in the right of (149).

**23)** We apply Lemma 17 for  $k = 1, \ell = 1$ :

$$\begin{aligned} \|\zeta_1 D_\phi u\|_{L_T^2(H^3)} + \|\zeta_1 D_\phi \rho\|_{L_T^2(H^2)} &\leq C \left( \left\| \zeta_1 D_\phi \frac{d\rho}{dt} \right\|_{L_T^2(H^2)} + \|u_t\|_{L_T^2(H^2)} + \|u\|_{L_T^2(H^3)} + \|\rho\|_{L_T^2(H^2)} \right. \\ &\quad \left. + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} \right). \end{aligned} \quad (150)$$

Observe that  $\|\zeta_1 D^2 D_\phi \frac{d\rho}{dt}\|_{L_T^2(L^2)}$  is estimated with the help of the third term in the left of (149), the third in the left of (147) and the fourth in the left of (146).

**24)** We apply Lemma 16 for  $k = 0, \ell = 2$ :

$$\begin{aligned} \|\zeta_1 D_r^3 \rho\|_{L_T^\infty(L^2)} + \|\zeta_1 D_r^3 \rho\|_{L_T^2(L^2)} + \|\zeta_1 D_r^3 \frac{d\rho}{dt}\|_{L_T^2(L^2)} &\leq C(\|\zeta_1 D_\phi D_r^2 u\|_{L_T^2(H^1)} \\ &\quad + \|u_t\|_{L_T^2(H^2)} + \|u\|_{L_T^2(H^3)} + \|\rho_0\|_{H^3} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (151)$$

The term  $\|\zeta_1 D_\phi D_r^2 u\|_{L_T^2(H^1)}$  is estimated with the help of the first term in the left of (150).

**25)** We apply Lemma 17 for  $k = 0, \ell = 2$ :

$$\begin{aligned} \|\zeta_1 u\|_{L_T^2(H^4)} + \|\zeta_1 \rho\|_{L_T^2(H^3)} &\leq C \left( \left\| \zeta_1 \frac{d\rho}{dt} \right\|_{L_T^2(H^3)} + \|u_t\|_{L_T^2(H^2)} + \|u\|_{L_T^2(H^3)} + \|\rho\|_{L_T^2(H^2)} \right. \\ &\quad \left. + \|\dot{a}\|_{L_T^2} + \|\omega\|_{L_T^2} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} \right). \end{aligned} \quad (152)$$

The term  $\|\zeta_1 D^3 \frac{d\rho}{dt}\|_{L_T^2(L^2)}$  (contained in the first norm of the right hand side of (152)) is estimated with the third terms in the left hand sides of (151), (149) and (147) and the fourth term in the left of (146).

Combining estimates (145)-(152), we get:

$$\begin{aligned} &\|u\|_{L_T^2(H^4)} + \|\rho\|_{L_T^2(H^3)} + \|D^3 \rho\|_{L_T^\infty(L^2)} + \|D^3 \frac{d\rho}{dt}\|_{L_T^2(L^2)} \\ &\leq C(\|u\|_{L_T^2(H^3)} + \|u\|_{L_T^\infty(H^2)} + \|u\|_{H_T^1(H^2)} + \|\rho\|_{L_T^2(H^2)} + \|\frac{d\rho}{dt}\|_{L_T^2(H^2)} + \|\dot{a}\|_{L_T^2} \\ &\quad + \|\omega\|_{L_T^2} + \|u_0\|_{H^3} + \|\rho_0\|_{H^3} + \|f_0\|_{L_T^2(H^3)} + \|f_1\|_{L_T^2(H^2)} + N^{3/2}(0, T) + N^2(0, T)). \end{aligned} \quad (153)$$

Together with (144), this estimate yields

$$\begin{aligned}
& \|u\|_{W_T^{1,\infty}(H^1)} + \|u\|_{L_T^\infty(H^3)} + \|u\|_{L_T^2(H^4)} + \|u\|_{H_T^2(L^2)} + \|u\|_{H_T^1(H^2)} \\
& + \|\rho\|_{L_T^2(H^3)} + \|D^3\rho\|_{L_T^\infty(L^2)} + \|\rho\|_{H_T^2(L^2)} + \|D^3\frac{d\rho}{dt}\|_{L_T^2(L^2)} + \|\ddot{a}\|_{H_T^1} + \|\dot{\omega}\|_{H_T^1} \\
& \leq C(\|u\|_{L_T^2(H^3)} + \|u\|_{H_T^1(H^1)} + \|\rho\|_{H_T^1(H^1)} + \|\rho\|_{L_T^2(H^2)} + \|\rho\|_{W_T^{1,\infty}(L^2)} + \|\rho\|_{L_T^\infty(H^2)} + \|\frac{d\rho}{dt}\|_{L_T^2(H^2)}) \quad (154) \\
& + \|\dot{a}\|_{L_T^\infty} + \|\dot{a}\|_{H_T^1} + \|\omega\|_{H_T^1} + \|\omega\|_{L_T^\infty} + \|f_0\|_{L_T^2(H^3)} + \|f_0\|_{H_T^1(L^2)} + \|f_1\|_{L_T^2(H^2)} + \|f_1\|_{H_T^1(L^2)} \\
& + \|f_1\|_{L_T^\infty(H^1)} + \|f_2\|_{H_T^1} + \|f_3\|_{H_T^1} + \|u_0\|_{H^3} + \|\rho_0\|_{H^3} + \|\partial_t u(0, \cdot)\|_{H^1} + N^{3/2}(0, T) + N^2(0, T).
\end{aligned}$$

Finally, we combine this estimate with (140) and we conclude inequality (64). Similarly as we did for  $\|\rho_t\|_{L^\infty(H^1)}$  (just after (140)), the norm  $\|\rho_t\|_{L^\infty(H^2)}$  is estimated using the equation of  $\rho$  and thanks to  $f_0 \in L_T^\infty(H^2)$ .

## References

- [1] H. BEIRÃO DA VEIGA, *On the existence of strong solutions to a coupled fluid-structure evolution problem*, J. Math. Fluid Mech., **6** (2004), no. 1, 21–52.
- [2] M. BOULAKIA, *Existence of weak solutions for an interaction problem between an elastic structure and a compressible viscous fluid*, J. Math. Pures et Appliquées, **84** (2005), no. 11, 1515–1554.
- [3] M. BOULAKIA, *Existence of weak solutions for the three dimensional motion of an elastic structure in an incompressible fluid*, J. Math. Fluid Mech., **9** (2007), no. 2, 262–294.
- [4] A. CHAMOLLE, B. DESJARDINS, M.J. ESTEBAN, C. GRANDMONT, *Existence of weak solutions for the unsteady interaction of a viscous fluid with an elastic plate*, J. Math. Fluid Mech., **7** (2005), no. 3, 368–404.
- [5] C. CONCA, J. SAN MARTIN, M. TUCSNAK, *Existence of solutions for the equations modelling the motion of a rigid body in a viscous fluid*, Comm. Partial Differential Equations, **25** (2000), no. 5-6, 1019–1042.
- [6] D. COUTAND, S. SHKOLLER, *Motion of an elastic solid inside an incompressible viscous fluid*, Arch. Ration. Mech. Anal., **176** (2005), no. 1, 25–102.
- [7] B. DESJARDINS, M. J. ESTEBAN, *On weak solutions for fluid-rigid structure interaction: compressible and incompressible models*, Comm. Partial Differential Equations, **25** (2000), no. 7-8, 1399–1413.
- [8] B. DESJARDINS, M. J. ESTEBAN, C. GRANDMONT, P. LE TALLEC, *Weak solutions for a fluid-elastic structure interaction model*, Rev. Mat. Complut. **14** (2001), no. 2, 523–538.
- [9] E. FEIREISL, A. NOVOTNÝ, H. PETZELTOVÁ, *On the existence of globally defined weak solutions to the Navier-Stokes equations*, J. Math. Fluid Mech. **3** (2001), no. 4, 358–392.
- [10] E. FEIREISL, *On the motion of rigid bodies in a viscous compressible fluid*, Arch. Ration. Mech. Anal. **167** (2003), no. 4, 281–308.
- [11] E. FEIREISL, *Dynamics of viscous compressible fluids*, Oxford Science Publications, Oxford (2004).
- [12] C. GRANDMONT, Y. MADAY, *Existence for an unsteady fluid-structure interaction problem*, M2AN Math. Model. Numer. Anal., **34** (2000), no. 3, 609–636.
- [13] D. HOFF, *Global solutions of the Navier-Stokes equations for multidimensional compressible flow with discontinuous initial data*, J. Differential Equations **120** (1995), no. 1, 215–254.

- [14] D. HOFF, *Strong convergence to global solutions for multidimensional flows of compressible, viscous fluids with polytropic equations of state and discontinuous initial data*, Arch. Rational Mech. Anal. **132** (1995), no. 1, 1–14.
- [15] P.-L. LIONS, *Existence globale de solutions pour les équations de Navier-Stokes compressibles isentropiques*, C. R. Acad. Sci. Paris Sr. I Math. **316** (1993), no. 12, 1335–1340.
- [16] P.L. LIONS, *Mathematical Topics in Fluid Mechanics*, Oxford Science Publications, Oxford (1996).
- [17] A. MATSUMURA, T. NISHIDA, *The initial value problem for the equations of motion of viscous and heat-conductive gases*, J. Math. Kyoto Univ., **20** (1980), no. 1, 67–104.
- [18] A. MATSUMURA, T. NISHIDA, *Initial-boundary value problems for the equations of motion of general fluids*, Computing methods in applied sciences and engineering, V (Versailles, 1981) North-Holland, Amsterdam (1982), 389–406.
- [19] J. SAN MARTIN, V. STAROVOITOV, M. TUCSNAK, *Global weak solutions for the two dimensional motion of several rigid bodies in an incompressible viscous fluid*, Arch. Rational Mech. Anal., **161** (2002), no. 2, 93–112.
- [20] T. TAKAHASHI, *Analysis of strong solutions for the equations modeling the motion of a rigid-fluid system in a bounded domain*, Adv. Differential Equations, **8** (2003), no. 12, 1499–1532.
- [21] A. TANI, *On the first initial-boundary value problem of compressible viscous fluid motion*, Publ. RIMS, Kyoto Univ. **13** (1977), 193–253.
- [22] R. TEMAM, *Navier-Stokes equations. Theory and numerical analysis. Studies in Mathematics and its applications*, **2**; North Holland Publishing Co., Amsterdam-New York-Oxford (1977).
- [23] E. ZEIDLER, *Nonlinear functional analysis and its applications. I. Fixed-point theorems*, Translated from the German by Peter R. Wadsack. Springer-Verlag, New York, 1986.